

Precise Measurements of $\pi\pi$ Scattering Lengths from Kaon Decays

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Kaon Physics in the New Millennium

Since about a decade a new generation of high-statistics kaon experiments are in operation

NA48, KTeV, KLOE, ISTRA+, E787/949, NA62, ...



Significant progress in many fields in the last few years:

- Direct CP violation (ε'/ε)
- Precision determination of $|V_{us}|$
- **Tests of Chiral Perturbation Theory** ← **This Talk**
- Search for New Physics phenomena (LFV, ...)

Outline

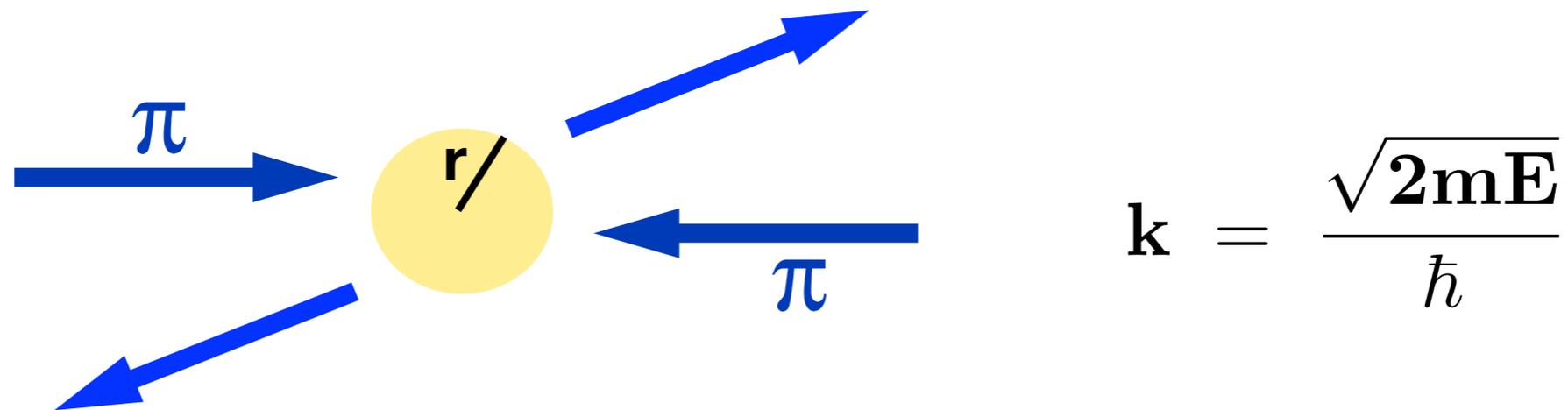
Introduction

Measurement of $\pi\pi$ scattering lengths

- **Wigner-cusp in $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$**
- **Measurement of $K^\pm \rightarrow \pi^+ \pi^- e^\pm \nu$ (K_{e4}) decays**
- **Future Prospects**

Conclusions

$\pi\pi$ Scattering Lengths



At low energy $kr \ll 1$: **S-wave** dominates scattering amplitude.
Isospin $I = 0, 2$ because of Bose statistics.

- Scattering matrix $\mathbf{S}|\pi\pi\rangle = e^{2i\delta}|\pi\pi\rangle$ parametrized by two phases:

$$\delta_{0,2} = -\mathbf{a}_{0,2} \cdot \mathbf{k} + \mathcal{O}(k^2)$$

- At low energy **S-wave scattering lengths $\mathbf{a}_0, \mathbf{a}_2$** are essential parameters of **Chiral Perturbation Theory (ChPT)**.

(In the following: $a_{0,2}$ quoted in units of m_π)

Chiral Symmetry Breaking

Chiral symmetry: (considering only u, d quarks)

- Vacuum ground state not invariant under chiral symmetry.

⇒ $SU(2)_L \otimes SU(2)_R$ becomes $SU(2)_{L+R}$.

⇒ Three massless **Goldstone bosons** are created.

(π^+, π^-, π^0)

Spontaneous symmetry breaking:

- Assumption: u, d, s quarks massless.

⇒ Lagrange function \mathcal{L}_{QCD} invariant under $SU(2)_L \otimes SU(2)_R$ transformations.

Chiral Symmetry Breaking

Symmetry breaking by finite quark masses:

- Symmetry broken by quark masses $m_u, m_d \neq 0$
⇒ Pions acquire masses.

- Gell-Mann, Oakes, Renner:

(PR 175 (1968) 2195)

$$M_\pi^2 \approx \frac{1}{F_\pi^2} (m_u + m_d) |\langle 0 | q\bar{q} | 0 \rangle|$$

with: $F_\pi =$ pion decay constant = 92.4 MeV

and: $\langle 0 | q\bar{q} | 0 \rangle =$ **quark condensate.**

- Stern *et al.*: $\langle 0 | q\bar{q} | 0 \rangle$ not necessarily $\neq 0$ for $m_u, m_d \neq 0$,
(NP B457 (1995) 513) as expansion in the quark masses is:

$$M_\pi^2 = \frac{1}{F_\pi^2} (m_u + m_d) |\langle 0 | q\bar{q} | 0 \rangle| - \frac{\bar{l}_3}{32\pi^2 F_\pi^2} M^4 + \mathcal{O}(M^6)$$

with $\bar{l}_3 =$ coupling constant, could be unexpectedly large.

Theory predictions for a_0 and a_2

- $a_0, a_2 = 0$ for $m_d, m_u = 0$ (chiral limit)
- Scattering lengths a_0, a_2 are directly connected to m_π :

$$a_0 \sim \frac{7 m_\pi^2}{32\pi F_\pi^2} = 0.16$$

$$a_2 \sim \frac{-m_\pi^2}{16\pi F_\pi^2} = -0.045$$

(Weinberg, PRL 17 (1966) 216)

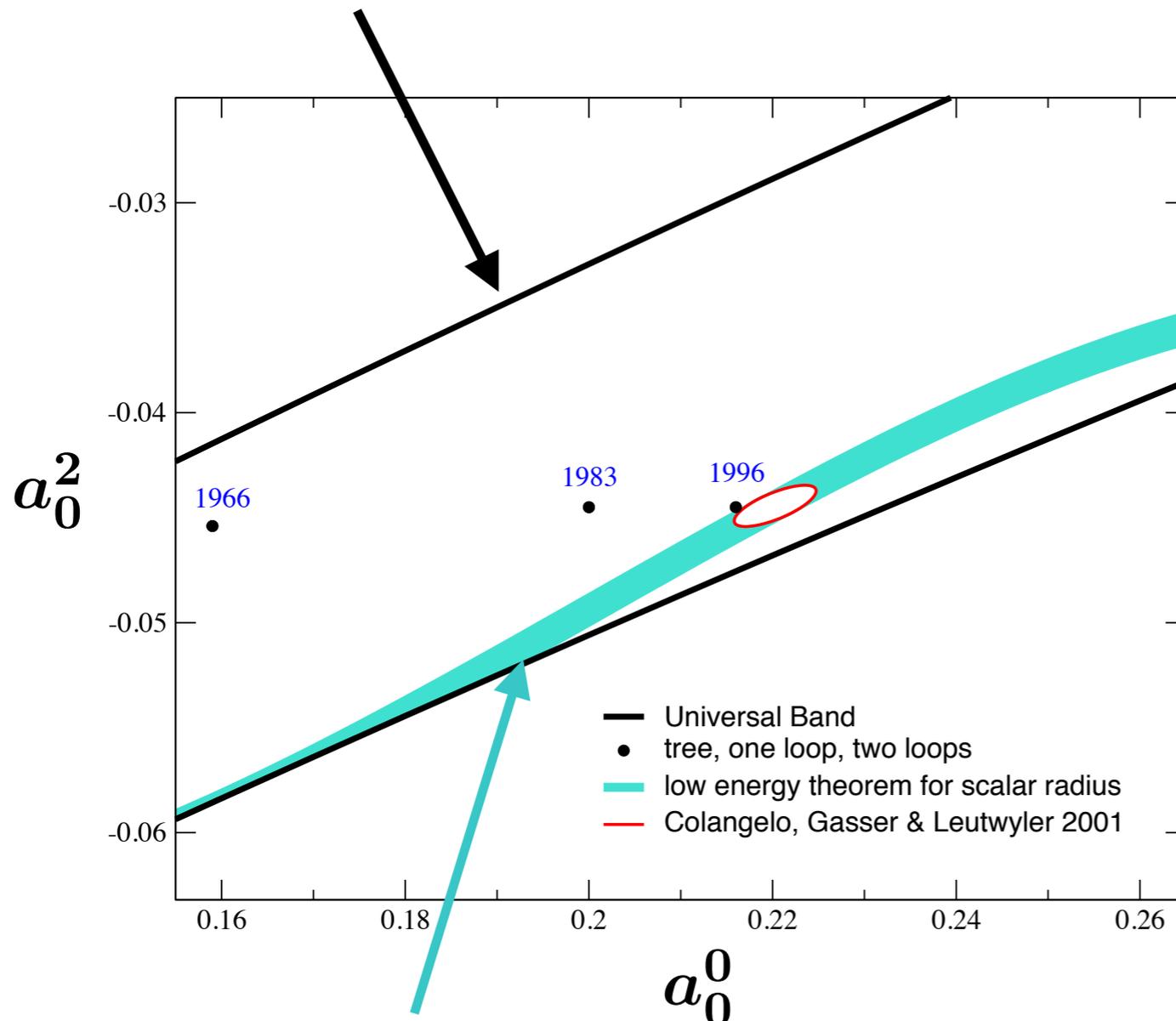
- Precise prediction within Chiral Perturbation Theory:

$$a_0 = 0.220 \pm 0.005$$

$$a_2 = -0.0444 \pm 0.0010$$

(Colangelo, Gasser, Leutwyler, PRL 86 (2001) 5008)

Universal band (from Roy equations)



ChPT constraint:

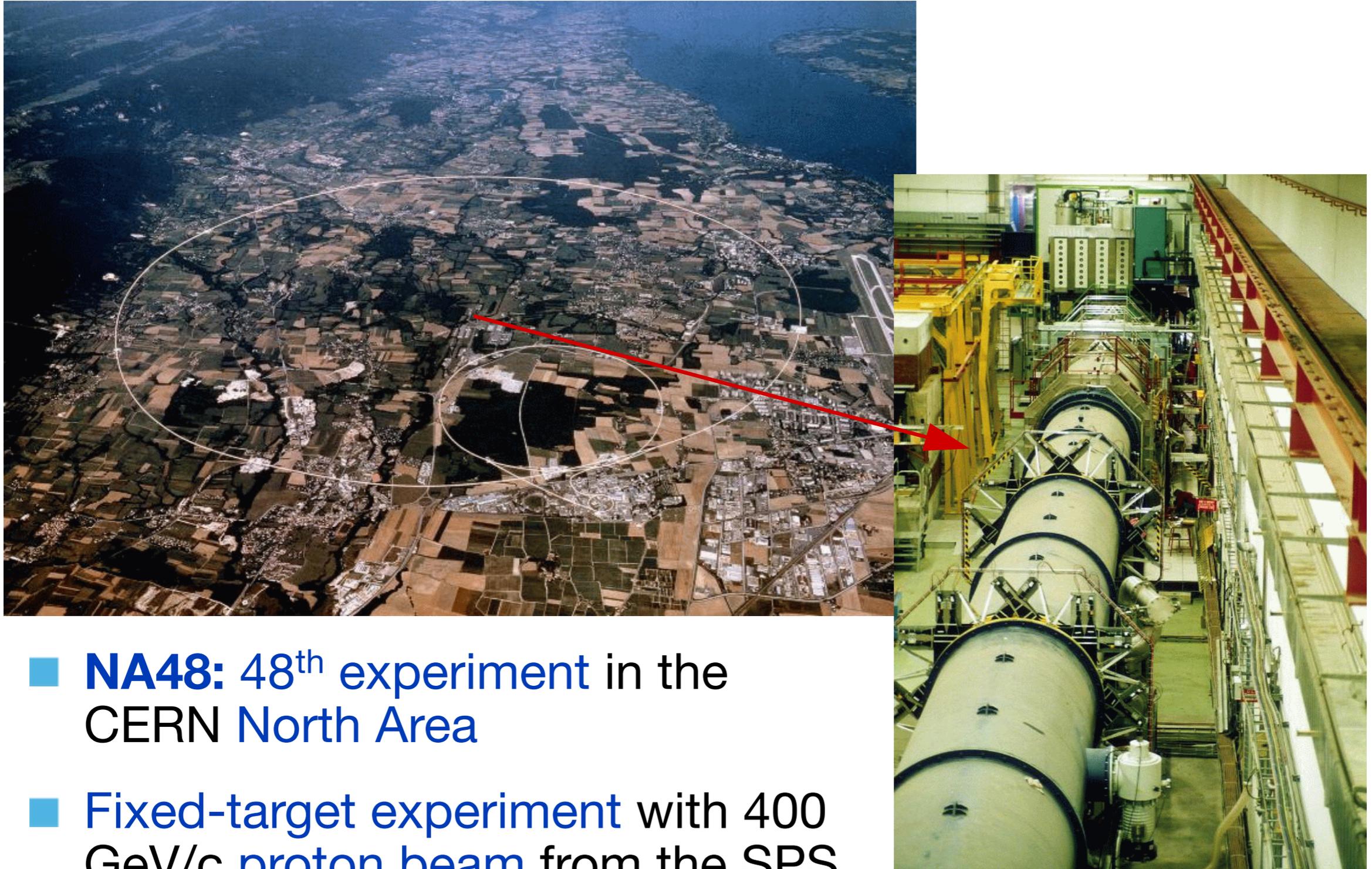
$$a_2 = -0.0444(8) + 0.236 (a_0 - 0.22) - 0.61 (a_0 - 0.22)^2 - 9.9 (a_0 - 0.22)^3$$

Measuring $\pi\pi$ Scattering Lengths

Three kinds of measurements have been performed:

- **Pionium lifetime $(\pi^+\pi^-)_{\text{atom}}$:** Measurement of $|a_2 - a_0|$
→ DIRAC experiment
- **Cusp in $K \rightarrow \pi\pi\pi$ ($K_{3\pi}$) decays:** Measurement of $a_2 - a_0$, a_2
 - $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ → NA48/2: ~60 million events
 - $K_L \rightarrow \pi^0 \pi^0 \pi^0$ → KTeV + NA48/2: (70+100) million events
- **$K^\pm \rightarrow \pi^+ \pi^- e^\pm \nu$ (K_{e4}) decays:** Measurement of a_0 , a_2
 - S118 (Geneva-Saclay, 1977): ~30 000 events
 - BNL E685 (2003): ~400 000 events
 - NA48/2 (2009): ~1.1 million events**

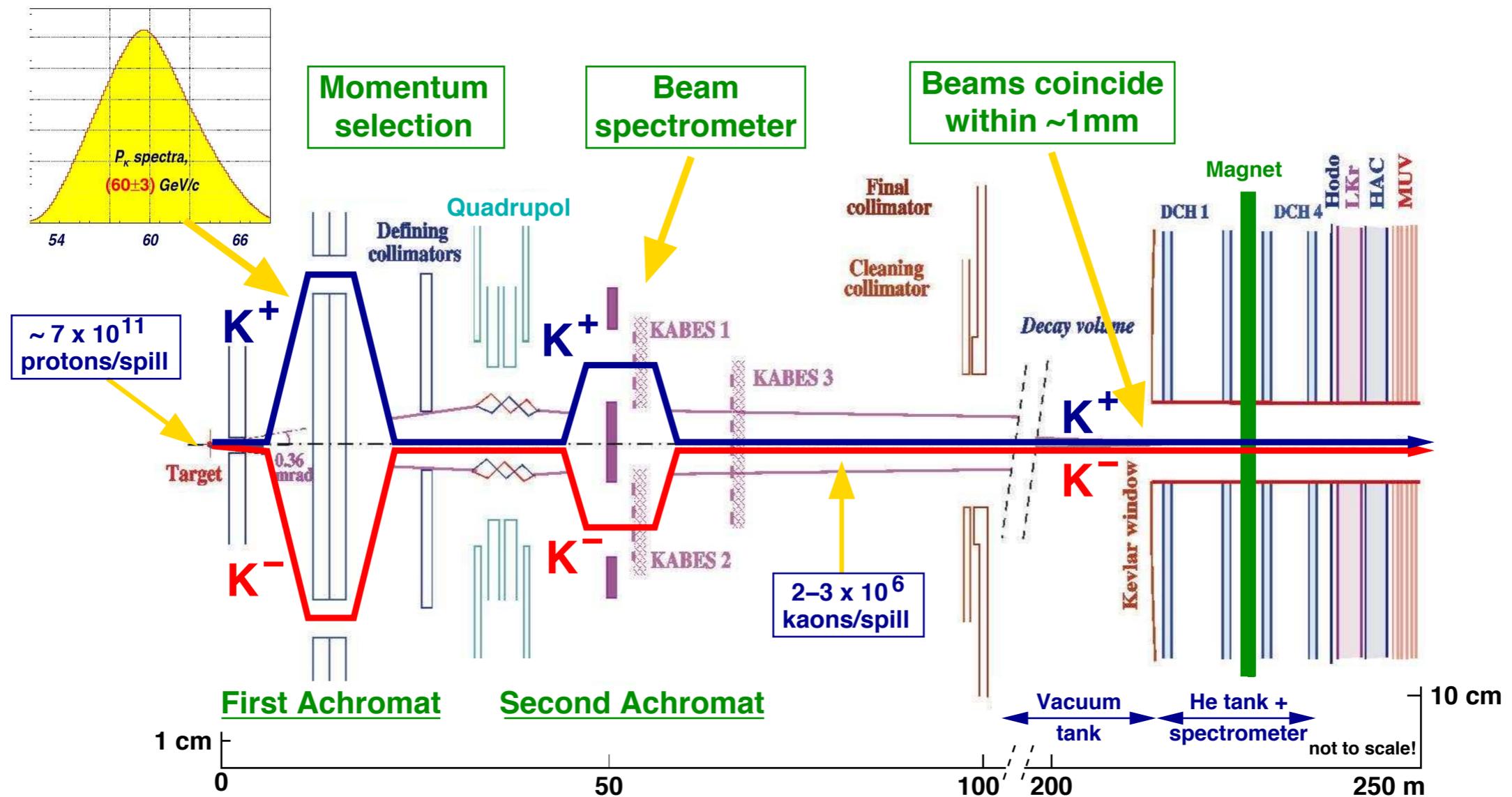
The NA48/2-Experiment



- **NA48:** 48th experiment in the CERN North Area
- Fixed-target experiment with 400 GeV/c proton beam from the SPS

NA48/2 in 2003/2004

- Simultaneous K^+ and K^- beams with $p_{K^\pm} = (60 \pm 3) \text{ GeV}/c$.



- Trigger:**
- 3 charged tracks or
 - 1 charged track + missing p_T
- Efficiencies > 99 %

NA48/2 Detector

Main detector components:

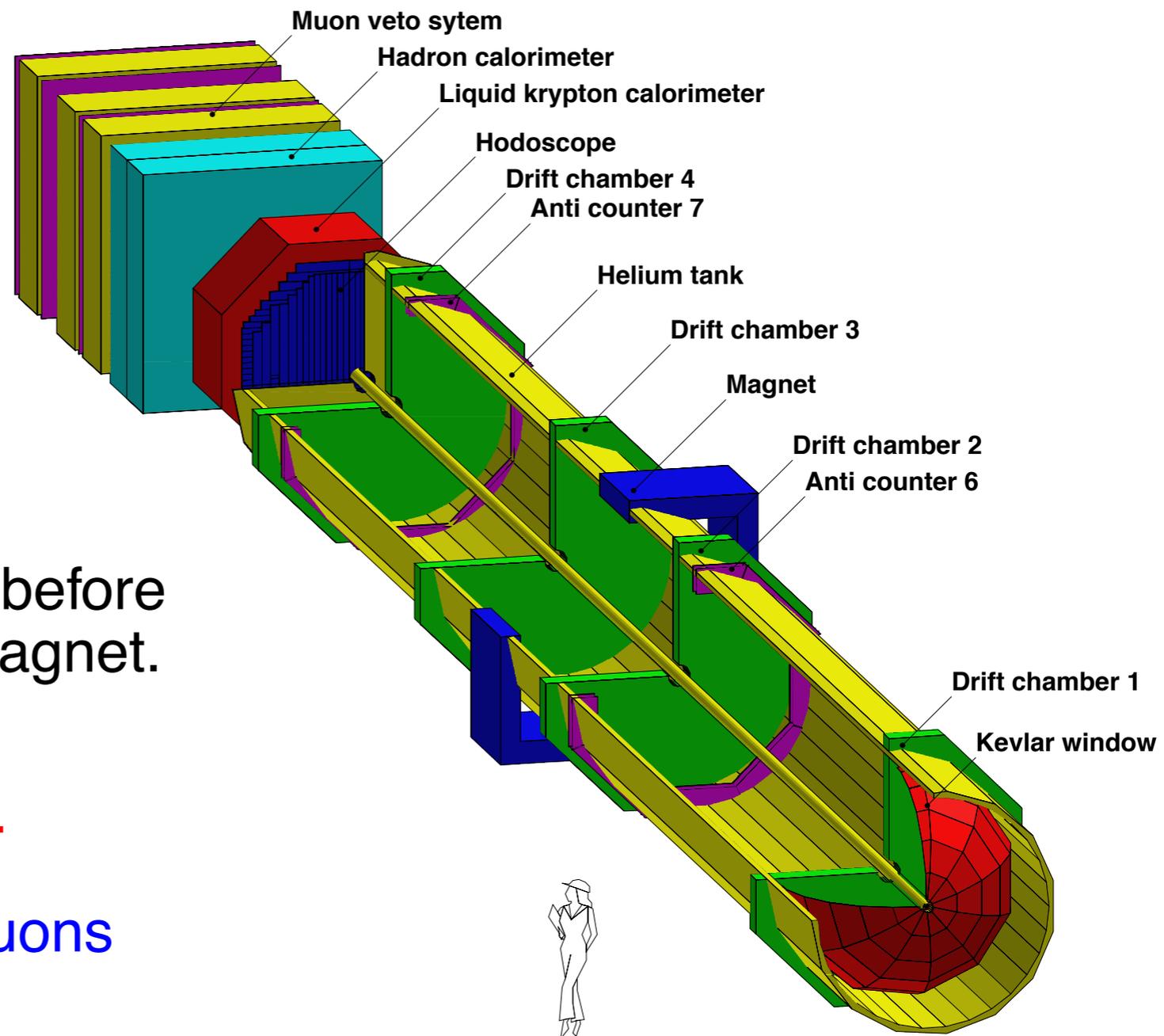
■ Magnet spectrometer

■ Two drift chambers each before and after spectrometer magnet.

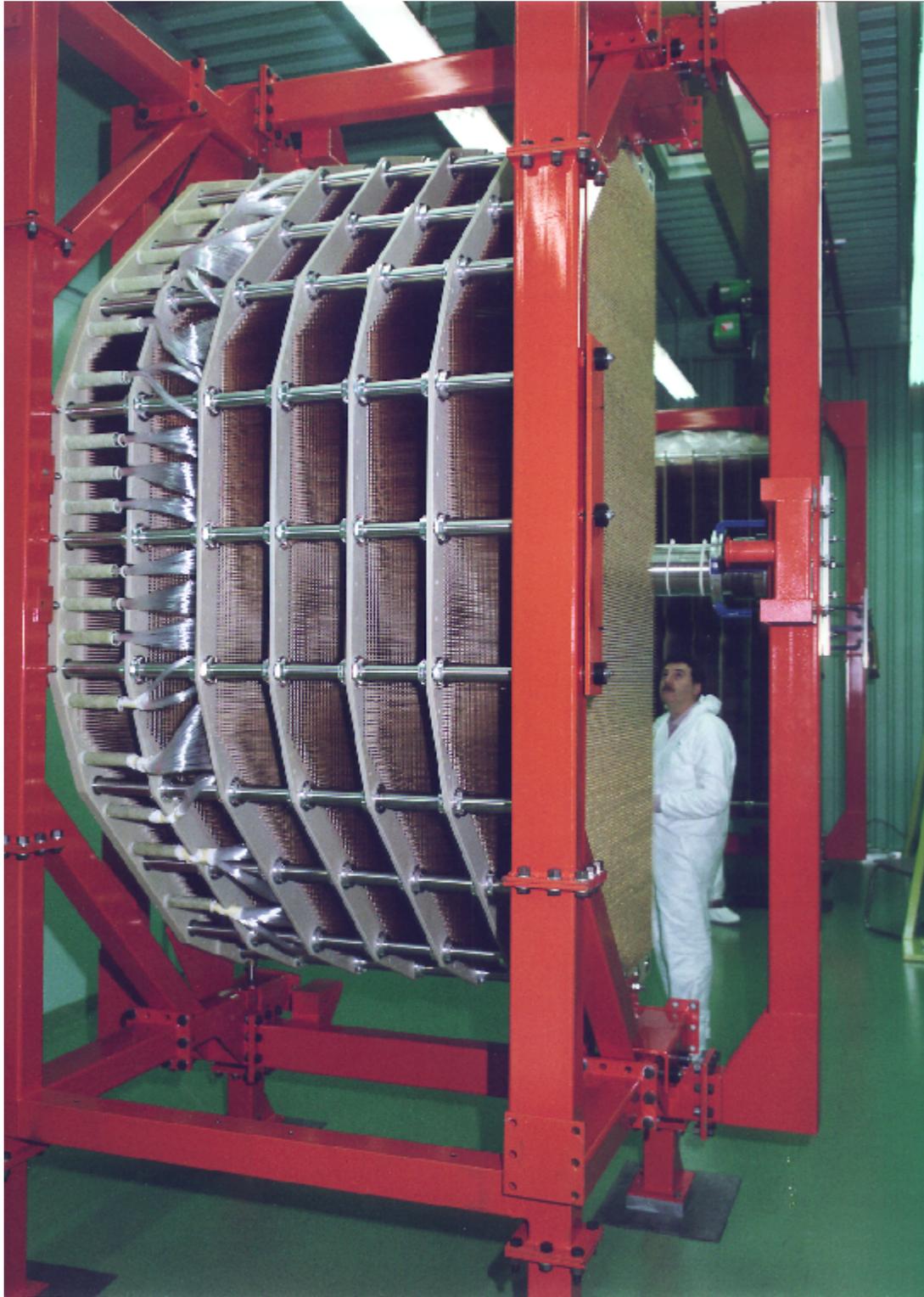
■ Momentum resolution:
 $\leq 1\%$ for $20 \text{ GeV}/c$ tracks.

■ Anti-counters for photons, muons

■ Liquid Krypton Calorimeter



Liquid-Krypton Calorimeter (LKr)



■ Structure:

13212 cells of $2 \times 2 \text{ cm}^2$ along beam axis in $\sim 10 \text{ m}^3$ liquid krypton.

■ Energy resolution:

$$\frac{\Delta E}{E} = \frac{3.2\%}{\sqrt{E[\text{GeV}]}} \oplus \frac{90 \text{ MeV}}{E} \oplus 0.42\%$$

$\Rightarrow \sim 1\%$ for 20 GeV photons.

■ Position resolution:

$$\sigma_x = \sigma_y = \frac{4.2 \text{ mm}}{\sqrt{E[\text{GeV}]}} \oplus 0.6 \text{ mm}$$

$\Rightarrow \sim 1 \text{ mm}$ for 20 GeV photons.

Cusp in $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ Decays

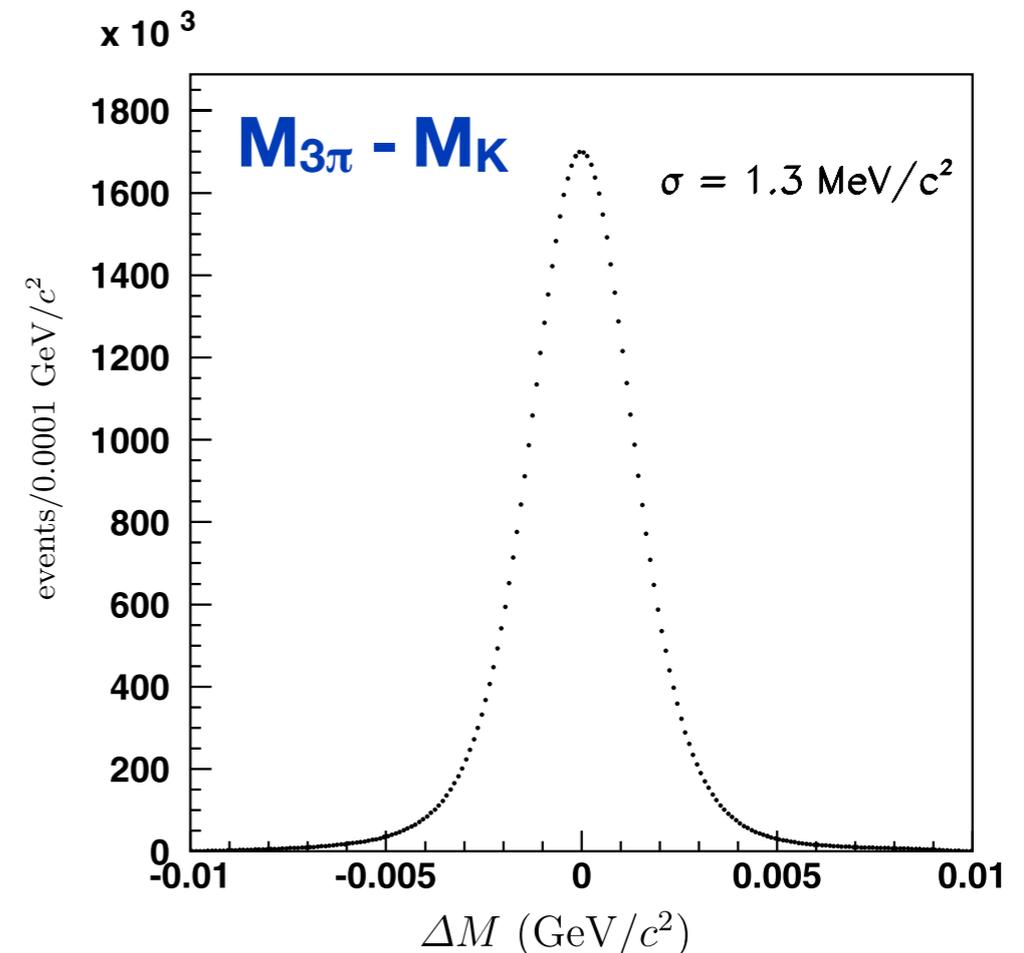
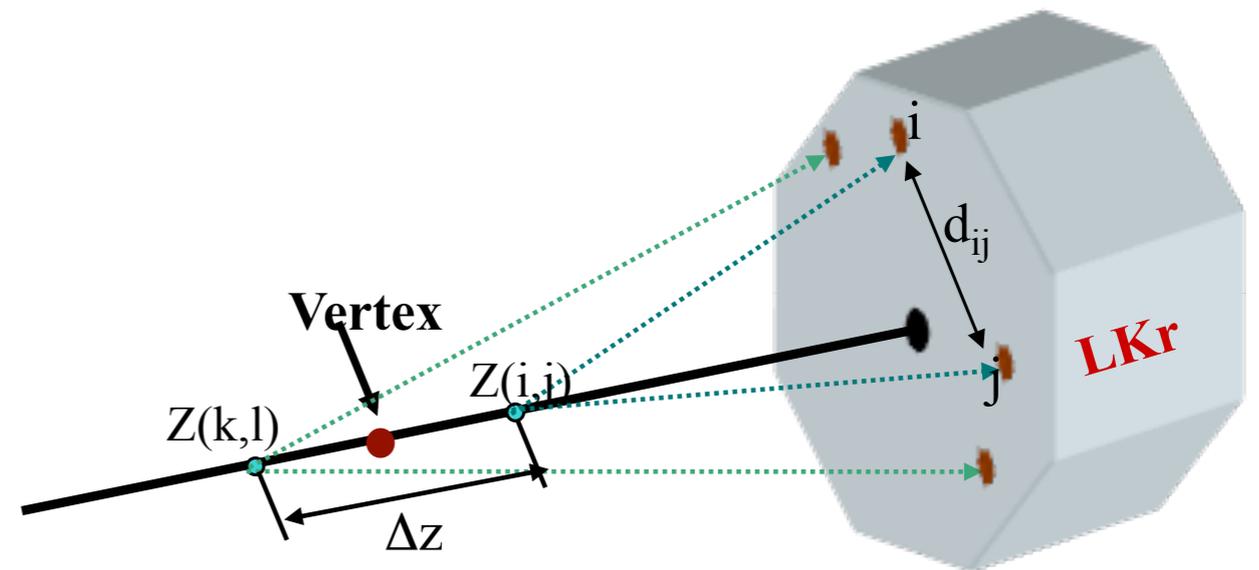
Reconstruction of $K \rightarrow \pi^\pm \pi^0 \pi^0$ Decays

$K \rightarrow \pi^\pm \pi^0 \pi^0$ selection:

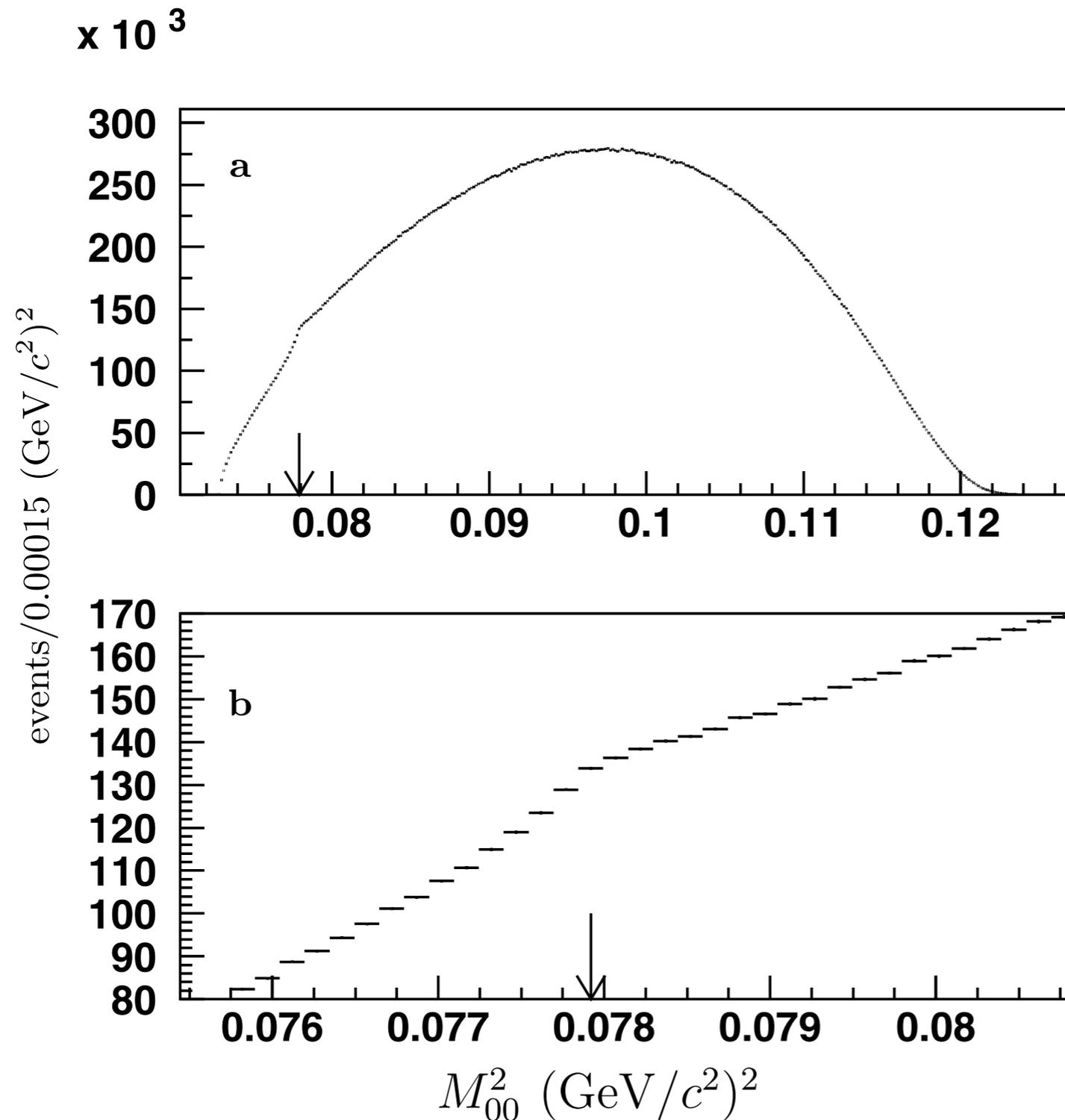
- 1 charged track + 4 e.m. calorimeter clusters
- $\pi^0 \rightarrow \gamma\gamma$ selection: consider all 3 pairings and minimize vertex difference Δz
- **invariant $\pi^0 \pi^0$ mass $M(\pi^0 \pi^0)$:** only calorimeter and vertex information used

→ **60 million events**

(mass resolution 1.3 MeV,
negligible background)



Cusp in $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ Decays



$M(\pi^0\pi^0)$ distribution:

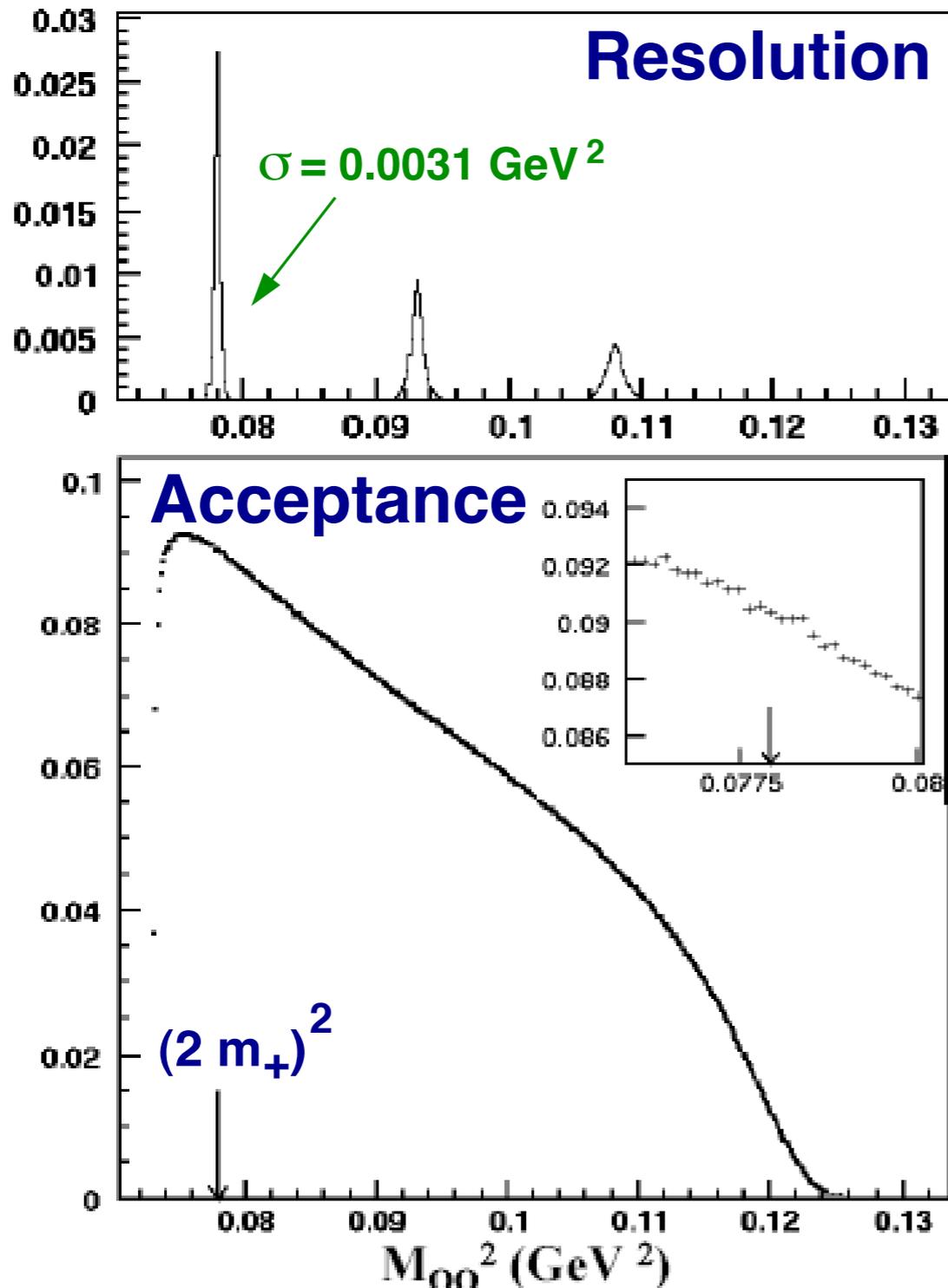
Clear cusp at

$$M(\pi^0\pi^0) = 2 m(\pi^\pm)$$

(were expecting peak from ponium formation)



$K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ Acceptance and Resolution



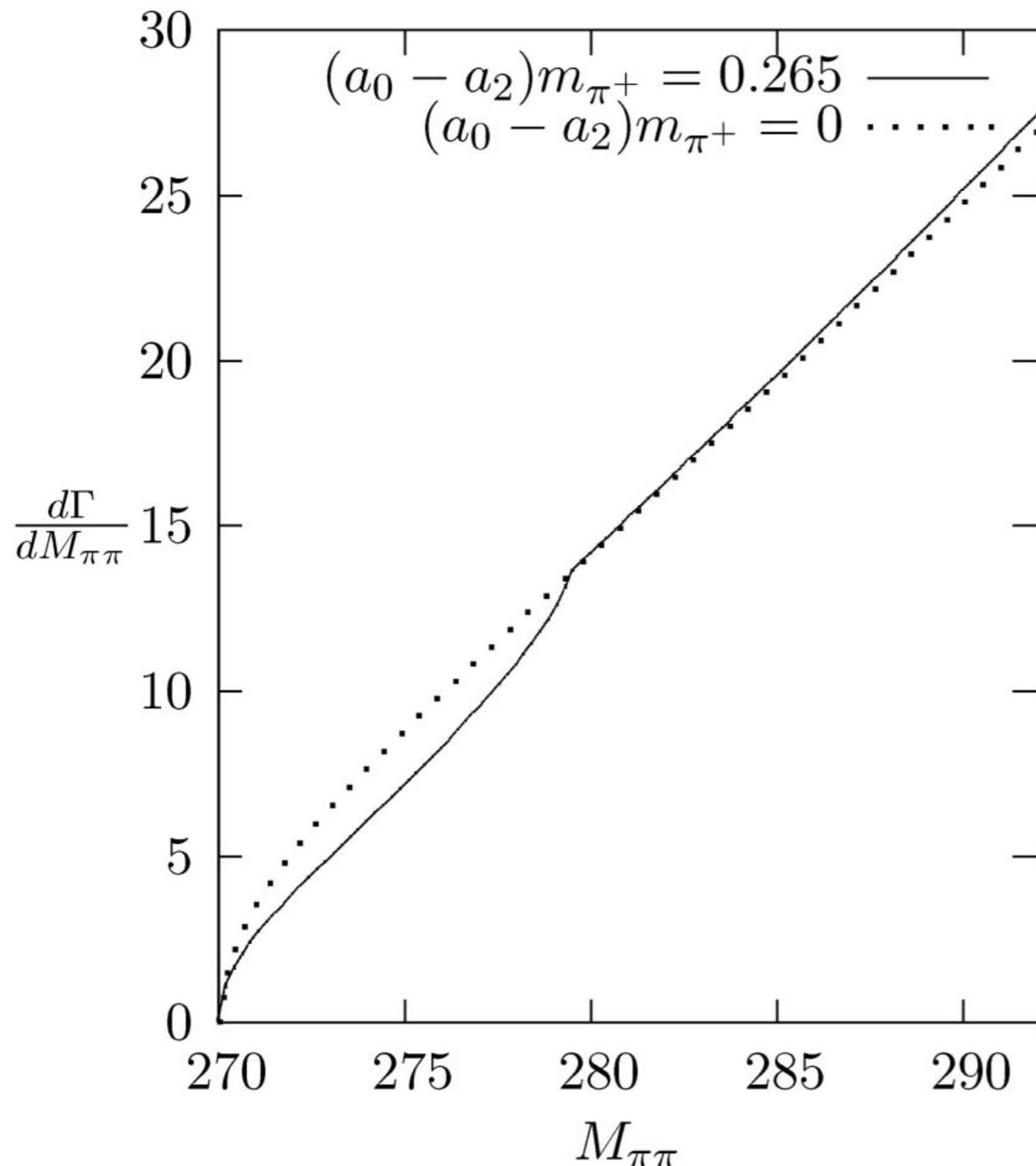
Resolution on $m_{\pi^0 \pi^0}^2$:

- Determined by **LKR resolution.**
- Best at **low $m_{\pi^0 \pi^0}^2$**
(due to kinematical constraints).
- At $m_{\pi^0 \pi^0}^2 = (2 m_{\pi^+})^2$:
 $\sigma = 0.0031 \text{ GeV}^2$

Acceptance:

- Acceptance \approx **linearly varying** around $(2 m_{\pi^+})^2$.
- Modelled by **Monte Carlo simulation.**

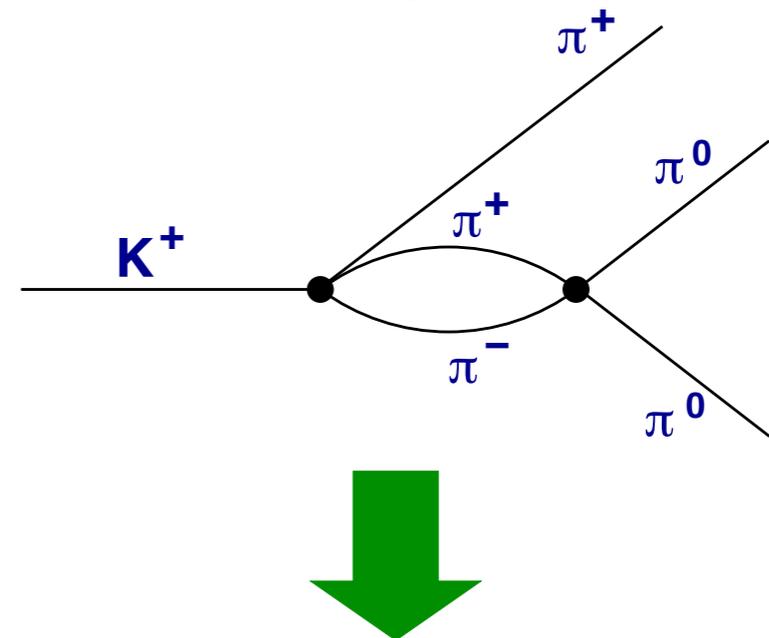
Cusp Explanation



N. Cabibbo (2004):

(PRL 93 (2004) 121801)

Rescattering of $\mathbf{K}^+ \rightarrow \pi^+ \pi^+ \pi^-$
to $\mathbf{K}^+ \rightarrow \pi^+ \pi^0 \pi^0$.



Measurement of $\pi\pi$ scattering
length $a_0 - a_2$ possible!

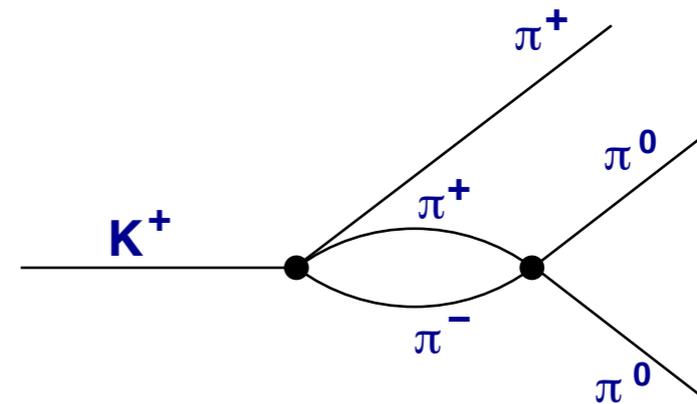
(Effect already predicted much earlier:
Budini, Fonda, PRL 6 (1961) 419)

a_0 - a_2 from $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ Decays

■ $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ decays:

$K^+ \rightarrow \pi^+ \pi^+ \pi^-$ amplitude

contributes to $K^+ \rightarrow \pi^+ \pi^0 \pi^0$ via charge exchange $\pi^+ \pi^- \rightarrow \pi^0 \pi^0$.



■ Matrix element: $\mathcal{M}(K^+ \rightarrow \pi^+ \pi^0 \pi^0) = \mathcal{M}_0 + \mathcal{M}_1$

with: \mathcal{M}_0 = "unperturbed" amplitude

\mathcal{M}_1 = rescattered amplitude

■ Above threshold: $s_\pi = m_{\pi^0 \pi^0}^2 > 4m_{\pi^+}^2$ \mathcal{M}_1 imaginary

$$\mathcal{M}_1 = -\frac{2}{3} (a_0 - a_2) m_{\pi^+} \mathcal{M}_{K^+ \rightarrow \pi^+ \pi^+ \pi^-}^{\text{threshold}} \cdot \pi \sqrt{4m_{\pi^+}^2 / s_\pi - 1}$$

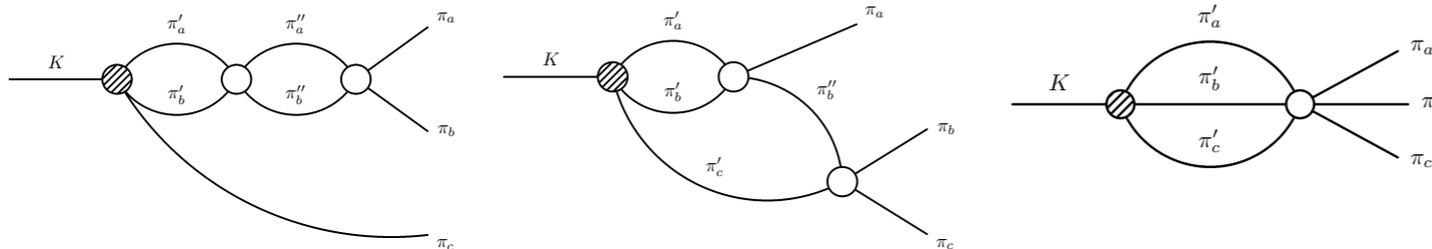
Below threshold: $s_\pi < 4m_{\pi^+}^2$ \mathcal{M}_1 real and negative

$$\mathcal{M}_1 = -\frac{2}{3} (a_0 - a_2) m_{\pi^+} \mathcal{M}_{K^+ \rightarrow \pi^+ \pi^+ \pi^-}^{\text{threshold}} \cdot (-i)\pi \sqrt{1 - 4m_{\pi^+}^2 / s_\pi}$$

Theoretical Approach (CI)

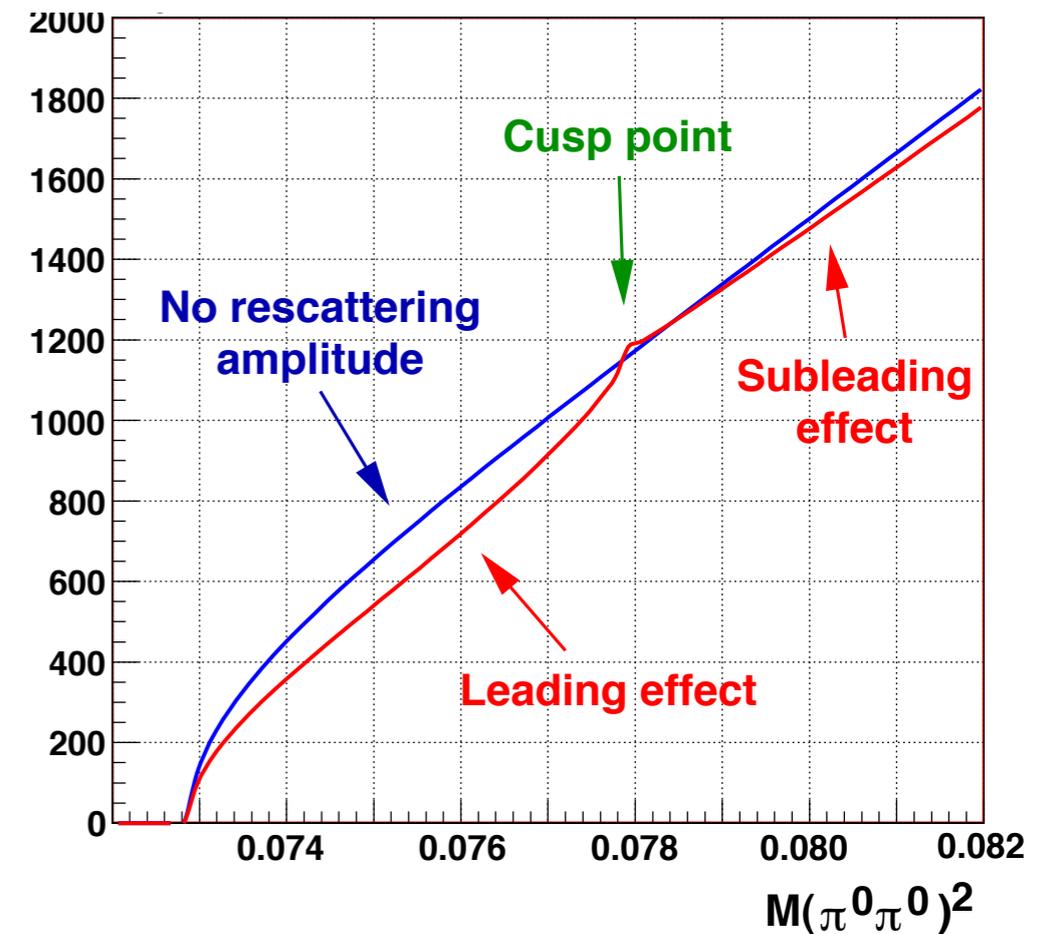
Second order computations: (Cabibbo, Isidori, JHEP03 (2005) 21)

- Other rescattering corrections: $\pi^0\pi^0 \rightarrow \pi^0\pi^0$, $\pi^+\pi^0 \rightarrow \pi^+\pi^0$, ...
- Two-loop level $\mathcal{O}(a_i^2)$ corrections:



⇒ Separate determination of a_2 possible.

Also: Coverage of all $K \rightarrow \pi\pi\pi$ decays (e.g. $K_L \rightarrow \pi^0\pi^0\pi^0$).



- $\mathcal{O}(a_i^3)$ corrections.
- Radiative corrections.

⇒ About 5% uncertainty on $a_0 - a_2$.

Theoretical Approach (BB)

Approach by the Bern-Bonn group:

- based on an **effective non-relativistic lagrangian**
- **different structure** of the expansion (w.r.t. CI)
- simultaneous fitting of neutral and charged amplitudes to **extract Dalitz plot slope parameters** (modified w.r.t. PDG parametrization)
- **electromagnetic effects** and **radiative corrections** outside the cusp point are included

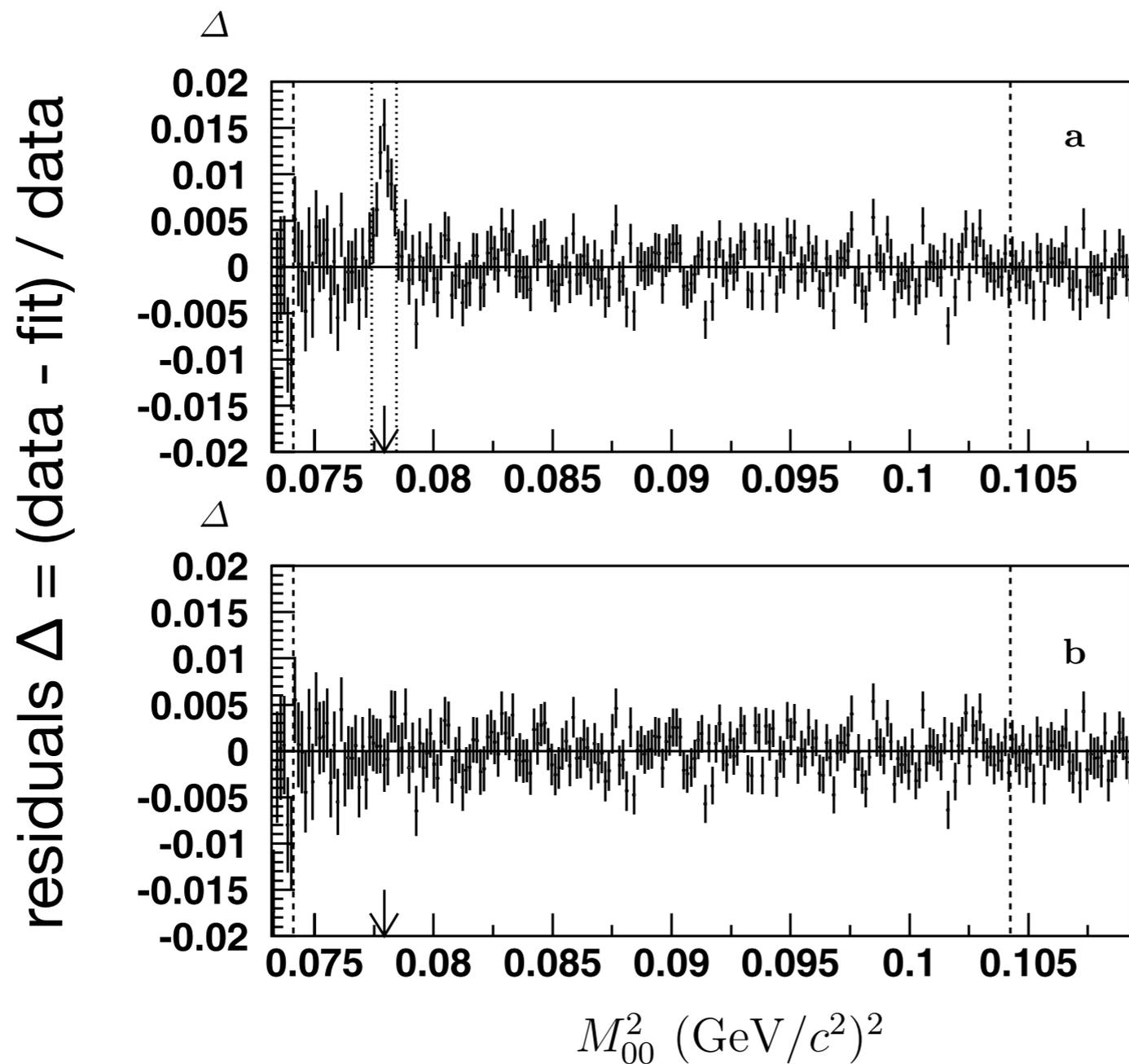
(Colangelo, Gasser, Kubis, Rusetsky, PLB 638 (2006) 187;

Bissinger, Fuhrer, Gasser, Kubis, Rusetsky, PLB 659 (2008) 576; NPH B806 (2009) 178)

→ provides so far most complete description of rescattering effect

Fit to the $M(\pi^0\pi^0)$ Spectrum (BB Model)

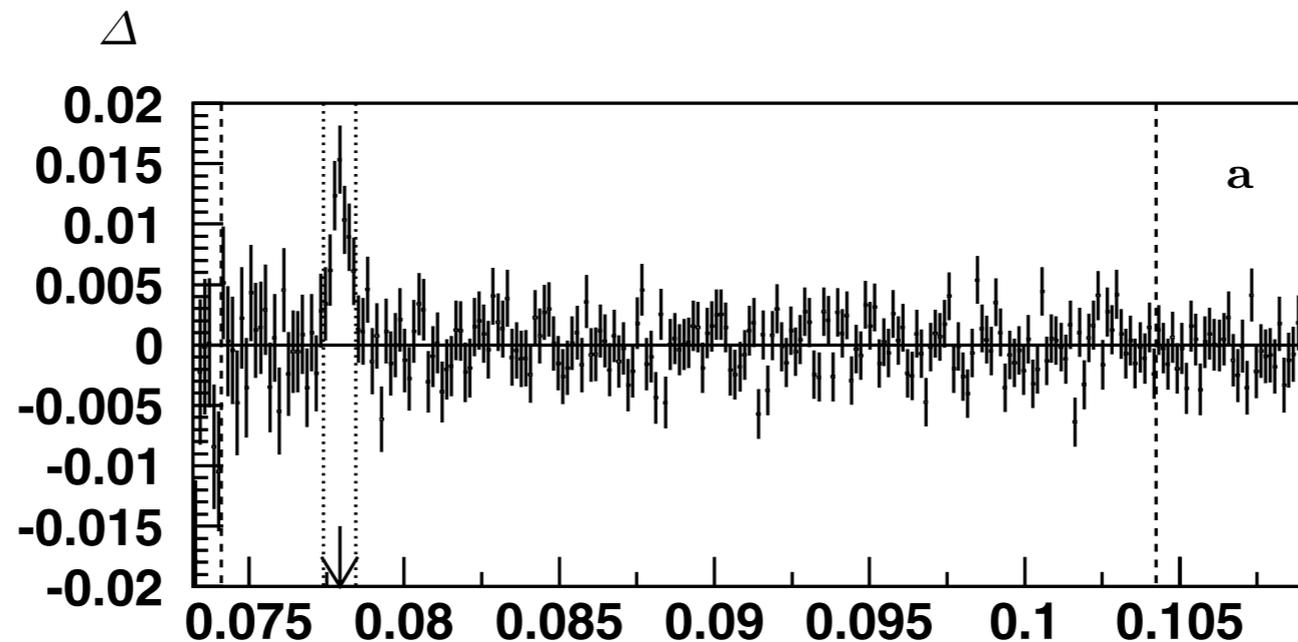
Free fit parameters: a_0 - a_2 , a_2 , Dalitz plot parameters, normalizations
(fit also includes $K \rightarrow \pi^\pm \pi^+ \pi^-$ decays)



7 bins around
cusp/pionium excluded

pionium fraction f_{atom}
left free in the fit

Brief Excursion: Pionium



Pionium $\pi^+\pi^-$ atoms decaying to $\pi^0\pi^0$?

- Observe an excess of events at cusp point $m_{\pi^0\pi^0} = 2m_{\pi^+}$ of $(1.8 \pm 0.2) \times 10^{-5}$ per $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ decay.
- Prediction for pionium formation in $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ (Silagadze, '94): 0.8×10^{-5}
- But: have to take **electromagn. corrections to FSI** into account!
 $\Rightarrow \pi^+\pi^-$ unbound states with resonant structure.
(Gevorkian, Tarasov, Voskresenskaya, PLB 649 (2007) 159)

Fit Results on $a_0 - a_2$ and a_2

fit	χ^2/ndf	$a_0 - a_2$	a_2	f_{atom}
CI	206.3/195	0.2727(46)	-0.0392(80)	0.0533(91)
CI (a)	201.6/189	0.2689(50)	-0.0344(86)	0.0533
CI (c)	210.6/196	0.2749(21)	-0.0413	0.0441(76)
CI (a,c)	207.6/190	0.2741(21)	-0.0415	0.0441
BB	462.9/452	0.2815(43)	-0.0693(136)	0.0530(95)
BB (a)	458.5/446	0.2775(48)	-0.0593(142)	0.0542
BB (c)	467.3/453	0.2737(26)	-0.0417	0.0647(76)
BB (a,c)	459.8/447	0.2722(27)	-0.0421	0.0647

CI	205.6/195	0.2483(45)	-0.0092(91)	0.0625(92)
CI (a)	202.9/189	0.2461(49)	-0.0061(98)	0.0625
CI (c)	222.1/196	0.2646(21)	-0.0443	0.0420(77)
CI (a,c)	219.7/190	0.2645(22)	-0.0444	0.0420
BB	477.4/452	0.2571(48)	-0.0241(129)	0.0631(97)
BB (a)	474.4/446	0.2544(51)	-0.0194(132)	0.0631
BB (c)	479.8/453	0.2633(24)	-0.0447	0.0538(77)
BB (ac)	478.1/447	0.2627(25)	-0.0449	0.0538

Rad. corr. off

a: pionium f_{atom} fixed

c: with ChPT constraint

a,c: both

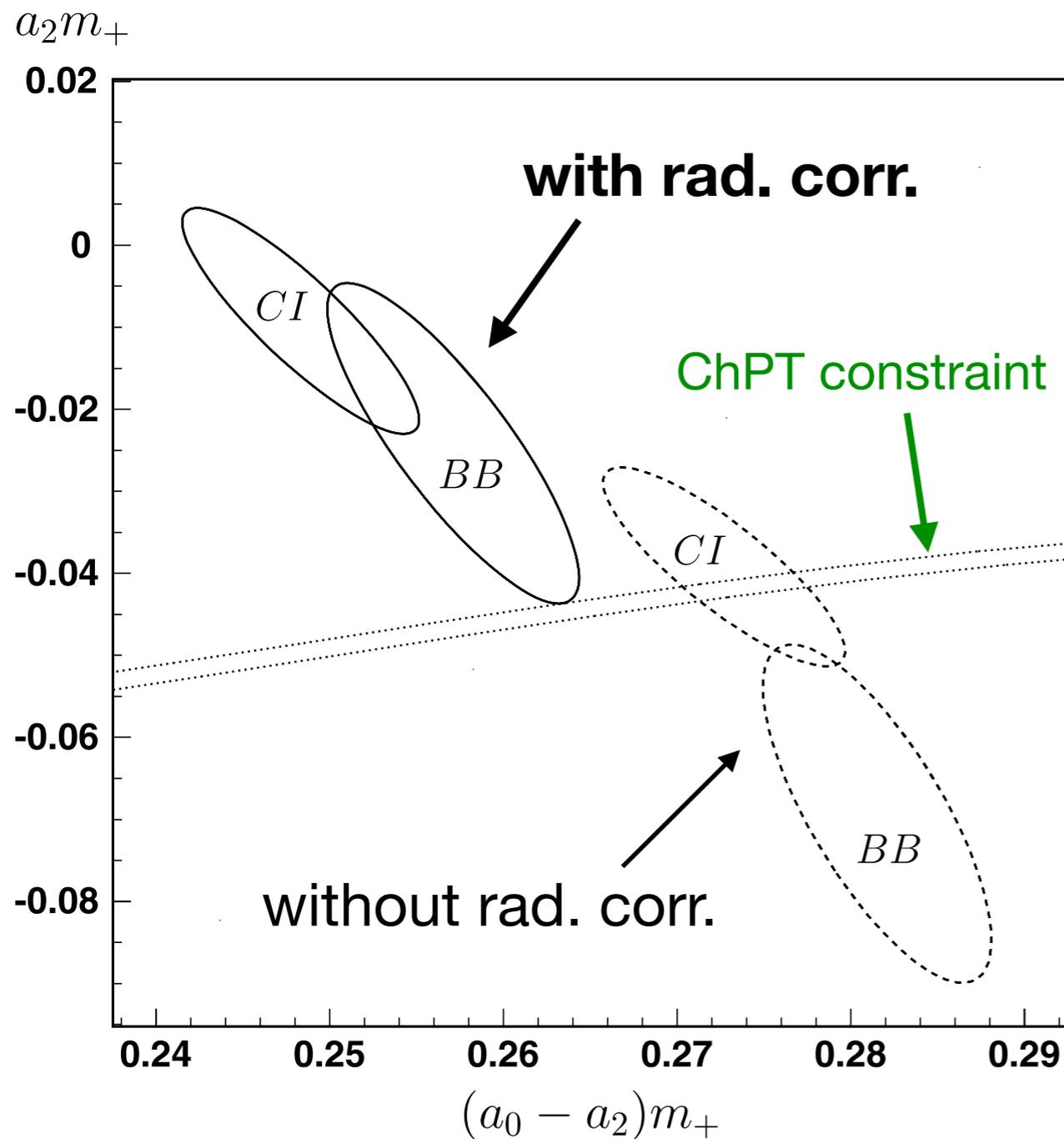
(...): statistical error

Rad. corr. on

← **final result**

(f_{atom} , $a_0 - a_2$, a_2
free in the fit)

Cusp Results on $a_0 - a_2$ and a_2



(only statistical uncertainties shown)

Final result: (EPJC 64 (2009) 589)

$$a_0 - a_2 = 0.257(5)_{\text{stat}}(3)_{\text{sys}}(1)_{\text{ext}}$$

$$a_2 = -0.024(13)_{\text{stat}}(9)_{\text{sys}}(2)_{\text{ext}}$$

(statistical correlation -0.839)

With ChPT constraint:

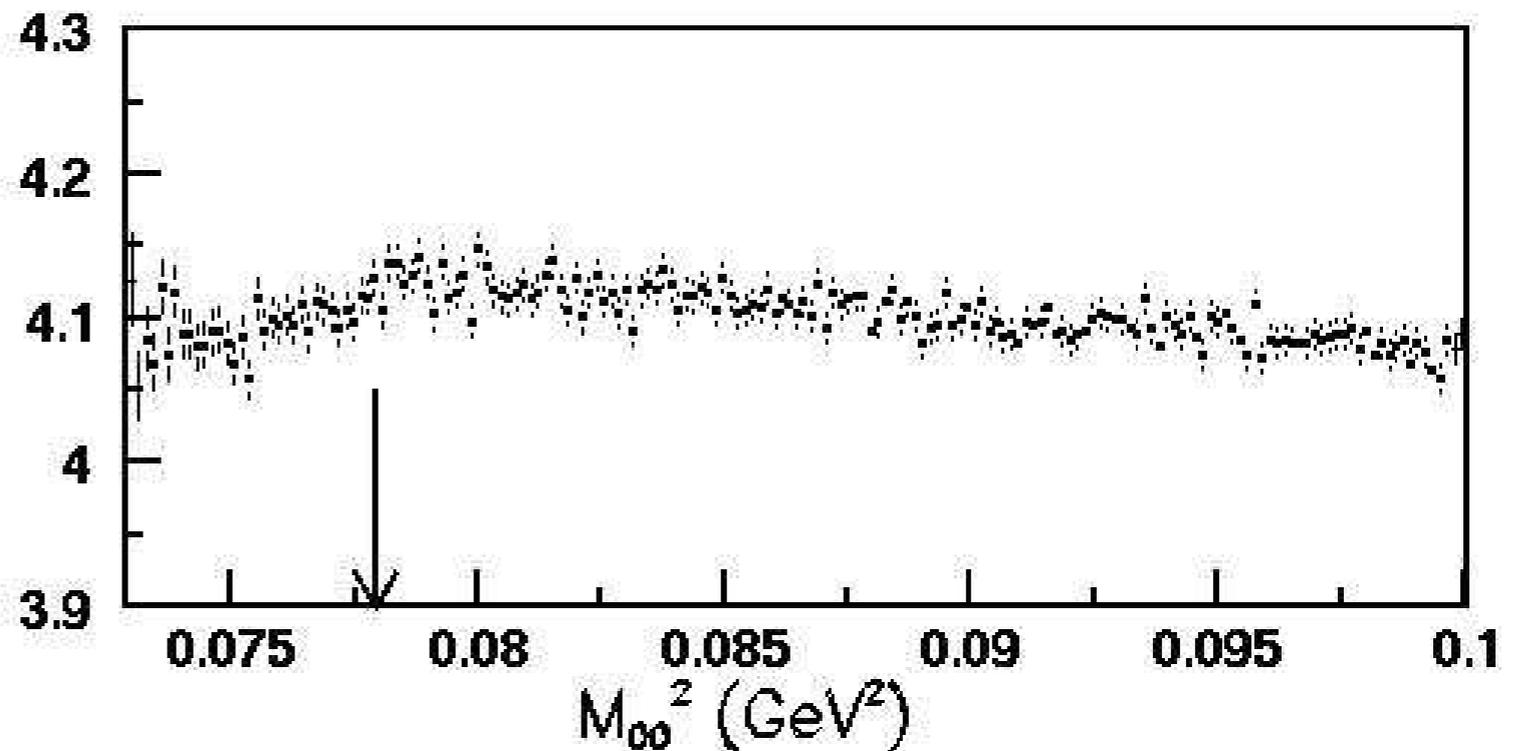
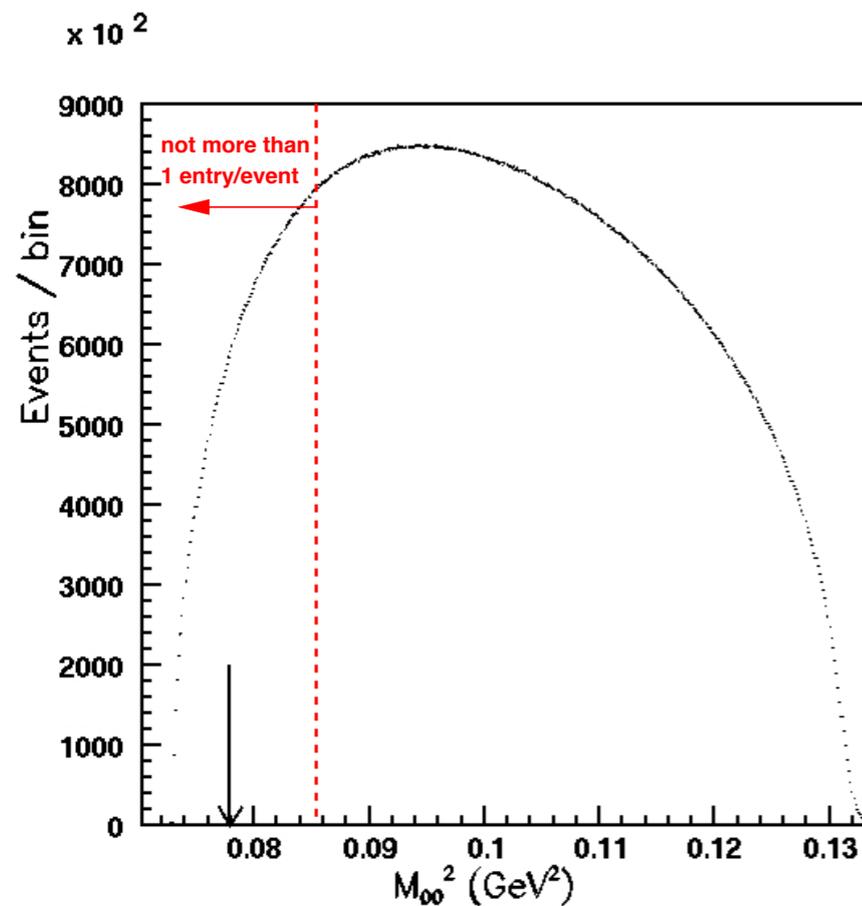
$$a_0 - a_2 = 0.2633(24)_{\text{stat}}(14)_{\text{sys}}(19)_{\text{ext}}$$

ChPT prediction:

$$a_0 - a_2 = 0.265(4)$$

Rescattering in $K_L \rightarrow \pi^0\pi^0\pi^0$ (NA48/2)

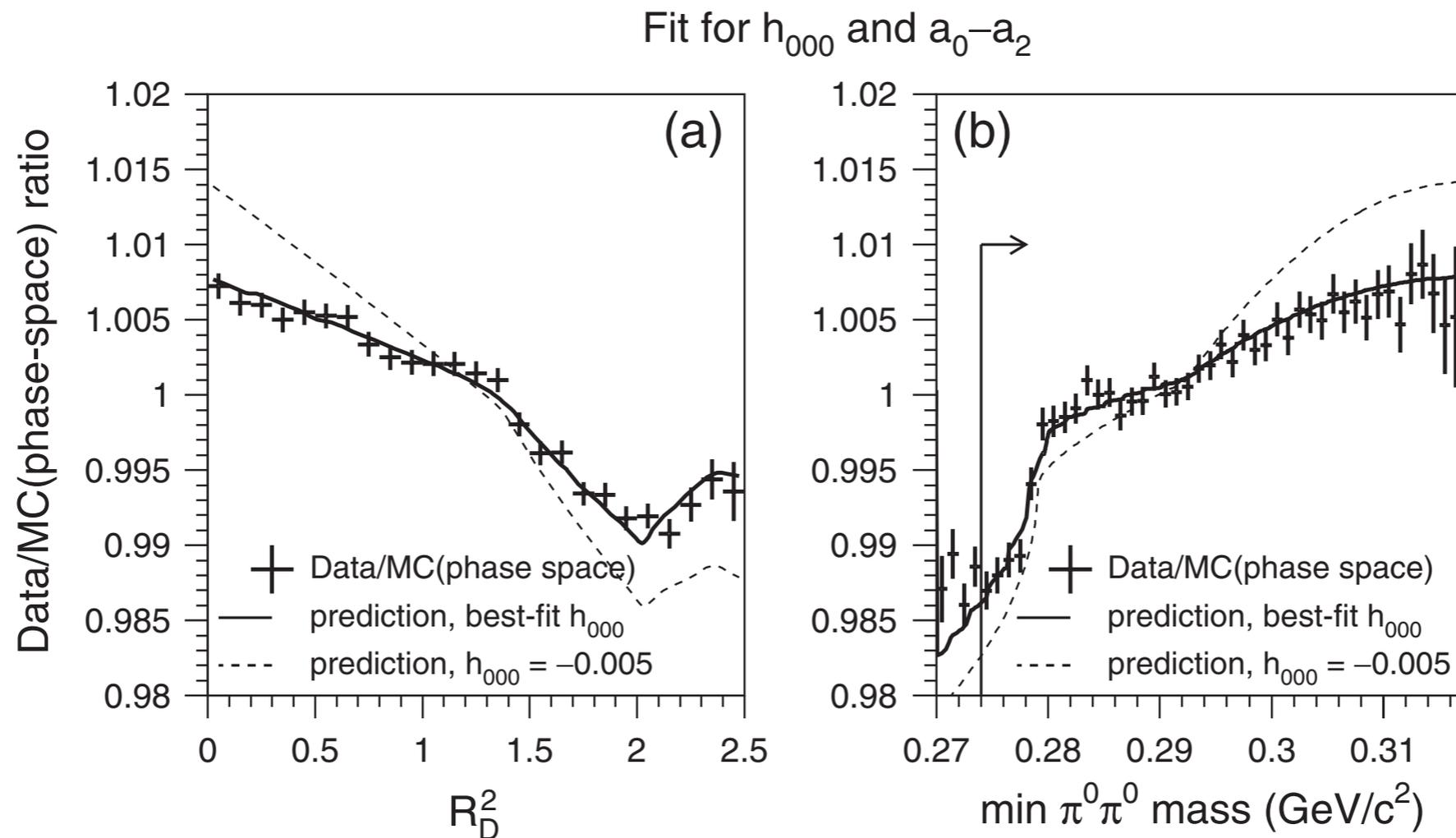
NA48/2 data taking in 2000:



Evidence for a change in slope near the cusp point.

Rescattering in $K_L \rightarrow \pi^0\pi^0\pi^0$ (KTeV)

(PRD 78 (2008) 032009)



$$m_{\pi^+}(a_0 - a_2) = 0.215 \pm 0.014_{\text{stat}} \pm 0.025_{\text{syst}} \pm 0.006_{\text{ext}}$$

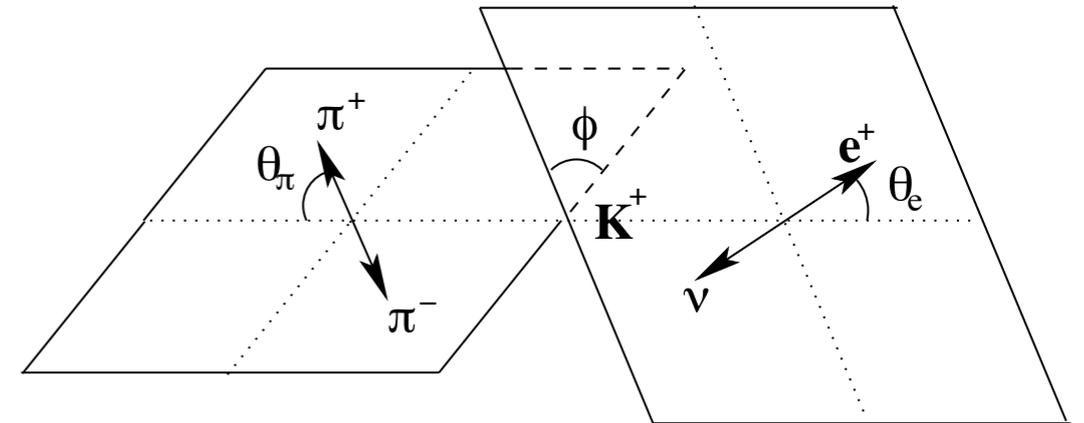
$$h_{000} = (-2.09 \pm 0.62_{\text{stat}} \pm 0.72_{\text{syst}} \pm 0.28_{\text{ext}}) \times 10^{-3}$$

Scattering Lengths from $K^\pm \rightarrow \pi^+ \pi^- e^\pm \nu$ Decays

$\pi\pi$ Scattering Lengths from K_{e4} Decays

K_{e4} decay:

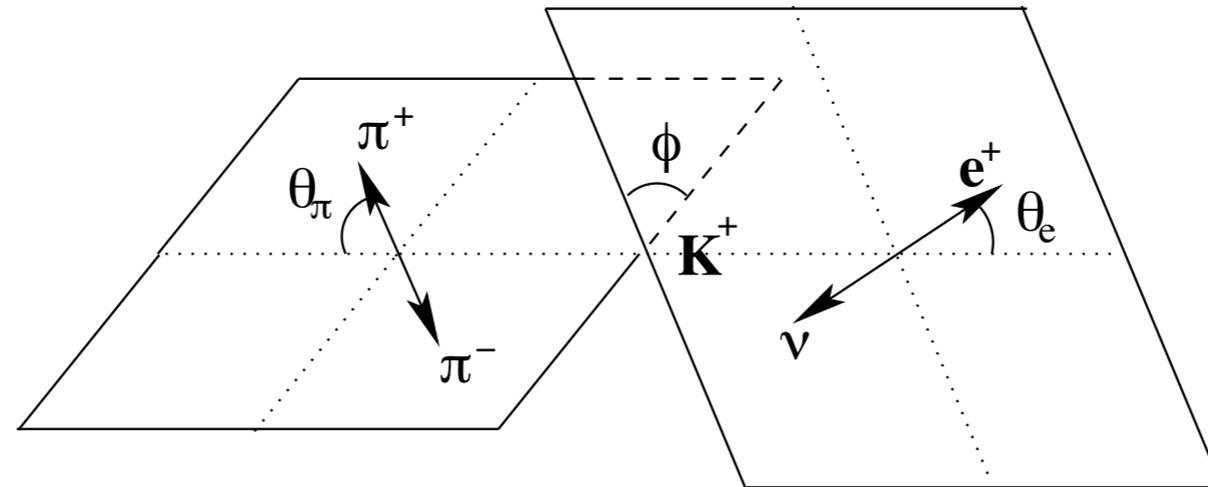
- $K_{e4} = K^\pm \rightarrow \pi^+ \pi^- e^\pm \nu$
- Very rare: $\text{Br}(K_{e4}) \sim 4 \times 10^{-5}$



Measurement of $\pi\pi$ scattering in K_{e4} :

- K_{e4} decay amplitude depends on two complex phases:
 - $\delta_0 = \pi\pi$ scattering phase shift for $I = 0, l = 0$ (*S-wave*)
 - $\delta_1 = \pi\pi$ scattering phase shift for $I = 1, l = 1$ (*P-wave*)($I = 2$ suppressed by $\Delta I = \frac{1}{2}$ rule)
- Decay rate depends on **difference** $\delta = \delta_0 - \delta_1$, with $\delta = \delta(m_{\pi\pi})$.
- $\delta \neq 0$ implies **asymmetric distribution** of lepton w.r.t. $\pi\pi$ plane.

Kinematic Variables in K_{e4}



K_{e4} is 4-body decay \implies 5 independent kinematic variables.
(Cabibbo-Maksymowicz variables)

$s_\pi = M_{\pi\pi}^2$ Invariant **di-pion mass** squared.

$s_e = M_{e\nu}^2$ Invariant **di-lepton mass** squared.

θ_π **Angle of π^+** w.r.t. $\pi\pi$ direction of flight in $\pi\pi$ rest frame.

θ_e **Angle of e^+** w.r.t. $e\nu$ direction of flight in $e\nu$ rest frame.

ϕ **Angle of $\pi\pi$ decay plane** w.r.t. $e\nu$ decay plane.

K_{e4} Selection

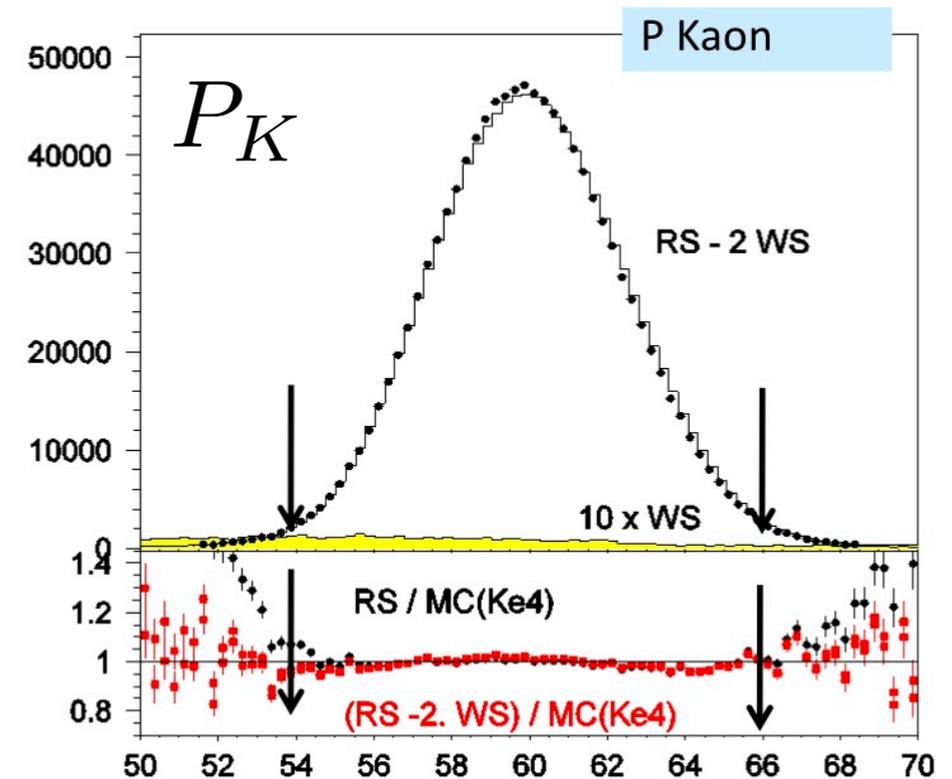
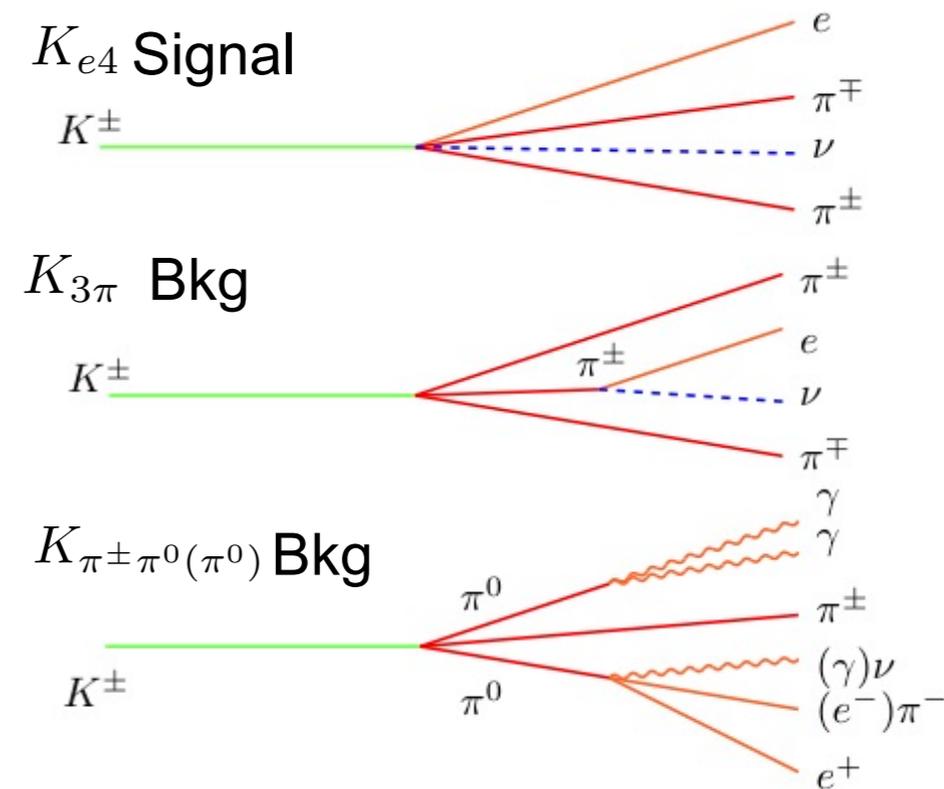
K_{e4} selection:

- 3 charged tracks and 1 good vertex
- 2 opposite-sign pions, 1 electron ($E/p \sim 1$)
- missing transverse momentum
- kaon momentum close to 60 GeV/c

Background:

- $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ with $\pi \rightarrow e\nu$ or mis-identified pion
- $K^\pm \rightarrow \pi^\pm \pi^0 (\pi^0)$ with $\pi^0 \rightarrow e^+ e^- \nu$ and mis-identified electron
- **Background estimation** from wrong-sign $\pi^+ \pi^+ e^-$ events

→ **Background ~ 0.6 %**



K_{e4} Fitting Procedure

- Full event sample (2003+2004): **1.13 million K_{e4} decays**

- Fit in *iso-populated boxes* in the 5-dim. CM variables:

$$10(M_{\pi\pi}) \times 5(M_{e\nu}) \times 5(\cos\theta_e) \times 5(\cos\theta_\pi) \times 12(\phi) = 15000 \text{ Boxes}$$

- Assuming constant form factors, K^+ and K^- samples fitted separately in **10 independent $M_{\pi\pi}$ bins** and then combined in each $M_{\pi\pi}$ bin.

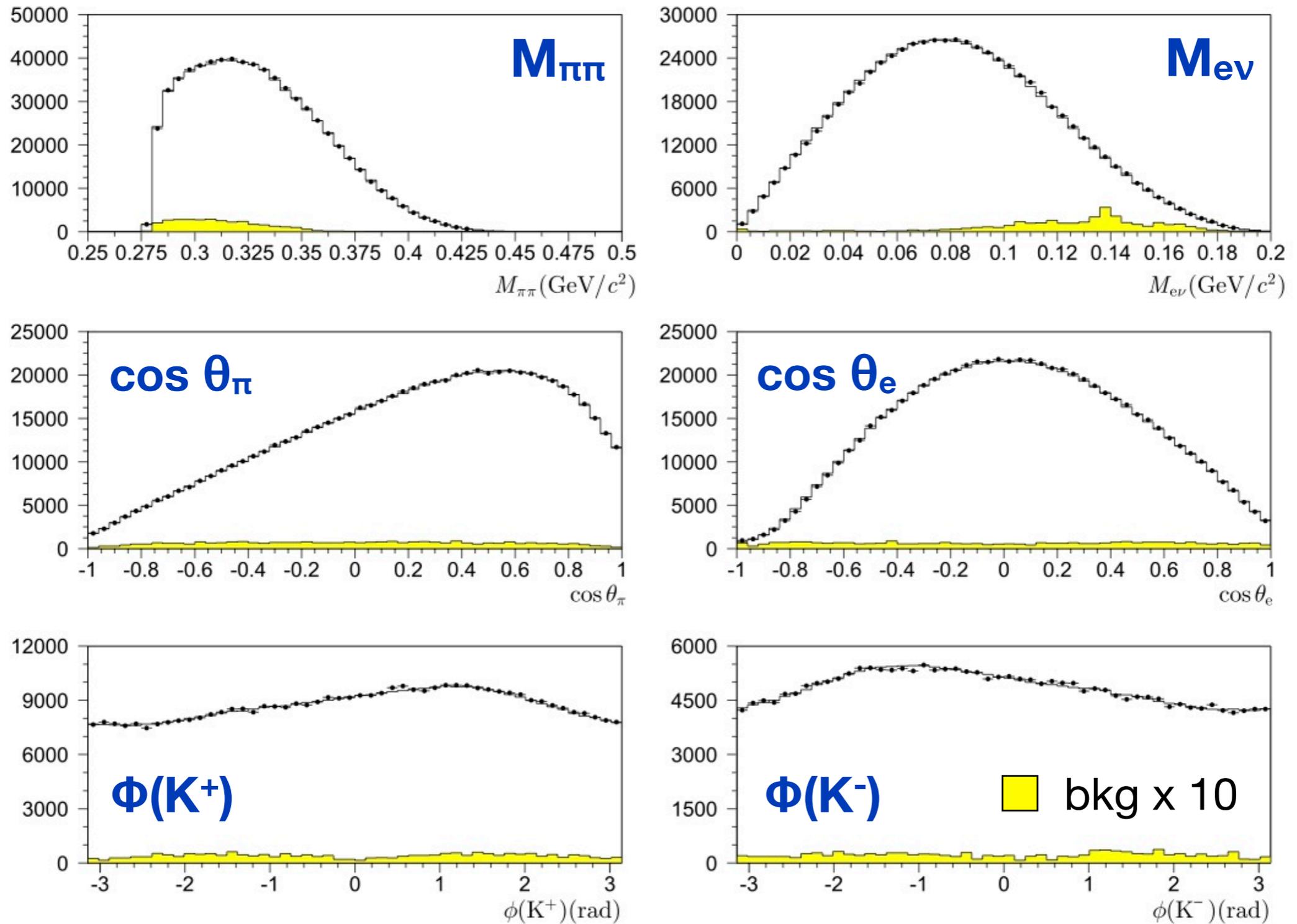
Data:

K^+ sample:	726 400 events	→	48 events/box
K^- sample:	404 400 events	→	27 events/box

MC:

K^+ sample:	17.4×10^6 events	→	1160 events/box
K^- sample:	9.7×10^6 events	→	650 events/box

K_{e4} Fit Results



K_{e4} Form Factor Results

Partial wave expansion of form factors:

$$F = F_s e^{i\delta_s} + F_p e^{i\delta_p} \cos \theta_\pi + d \text{ wave} \dots$$

$$G = G_p e^{i\delta_g} + d \text{ wave} \dots$$

$$H = H_p e^{i\delta_h} + d \text{ wave} \dots$$

Single form factors parametrized in Taylor expansion:

$$F_s = f_s + f'_s q^2 + f''_s q^4 + f'_e S_e / 4m_\pi^2 + \dots$$

$$F_p = f_p + f'_p q^2 + \dots$$

$$G_p = g_p + g'_p q^2 + \dots$$

$$H_p = h_p + h'_p q^2 + \dots$$

$$\delta(q^2) = \delta_s - \delta_p.$$

$$q^2 = (S_\pi / 4m_\pi^2) - 1$$

- All form factors measured w.r.t. f_s
- Systematics from acceptance and background control

$$f'_s / f_s = 0.152 \pm 0.007_{\text{stat}} \pm 0.005_{\text{syst}}$$

$$f''_s / f_s = -0.073 \pm 0.007_{\text{stat}} \pm 0.006_{\text{syst}}$$

$$f'_e / f_s = 0.068 \pm 0.006_{\text{stat}} \pm 0.007_{\text{syst}}$$

$$f_p / f_s = -0.048 \pm 0.003_{\text{stat}} \pm 0.004_{\text{syst}}$$

$$g_p / f_s = 0.868 \pm 0.010_{\text{stat}} \pm 0.010_{\text{syst}}$$

$$g'_p / f_s = 0.089 \pm 0.017_{\text{stat}} \pm 0.013_{\text{syst}}$$

$$h_p / f_s = -0.398 \pm 0.015_{\text{stat}} \pm 0.008_{\text{syst}}$$

→ **First evidence of non-zero f'_e and f_p !**

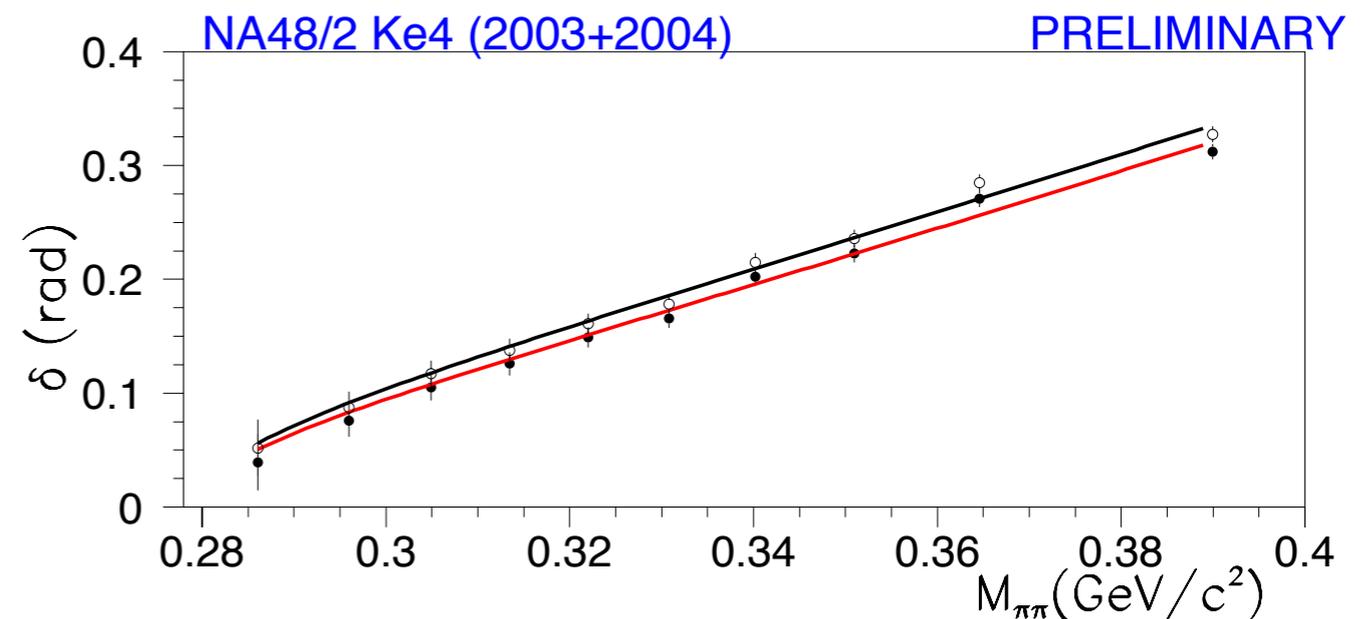
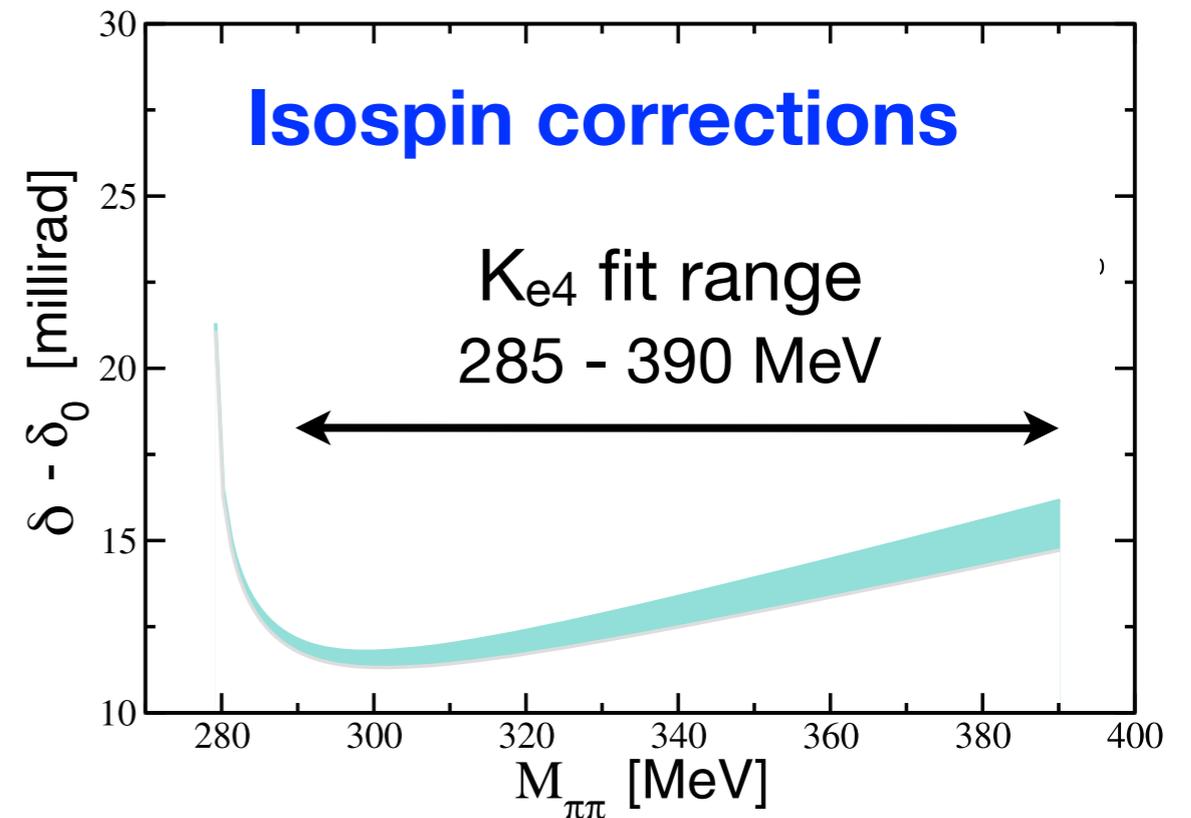
From phase shifts to scattering lengths

Corrections to be applied:

- **Radiative effects:**
Included in the simulation (Coulomb attraction, IB).
- **Mass effects:**
Isospin corrections have to be applied to δ .
Developed in close collaboration with NA48/2.
(Colangelo, Gasser, Rusetsky, EPJC 59 (2009) 777)



Effect of 10-15 mrad on δ
(stat. precision $\sim 7-8$ mrad)



NA48/2 Fit of Phase Shift $\delta = \delta_0 - \delta_1$

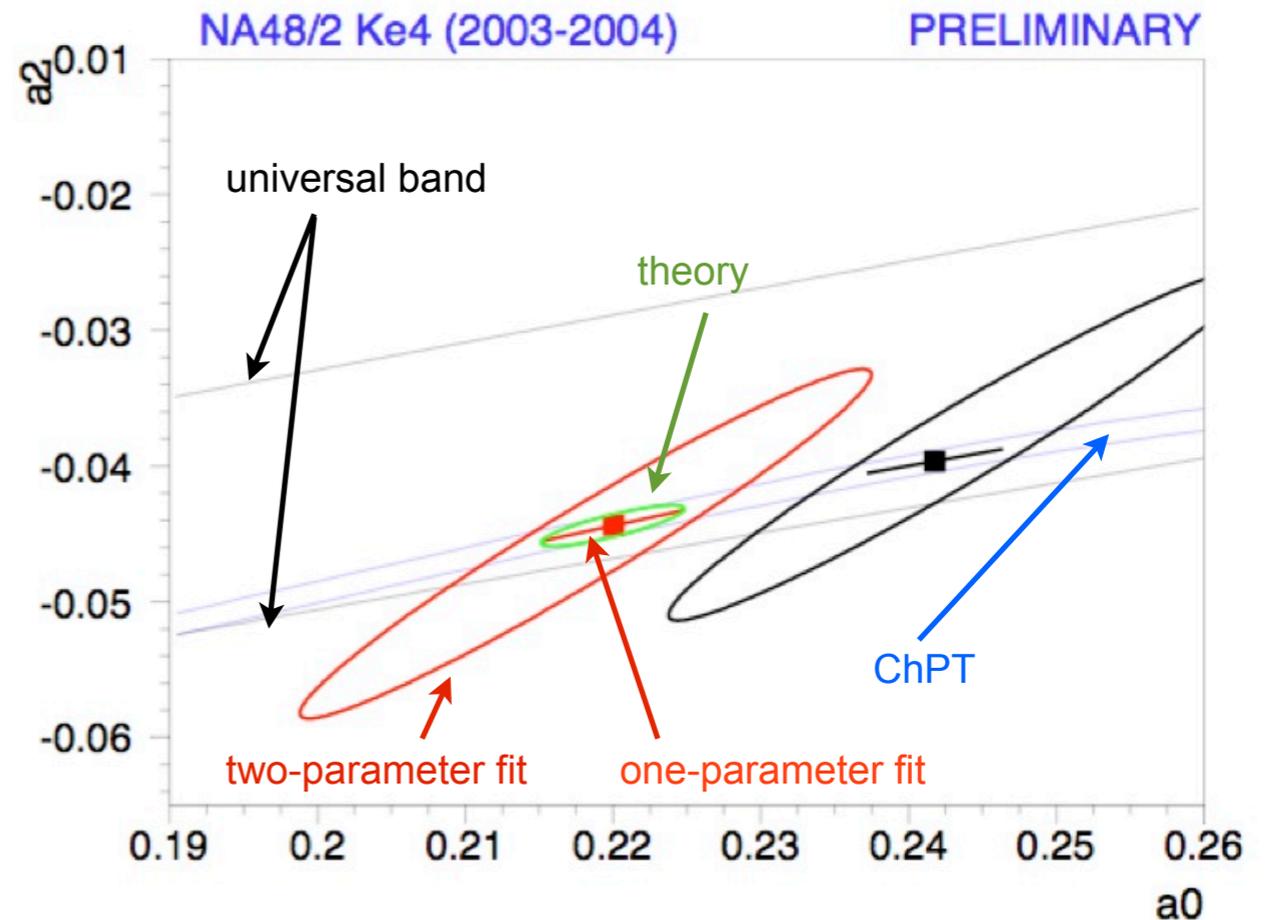
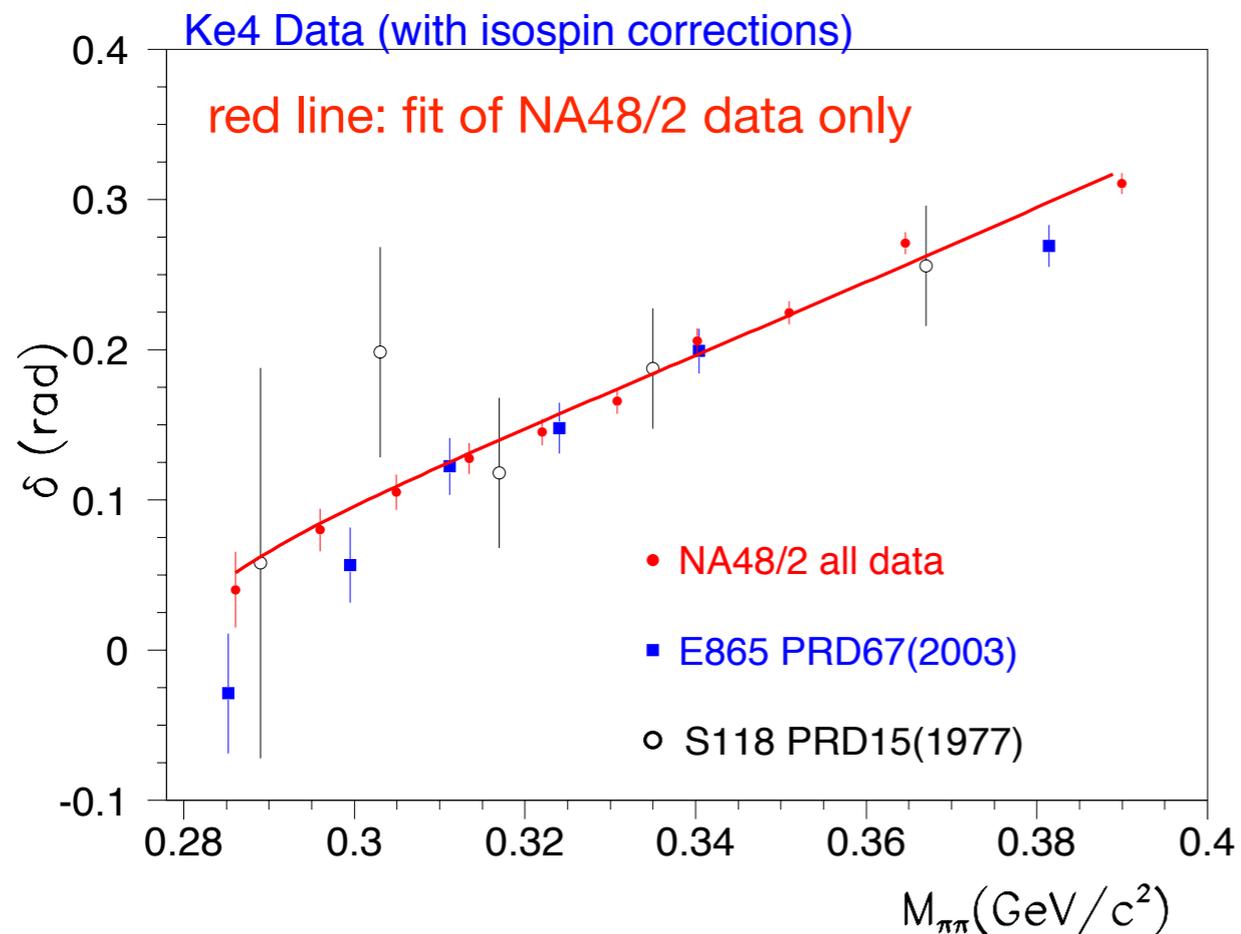
Two-parameter fit:

$$a_0 = 0.2220 \pm 0.0128_{\text{stat}} \pm 0.0050_{\text{syst}} \pm 0.0037_{\text{theo}}$$

$$a_2 = -0.0432 \pm 0.0086_{\text{stat}} \pm 0.0034_{\text{syst}} \pm 0.0028_{\text{theo}}$$

One-parameter fit: (with ChPT constraint)

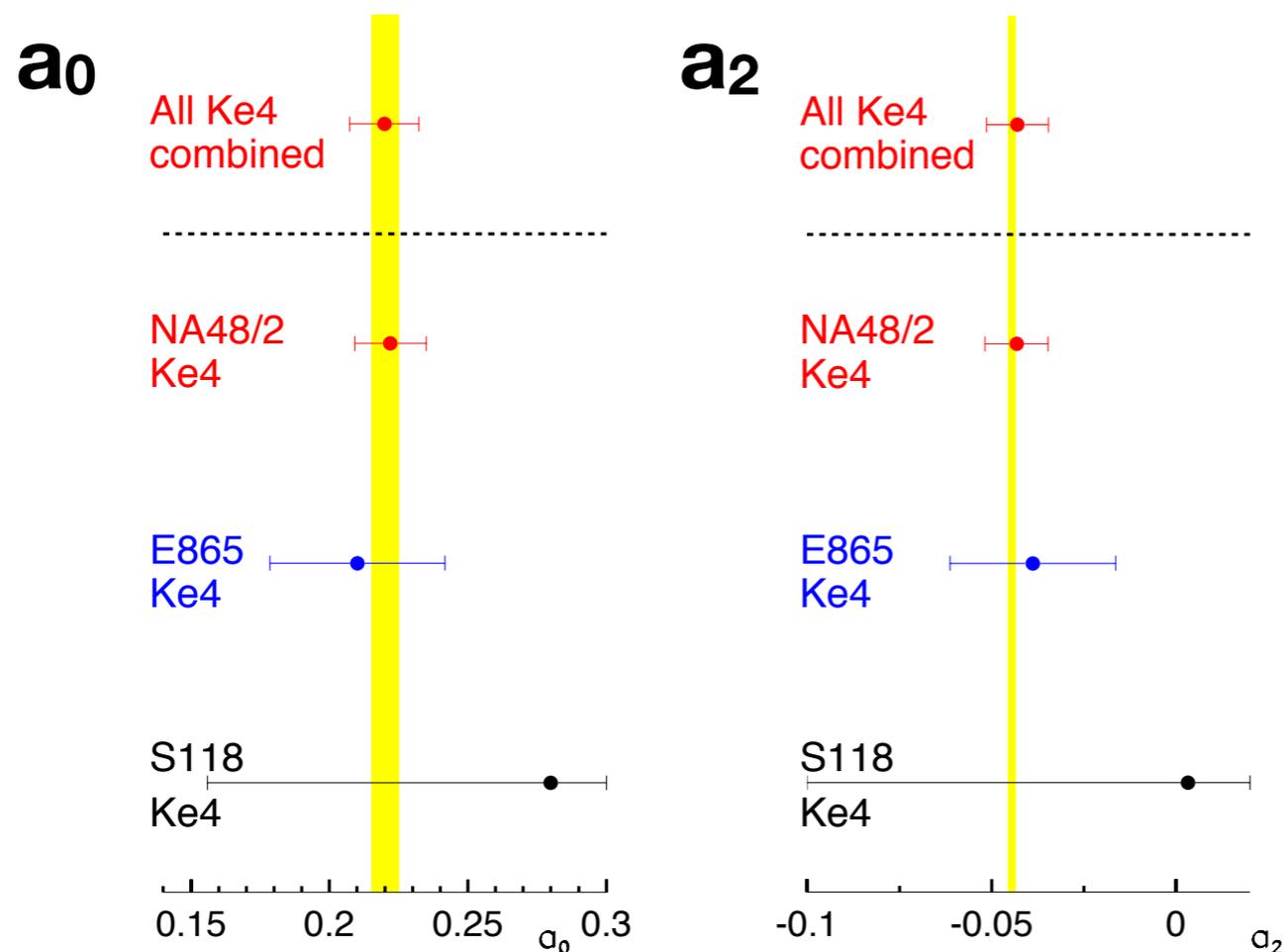
$$a_0 = 0.2206 \pm 0.0049_{\text{stat}} \pm 0.0018_{\text{syst}} \pm 0.0064_{\text{theo}}$$



Comparison with previous K_{e4} Results

Comparison of K_{e4} measurements

(without ChPT constraint, old experiments isospin corrected):



■ NA48/2 dominates K_{e4} measurements

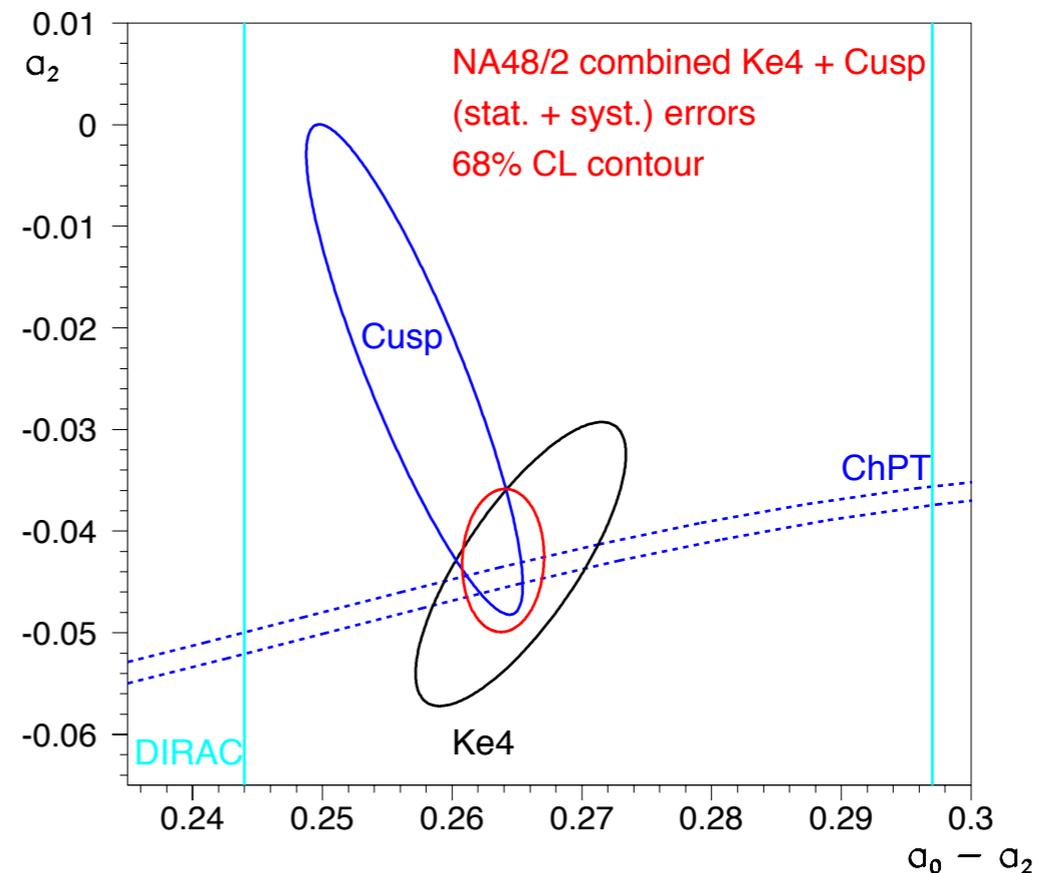
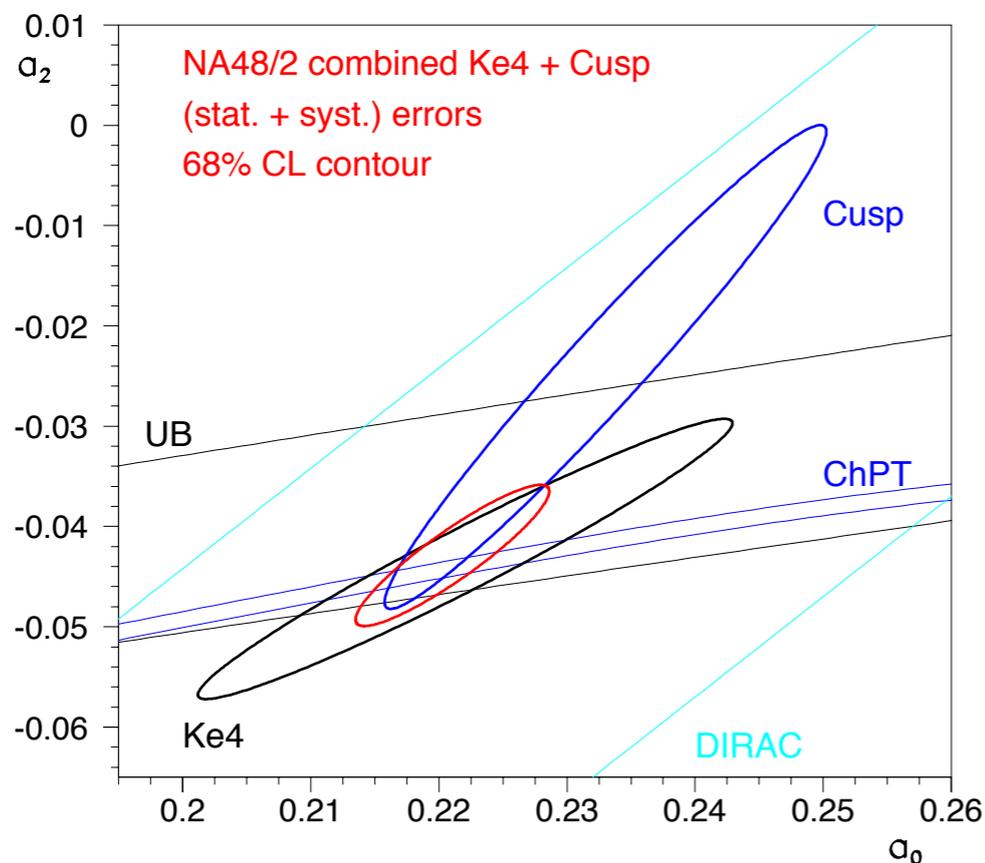
■ **Perfect agreement with ChPT prediction!**

Yellow lines: Theory prediction
(not combination of measurements)

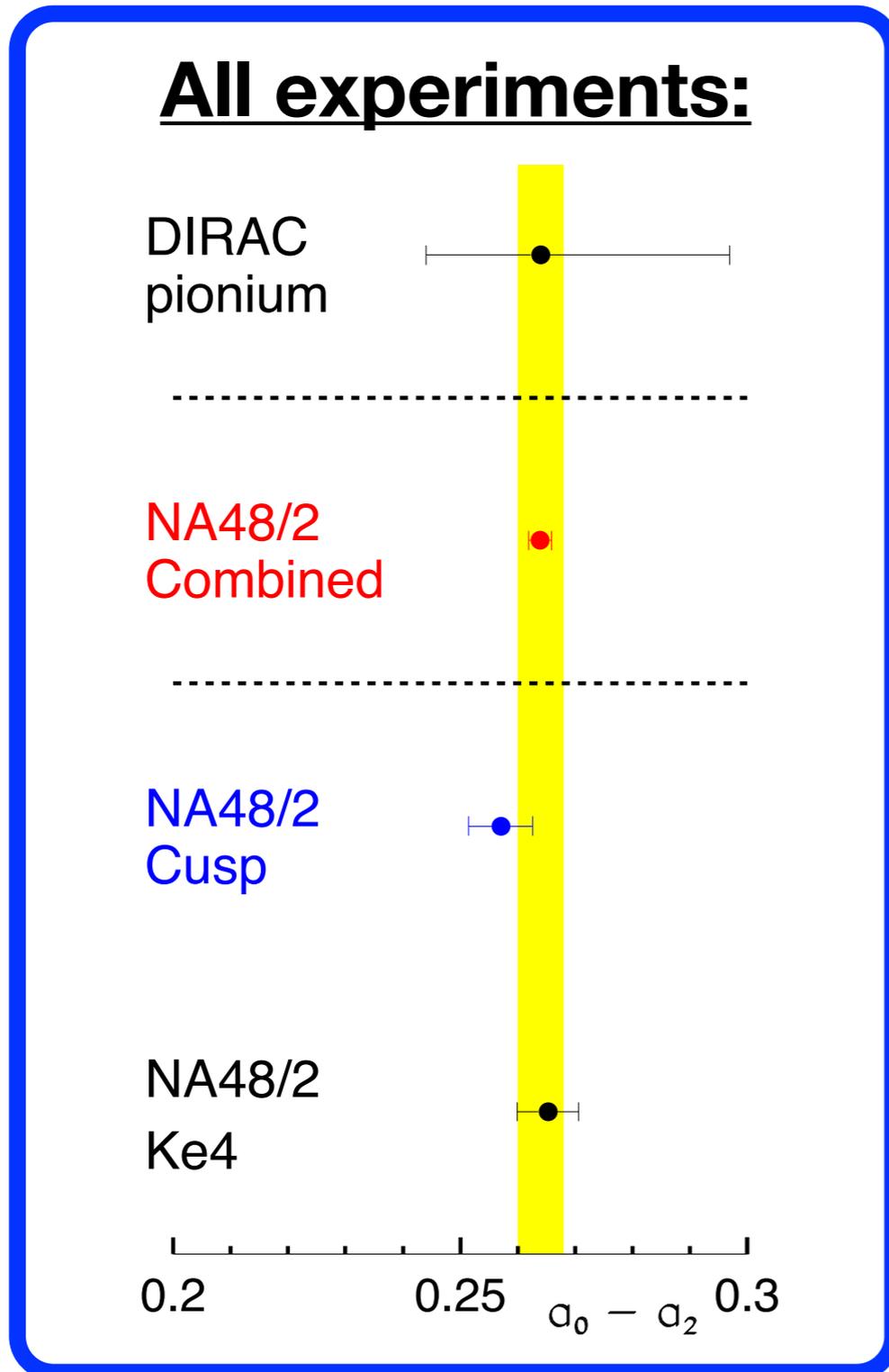
Combination of Cusp and Ke4

Two independent measurements with **different samples**,
different systematics and **different theory**:

$$\begin{aligned}
 a_0 &= 0.2210 \pm 0.0047_{\text{stat}} \pm 0.0040_{\text{syst}} \\
 a_2 &= -0.0429 \pm 0.0044_{\text{stat}} \pm 0.0028_{\text{syst}} \\
 a_0 - a_2 &= 0.2639 \pm 0.0020_{\text{stat}} \pm 0.0015_{\text{syst}}
 \end{aligned}$$



Combination with other Measurements



**Experimental data
has reached
theoretical precision**

**Perfect agreement
with ChPT prediction**

Yellow line: Theory prediction
(*not* combination of measurements)

(no uncertainty from theory on cusp result)

Future Prospects

Future for Scattering Lengths from Kaons

NA48/2: Results still not systematically limited, but need more kaons. Single event sensitivity $\sim 10^{-10}$ for NA48/2.

→ Need more than ~ 100 billion kaon decays to top this!
(if acceptance is about 10% and trigger efficiency ~ 1)

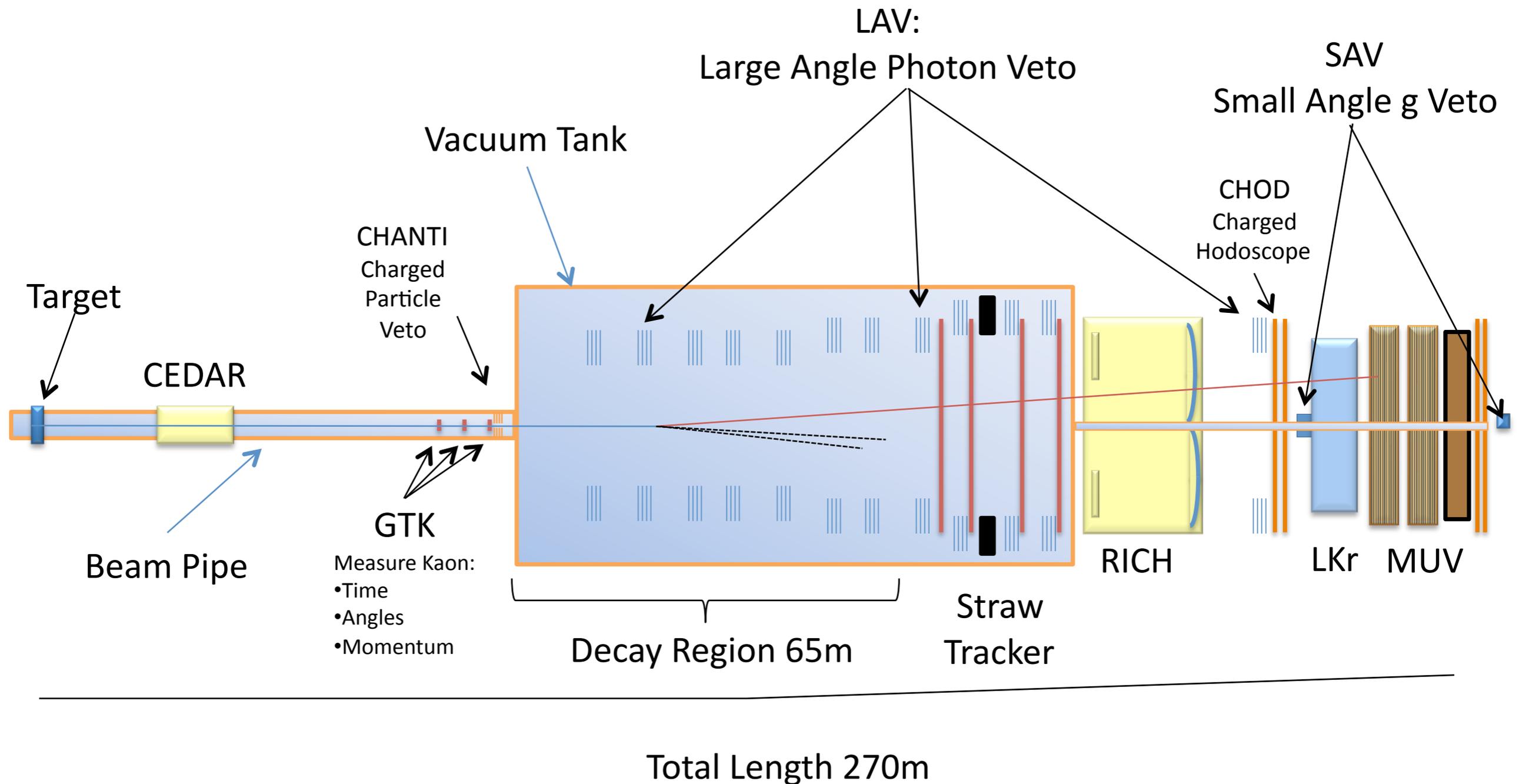
NA62: New experiment for measuring 100 $K^\pm \rightarrow \pi^\pm \nu \nu$ events, starting in 2013 ($\text{Br} \sim 10^{-10}$).

→ Single event sensitivity $\sim 10^{-12}$!

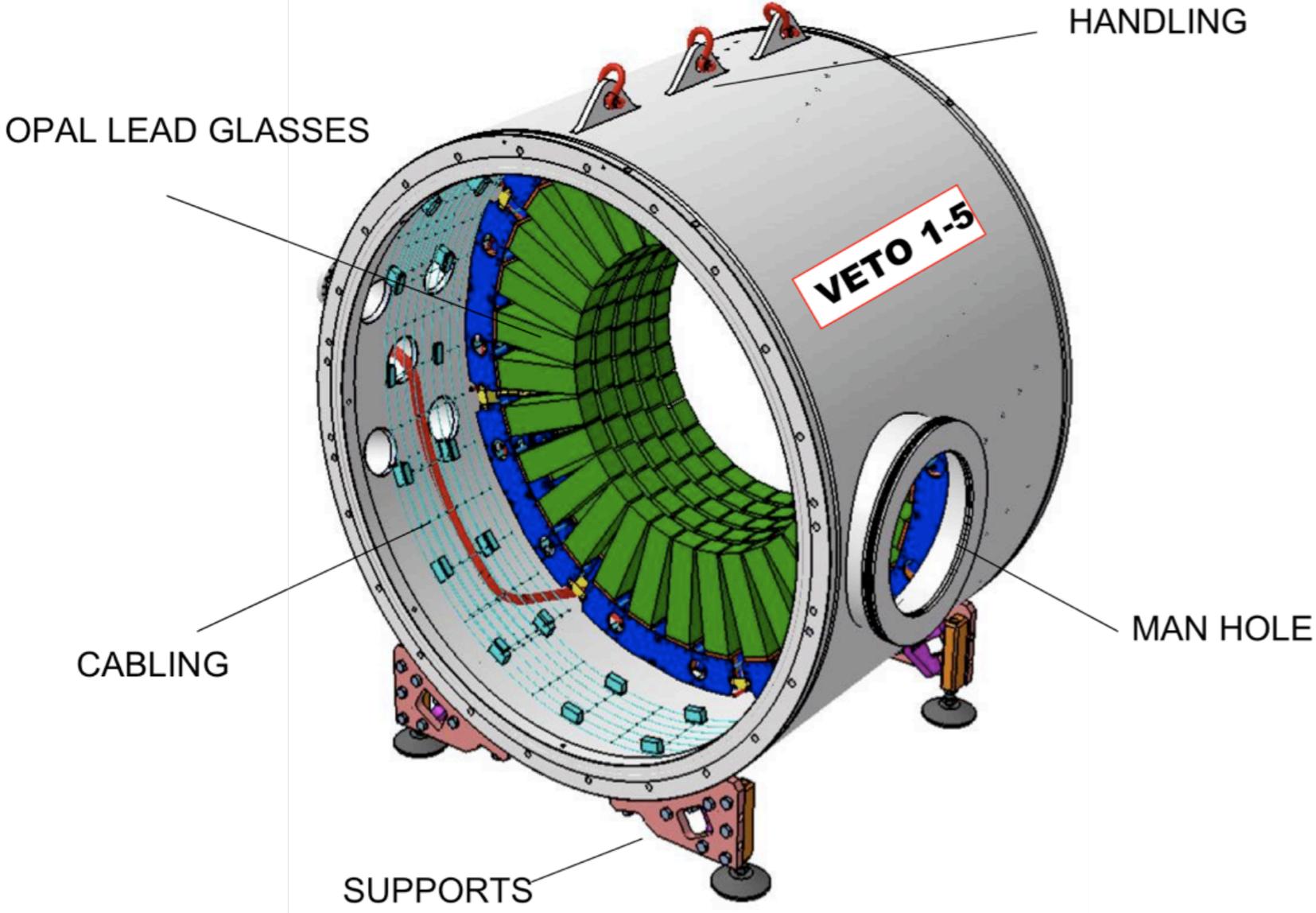
However: trigger conditions are not selecting $K_{3\pi}$, K_{e4}

→ Special running periods needed.

The NA62 Detector



First Detectors



Conclusions

- **Kaon decays** give unique possibility to study low-energy hadronic interactions with high precision.
- With very high statistics, **NA48/2** has checked **ChPT predictions** on **$\pi\pi$ scattering lengths** with high accuracy using both **$K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$** and **$K_{e4}$** decays.
→ **very strong test of the theory**
- **NA62**: Measurement of **100 $K^\pm \rightarrow \pi^\pm \nu \nu$ events**, starting in **2013**, also possibility to collect even more statistics for scattering length measurements.

Spares

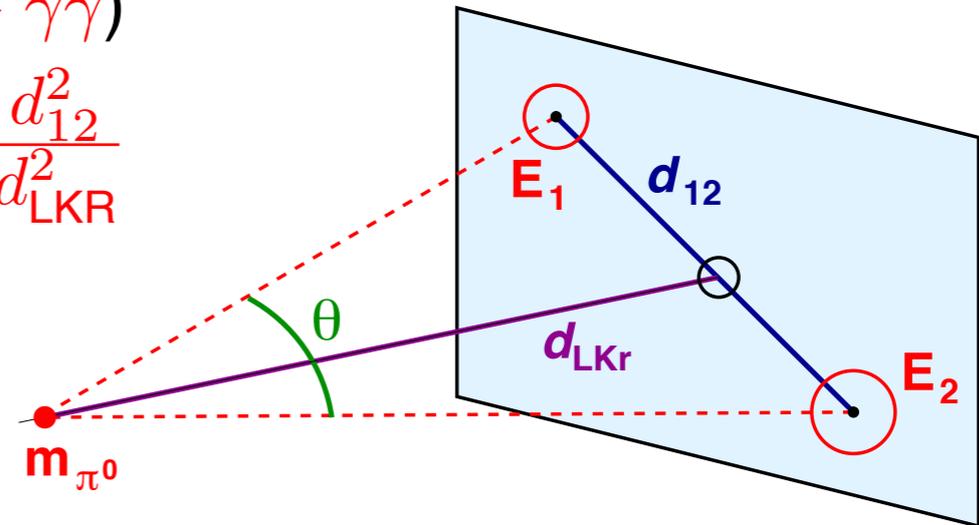
$K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$: $M(\pi^0 \pi^0)$ Reconstruction

Decays in two photons: (z.B. $\pi^0 \rightarrow \gamma\gamma$)

$$m_{\pi^0}^2 = 2 E_1 E_2 (1 - \cos \theta) = E_1 E_2 \frac{d_{12}^2}{d_{LKR}^2}$$

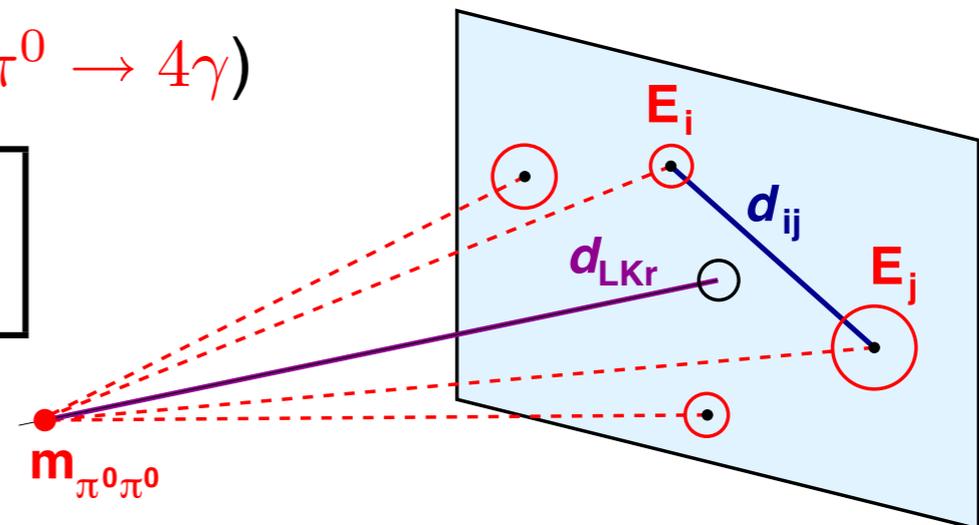
If mass m_{π^0} known:

$$d_{LKR} = \frac{1}{m_{\pi^0}} \sqrt{E_1 E_2 d_{12}^2}$$



Decays in many photons: (z.B. $\pi^0 \pi^0 \rightarrow 4\gamma$)

$$d_{LKR} = \frac{1}{m_{\pi^0 \pi^0}} \sqrt{\sum_{i,j;i>j} E_i E_j d_{ij}^2}$$



Turning it around:

Vertex d_{LKR} always the same!

⇒ Invariant mass $m_{\pi^0 \pi^0}$ only from LKr information:

$$m_{\pi^0 \pi^0} = m_{\pi^0} \sqrt{\sum_{i,j;i>j} E_i E_j d_{ij}^2 / \frac{1}{2} \left(\sqrt{E_1 E_2 d_{12}^2} + \sqrt{E_3 E_4 d_{34}^2} \right)}$$