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Nuclear Physics A 749 (2005) 243c–250c

NUCLEAR  
PHYSICS **A**

## Transverse momentum dependence of charmonium suppression in Pb-Pb collisions at the CERN SPS

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Charmonium suppression in Pb-Pb collisions at 158 GeV/c per nucleon is investigated in detail with the study of the transverse momentum distributions of  $J/\psi$  as a function of the centrality of the collision. It is shown that the observed  $J/\psi$  suppression in Pb-Pb interactions is particularly significant mainly at low transverse momentum where it strongly depends on centrality. For peripheral Pb-Pb collisions, the transverse momentum dependence of the  $J/\psi$  cross section is, as a function of centrality, qualitatively similar to the dependence observed in p-A and S-U collisions. Comparing peripheral and central

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Pb-Pb collisions, the data show a relative suppression in the whole  $p_T$  range although its amplitude significantly decreases with increasing  $p_T$  and becomes almost  $p_T$  independent for the highest  $p_T$  values.

## 1. INTRODUCTION

Charmonium suppression in ultrarelativistic heavy ion collisions is considered as a potential signature of the phase transition from normal nuclear matter to a deconfined state of quarks and gluons. Charmonium production has been measured by the NA50 Collaboration in Pb-Pb collisions at 158 GeV/c per nucleon and in proton-nucleus collisions at 400 and 450 GeV/c [1,2].

Normal nuclear absorption of  $J/\psi$  has been measured in proton-induced reactions. The corresponding cross-section, deduced in the frame of a Glauber calculation, amounts to  $4.24 \pm 0.4$  mb [3]. It provides thereby the  $J/\psi$  normal nuclear absorption reference as a function of the path in nuclear matter that the produced  $c\bar{c}$  pair has to go through in order to emerge and survive, a quantity which is directly related to the centrality of the collision. The main result of the NA50 experiment in the study of Pb-Pb collisions is that whereas peripheral Pb-Pb collisions approximately follow the normal nuclear absorption pattern, a departure from this normal behaviour is observed for semi-central reactions which increases in amplitude with increasing centrality.

New results have been obtained recently for the production of the  $\psi'$ . The absorption cross section of  $\psi'$  in nuclear matter, as deduced from p-A experiments, is  $7.9 \pm 0.6$  mb. As expected from a loosely bound state,  $\psi'$  absorption increases significantly already in S-U reactions. In Pb-Pb central collisions, the  $\psi'$  suppression is about two times stronger than for  $J/\psi$ .

Preliminary results obtained from our latest data samples collected under improved experimental conditions can be found in [4,5]. In this article we extend our analysis of  $J/\psi$  production and study the suppression as a function of the transverse momentum of the charmonium state.

## 2. TRANSVERSE MOMENTUM DISTRIBUTIONS OF CHARMONIUM

Some of the features of the  $J/\psi$  transverse momentum distributions obtained from the first data samples collected by the NA50 experiment can be found in [6]. In particular, the dependence, as a function of the centrality of the collision, of the mean square transverse momentum and of the slope of the  $M_T$  spectra were obtained from these data. When rescaled to the same energy and as a function of the mean length path of  $J/\psi$  in nuclear matter, the mean square transverse momentum of  $J/\psi$  exhibits the same behaviour for p-A, S-U and Pb-Pb collisions [7], which could be related to initial parton scattering. The data also show a change of the slope of the T dependence on the energy density near the value where the  $J/\psi$  production cross section starts to deviate from the normal absorption curve [8].

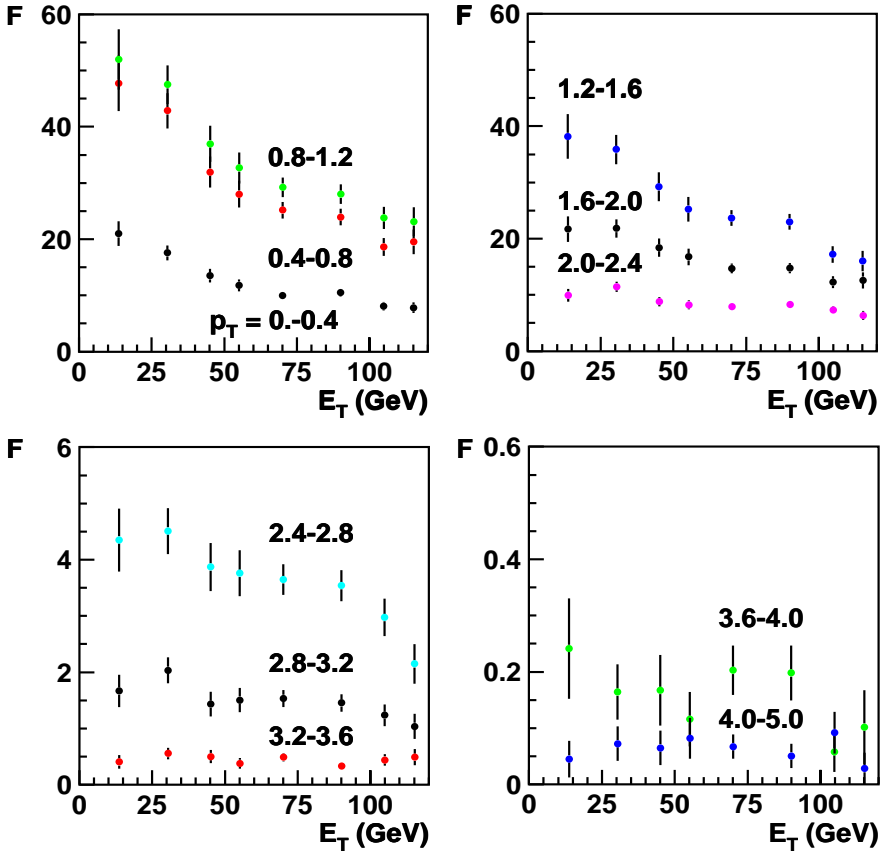


Figure 1. Ratio  $F$  of the  $J/\psi$  production cross section for Pb-Pb collisions at 158 GeV/c per nucleon in the  $p_T$  bins shown on the plots (in GeV/c) to the DY cross section, as a function of the measured neutral transverse energy in GeV.

The high quality and the size of the sample of data collected in year 2000 allows a more detailed study of the  $J/\psi$  transverse momentum. As in our previous analysis, we study the ratio of the  $J/\psi$  cross section to the Drell-Yan cross section (we consider here the Drell-Yan with invariant mass higher than  $4.2 \text{ GeV}/c^2$ ), which is proportional to the  $J/\psi$  yield per nucleon-nucleon collision. Events are binned according to the centrality of the collision in which they are produced, in fact, to the neutral transverse energy  $E_T$  which is experimentally measured, on an event by event basis, by an electromagnetic calorimeter with laboratory pseudorapidity coverage in the range [1.1-2.3].

We plot on Fig.1 the ratio  $F$  of the  $J/\psi$  to the DY cross section in the corresponding  $E_T$  bin as a function of the transverse energy  $E_T$  for 11 transverse momentum bins up to  $p_T = 5.0 \text{ GeV}/c$ . The figure shows that, whereas for low values of  $p_T$  there is a

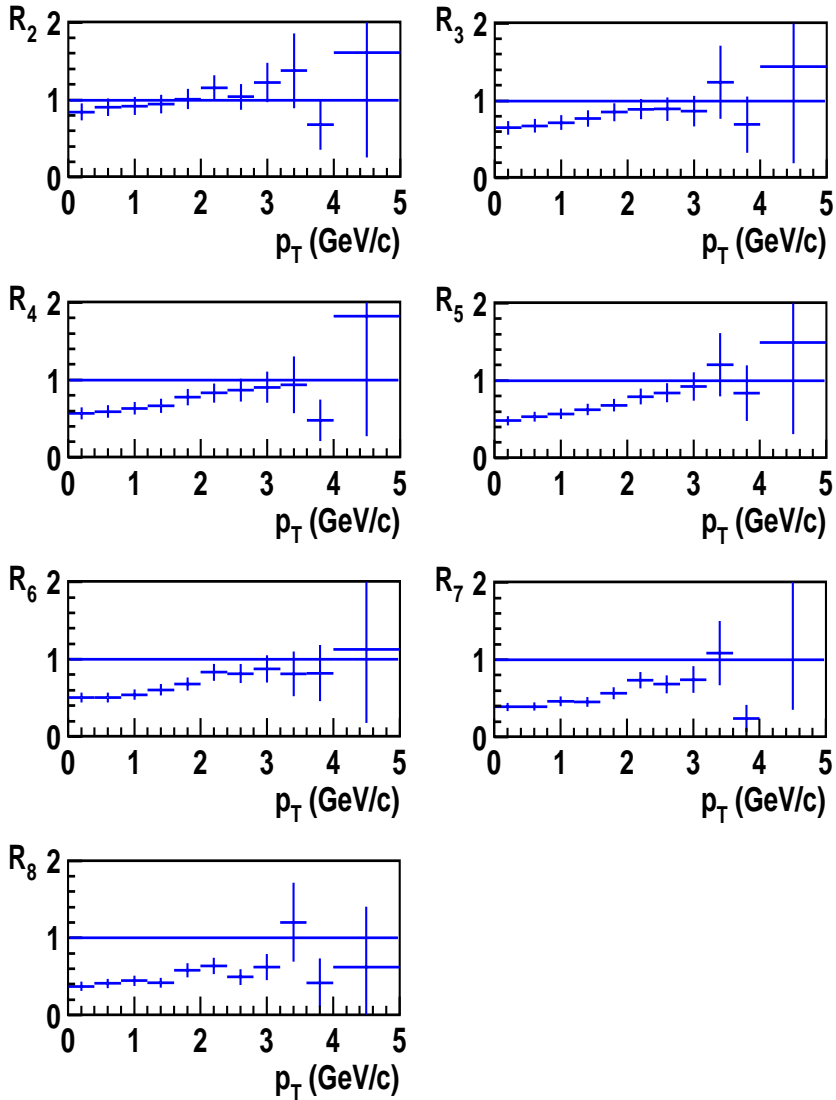


Figure 2. Ratios  $R_i$  of the  $J/\psi$  transverse momentum distribution normalized to the DY cross section in the  $E_i$  bin  $2 < i < 8$  to the first  $E_1$  bin.

significant  $J/\psi$  suppression which strongly increases with centrality, when  $p_T$  increases, the dependence of the  $J/\psi$  normalized yield on centrality becomes weaker and weaker. In other words, the suppression observed on the integrated  $p_T$  yield from peripheral to central collisions originates mainly from the suppression of  $J/\psi$  with low  $p_T$  values.

In order to better visualize this dependence we consider the ratio  $R_i$  of each  $p_T$  distribution corresponding to a given  $E_T$  bin  $i$  with respect to the first and most peripheral bin, namely:

$$R_i = ( J/\psi_i / DY_i ) / ( J/\psi_1 / DY_1 )$$

Fig.2 displays the eight ratios  $R_i$  as a function of  $p_T$ . It shows that with respect to the most peripheral collisions,  $J/\psi$  becomes more and more suppressed, with increasing centrality but also with decreasing  $p_T$  values. For high  $p_T$  values, above 3.5 GeV/c, the suppression although still increasing with centrality, exhibits no significant  $p_T$  dependence.

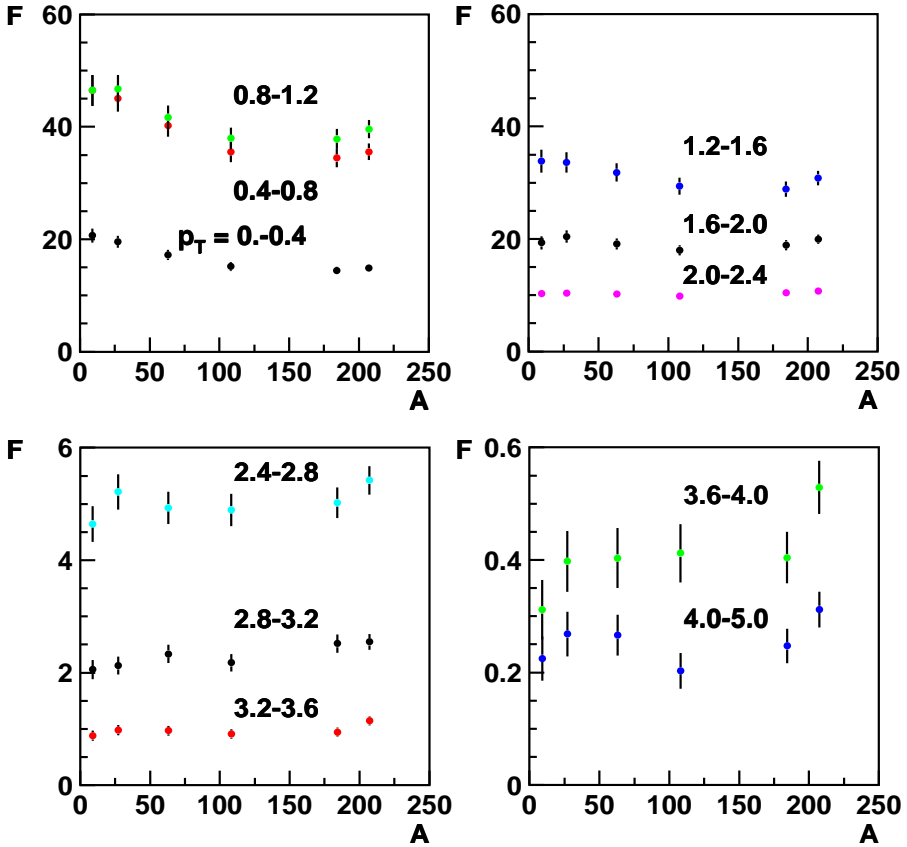


Figure 3. Ratio  $F$  of the  $J/\psi$  production cross section for proton-nucleus collisions in the  $p_T$  bins shown on the plots (in GeV/c) to the DY cross section, as a function of the atomic number of the target nucleus.

Pb-Pb collisions are compared, hereafter, with p-A reactions where the  $J/\psi$  survival probability is affected by normal nuclear absorption only. In this case, when the  $J/\psi$  yield is parametrized according to  $A^\alpha$ , nuclear absorption leads to a value of  $\alpha$  lower than unity reflecting the absorption of the  $c\bar{c}$  pair within the target. The above picture becomes more complex when the survival probability as a function of  $p_T$  is considered. Within the frame of the same NA50 experiment, we have therefore made a study of the  $J/\psi$  yield  $p_T$  dependence for 400 GeV p-induced reactions on 6 different target nuclei: Be, Al, Cu, Ag, W and Pb. We have considered the same 11  $p_T$  bins and have measured the ratio  $F$  in each of them for the six different targets. The results are shown in Fig.3.

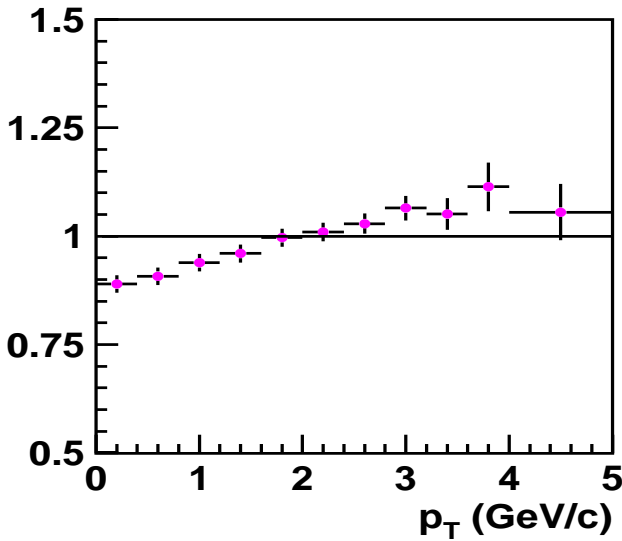


Figure 4. Parameter  $\alpha$  obtained from the fit of the proton-nucleus  $J/\psi$  production cross sections as a function of the transverse momentum (GeV/c).

We have used the above  $A^\alpha$  parametrization of the  $J/\psi$  cross section separately in each of the 11  $p_T$  bins in order to perform a  $p_T$  dependent analysis. The results of this study are illustrated in Fig.4. They show that whereas for low values of  $p_T$   $J/\psi$  production as a function of the atomic mass number  $A$  increases less than proportionally to  $A$  (Drell-Yan is proportional to  $A$  and both are proportional to the number of nucleus-nucleus collisions) leading to a value of  $\alpha$  lower than unity, for high  $p_T$  values  $J/\psi$  production increases faster than  $A$  so that the corresponding value of  $\alpha$  is higher than 1. There is a kind of normal nuclear absorption for the lower  $p_T$  values but the magnitude of this absorption decreases with increasing  $p_T$  then vanishes and turns to overproduction for high  $p_T$  already above 2 GeV/c. This is, in fact, a wellknown behaviour observed since long in the production of hadrons and known as the Cronin effect.

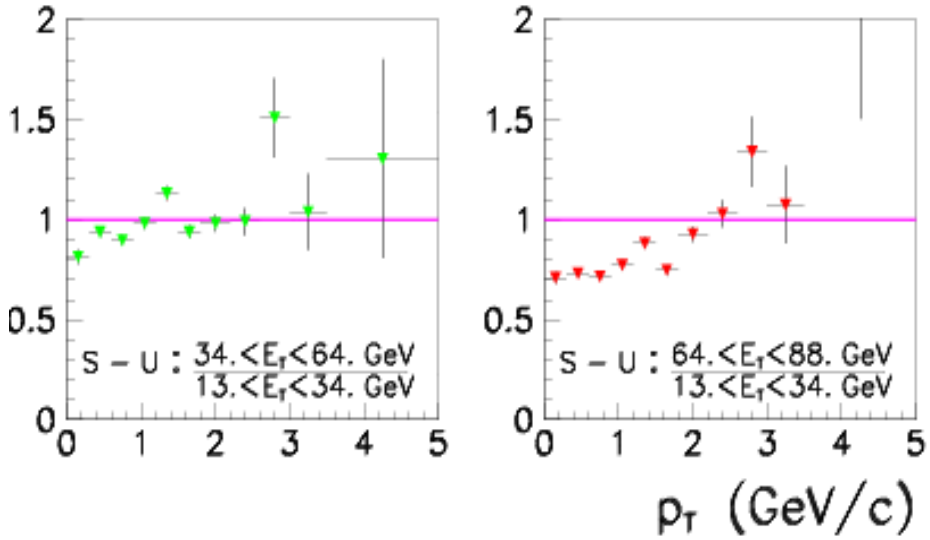


Figure 5. Ratios  $R_i$  of the  $J/\psi$  transverse momentum distribution normalized to the DY cross section for S-U collisions from the NA38 experiment for the case of three  $E_T$  intervals.

For comparison we show in Fig.5 the data for S-U collisions as obtained from the NA38 experiment, where the effect of absorption is seen for low  $p_T$  ( $R < 1$ ), together with some hints of enhancement for high  $p_T$  ( $R > 1$ ) suggesting, within errors, a behaviour similar to the Cronin effect observed in p-A collisions.

The Pb-Pb data can be rebinned using only 3 bins of transverse energy in order to minimize statistical fluctuations. Fig.6 shows that for the most central Pb-Pb collisions and with respect to the most peripheral bin, the suppression exists for all values of  $p_T$ . The centrality dependence decreases with increasing  $p_T$ . For the highest  $p_T$  values, no overproduction is observed: there is always an absorption which increases with centrality, although less pronounced than for small  $p_T$  and which, moreover, does not exhibit any significant  $p_T$  dependence.

### 3. CONCLUSIONS

The dependence of the  $J/\psi$  suppression pattern on  $p_T$  for Pb-Pb collisions is somewhat different from what is observed in the case of normal nuclear  $J/\psi$  absorption from p-induced reactions. In the latter case we see the change from absorption to enhancement with the increase of transverse momentum. For Pb-Pb collisions and for the whole  $p_T$  range, only absorption is observed with increasing centrality even if it is significantly stronger for low  $p_T$ . Moreover, the data suggest that absorption is almost  $p_T$  independent for the most central collisions and for the highest values of the transverse momentum.

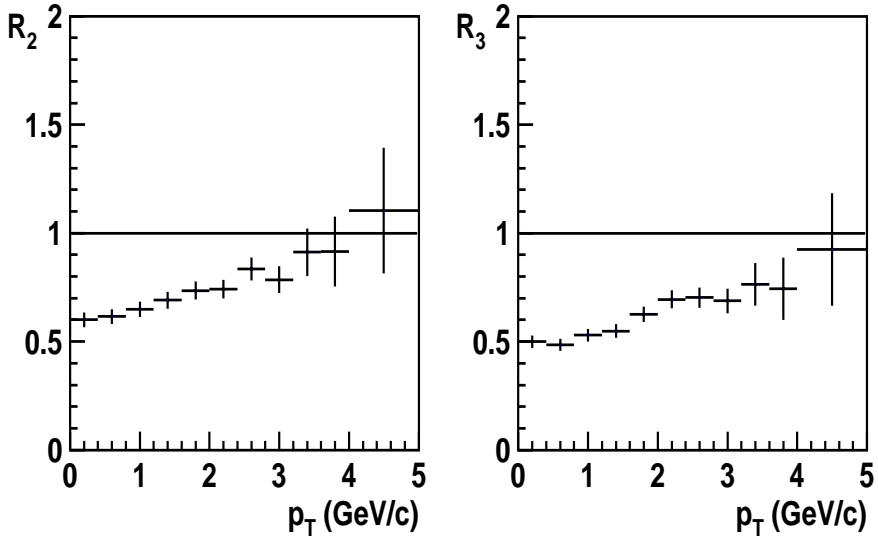


Figure 6. Ratios  $R_2$  and  $R_3$  of the  $J/\psi$  transverse momentum distribution normalized to the DY cross section for the second and third centrality bins with respect to the first and most peripheral one, in the case of three  $E_T$  intervals, for Pb-Pb collisions.

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