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# Fractions of $K^+K^-$ pairs with $K^+$ or $K^-$ mesons from the resonance decays in pp interactions at 24 GeV/c

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#### Abstract

The Coulomb  $K^+K^-$  correlations might be sensitive to non point-like  $K^+K^-$  pair production due to influence of kaons from decays of resonances with relatively small widths. The relative fractions of  $K^+K^-$  pairs with the  $K^+$  or  $K^-$  mesons from decays of the  $K^*(892)$ ,  $\Lambda(1520)$  and  $\phi$ produced in pp interactions at 24 GeV/c are determined using the experimental or estimated cross sections of these resonances.

In an attempt to estimate the number of produced  $K^+K^-$  atoms,  $A_{K^+K^-}$ , in proton-nucleus collisions at 24 GeV/c in the DIRAC experiment at the CERN PS by studying the Coulomb  $K^+K^$ correlations, it might be important to take into account non point-like  $K^+K^-$  pair production. Indeed if one of kaons in  $K^+K^-$  pair is the decay product of relatively narrow resonance, such as  $\Lambda(1520)$ or  $\phi$ , the diistance between K mesons in their center-of-mass system (with  $c\tau \approx 12.6$  fm and 46 fm for, respectively,  $\Lambda(1520)$  and  $\phi$ ) is although smaller but still not drastically different from the Bohr radius (109 fm) of  $A_{K^+K^-}$ . In this case the Coulomb correlations are weaker than for the point-like pair production. In order to take this into account it is necessary to know the fractions of  $K^+K^-$  pairs with  $K^+$  or  $K^-$  from such resonance decays. In this note, we determine the fractions of such  $K^+K^$ pairs with  $K^+$  or  $K^-$  from the  $K^*(892)$  (with  $c\tau \approx 4.2$  fm),  $\Lambda(1520)$  and  $\phi$  decays in pp interactions at 24 GeV/c.

The cross section of  $K^+K^-$  pair production in pp interactions at 24 GeV/c has been estimated in [1]:

$$\sigma(K^+K^-) = (0.500 \pm 0.028) \text{ mb.}$$
(1)

The cross sections of  $\phi$ ,  $\Lambda(1520)$  and  $K^*(892)$  in pp interactions at 24 GeV/c, either measured or estimated, are collected in Table 1.

Production mechanisms of the  $\Lambda(1520)$  and  $\Sigma^{*+}(1385)$  in proton fragmentation processes are very similar (as it is evident from the corresponding quark diagrams in Fig. 1). Therefore the  $\Lambda(1520)$ cross section was estimated in [4] from the measured  $\Sigma^{*+}(1385)$  cross section in pp interactions at 24 GeV/c,  $(0.276 \pm 0.026)$  mb [6], assuming the same ratio of the  $\Lambda(1520)$  and  $\Sigma^{*+}(1385)$  cross sections,  $(0.56 \pm 0.10)/(0.67 \pm 0.08) = 0.836 \pm 0.180$ , as measured in pp interactions at 400 GeV/c [7].

The  $\bar{K}^{*0}(892)$  cross section at 24 GeV/c is assumed [5] equal to the measured  $K^{*-}(892)$  cross section [3], as it is supported by results of pp experiment at 400 GeV/c [7], where  $\sigma(\bar{K}^{*0}(892)) = (2.96 \pm 0.54)$  mb and  $\sigma(K^{*-}(892)) = (2.87 \pm 0.39)$  mb. The  $K^{*0}(892)$  cross section at 24 GeV/c was estimated as explained in details in [5].

Other numerous resonances, some of which are observed only in the phase-wave analyses, either have large widths or small branching ratios into the final states with kaons and/or small production

Table 1: Mesured cross sections of  $\phi$  [2],  $K^{*+}(892)$  and  $K^{*-}(892)$  [3] and estimated cross sections of  $\Lambda(1520)$  [4],  $K^{*0}(892)$  and  $\bar{K}^{*0}(892)$  [5] in pp interactions at 24 GeV/c.

Resonance	$\phi$	$\Lambda(1520)$	$K^{*+}(892)$	$K^{*-}(892)$	$K^{*0}(892)$	$\bar{K}^{*0}(892)$
$\sigma$ (mb)	$0.158 \pm 0.035$	$0.231 \pm 0.054$	$0.66\pm0.06$	$0.19\pm0.09$	$0.44\pm0.04$	$0.19\pm0.04$



Figure 1: Quark diagrams of the  $\Lambda(1520)$  and  $\Sigma^{*+}(1385)$  production in proton fragmentation processes.

rates (such as  $f_1(1285)$  with  $\Gamma \approx 24 \text{ MeV}/c^2$  and  $Br(K\bar{K}) \approx 9\%$  or f'(1525) with  $\Gamma \approx 73 \text{ MeV}/c^2$  and  $Br(K\bar{K}) \approx 89\%$ ) and therefore are not relevant to the present analysis.

The fraction of  $K^+K^-$  pairs with the  $K^-$  from the  $\bar{K}^{*0}(892)$  decay is

$$f_{\bar{K}^{*0}}(K^+K^-) = Br(\bar{K}^{*0} \to K^-\pi^+)\sigma(\bar{K}^{*0})/\sigma(K^+K^-) = 25.3 \pm 5.5)\%.$$
(2)

The cross sections of  $K^{*+}(892)$  and  $K^{*0}(892)$  are much higher than those of  $K^{*-}(892)$  and  $\bar{K}^{*0}(892)$ . However it was explained [1] by proton fragmentation into  $K^+$  (or  $K^{*+}(892)$ ,  $K^{*0}(892)$ ) and hyperons. The associated production of such  $K^{*+}(892)$ ,  $K^{*0}(892)$  with the  $K\bar{K}$  pair production is presumably small. Therefore it is reasonably safe to assume that the fraction of  $K^+K^-$  pairs with the  $K^+$  from the  $K^{*0}(892)$  decay is the same as in (2):

$$f_{K^{*0}}(K^+K^-) = f_{\bar{K}^{*0}}(K^+K^-) = (25.3 \pm 5.5)\%.$$
(3)

The fractions (3) are quite large. However, the  $K^+K^-$  pairs with  $K^+$  or  $K^-$  from decays of the  $K^{*0}(892)$  or  $\bar{K}^{*0}(892)$  can not contribute significantly to non point-like  $K^+K^-$  pair production because of relatively large  $K^*(892)$  width.

The fraction of  $K^+K^-$  pairs with the  $K^-$  from the  $\Lambda(1520)$  decay amounts to

$$f_{\Lambda(1520)}(K^+K^-) = \epsilon \cdot Br(\Lambda(1520) \to pK^-) \cdot \frac{\sigma(\Lambda(1520))}{\sigma(K^+K^-)} = (8.07 \pm 2.00)\%, \tag{4}$$

where the branching ratio  $Br(\Lambda(1520) \to pK^-) = \frac{1}{2}(0.45 \pm 0.01)$  and factor  $\epsilon$  takes into account that proton may fragment either into  $\Lambda(1520)$  and  $K^+$  or into  $\Lambda(1520)$  and  $K^0\pi^+$ . If the latter process is ignored, then  $\epsilon = 1$ . If production rates of both processes are equal, then  $\epsilon = 1/2$ . The  $\epsilon$  value in (4) was determined using ratio of the measured cross sections of similar processes  $\sigma(K^+\Lambda) = (1.274 \pm 0.040)$  mb and  $\sigma(K^0\Lambda) = (0.820 \pm 0.028)$  mb [1]. This gave  $\epsilon = 0.777 \pm 0.036$ . Presumably small contribution of associated production of  $\Lambda(1520)$  with the "central"  $K\bar{K}$  pairs was ignored. With the known  $\phi$  and  $K^+K^-$  pair cross sections and branching ratio  $Br(\phi \to K^+K^-) = 0.489 \pm 0.005$ , the fraction of  $K^+K^-$  pairs with  $K^+$  and  $K^-$  mesons from the  $\phi$  decay is

$$f_{\phi}(K^{+}K^{-}) = Br(\phi \to K^{+}K^{-}) \cdot \frac{\sigma(\phi)}{\sigma(K^{+}K^{-})} = (15.45 \pm 3.53)\%.$$
(5)

However, if dominent mechanism of the  $\phi$  production is similar to the  $J/\psi$  production (as shown in Fig. 2), then  $K^+K^-$  pairs with both  $K^+$  and  $K^-$  from the  $\phi$  decay can not contribute to non point-like  $K^+K^-$  pair production.



Figure 2: Quark diagram of the  $\phi$  production similar to the  $J/\psi$  production.

The  $K^+K^-$  pairs with one kaon from the  $\phi$  decay and another from other sources (see quark diagrams in Fig. 3) could contibute to non point-like  $K^+K^-$  pair production. Such mechanism of the  $\phi$ production associated with pair of strange particles (dominantly kaons) has been shown to be supressed in pp interactions at 24 GeV/c, with estimated cross section  $\sigma_{ass}(\phi) = (0.018 \pm 0.004)$  mb [2]. Then the fraction of  $K^+K^-$  pairs with  $K^+$  (or  $K^-$ ) mesons from such  $\phi$  decay and associated  $K^-$  (or  $K^+$ ) mesons is

$$f_{\phi}(K^{+}K^{-}) = Br(\phi \to K^{+}K^{-}) \cdot \frac{\sigma_{ass}(\phi)}{\sigma(K^{+}K^{-})} = (1.8 \pm 0.4)\%, \tag{6}$$

assuming equal contribution of the  $\phi K^+ K^-$ ,  $\phi K^+ \bar{K}^0$ ,  $\phi K^- K^0$  and  $\phi K^0 \bar{K}^0$  final states.



Figure 3: Quark diagrams of the  $\phi$  production in association with two kaons.

Thus, in summary, main contribution to non point-like  $K^+K^-$  pair production arises from the  $\Lambda(1520)$  decays. Smaller contribution is due to the  $\phi$  decays, when the  $\phi$  is produced in association with the pair of kaons.

However, this concusion must be treated with great caution, since the fraction of  $\phi$  associated with any strange particle in pp interactions at 400 GeV/c is found to be  $0.91\pm0.11$  [7], in excellent agreement with OZI selection rule [8] and in contradiction with much smaller value of  $\sigma_{ass}(\phi)/\sigma(\phi) = 0.11\pm0.04$ at 24 GeV/c [2]. If this difference is not due to energy dependence, but can be explained by poor statistics and absence of particle identification in the bubble chamber experiment at 24 GeV/c, one is forced to admit that influence of  $\phi$  on non point-like  $K^+K^-$  pair production at 24 GeV/c can be comparable with that of  $\Lambda(1520)$ . These observations can be important in attempt to estimate the number of  $K^+K^-$  atoms in the data already accumulated by the DIRAC experiment.

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