

Coulomb and nuclear excitations of ^{70}Zn and ^{68}Ni at intermediate energy

S. Calinescu,¹ O. Sorlin,² I. Matea,³ F. Carstoiu,¹ D. D. Dao,⁴ F. Nowacki,⁴ G. de Angelis,⁵ R. Astabatyán,⁶ S. Bagchi,⁷ C. Borcea,¹ R. Borcea,¹ L. Cáceres,² M. Ciemála,⁸ E. Clément,² Z. Dombrádi,⁹ S. Franchoo,³ A. Gottardo,⁵ S. Grévy,¹⁰ H. Guerin,¹⁰ M. N. Harakeh,⁷ I. M. Harca,¹ O. Kamalou,² M. Kmiecik,⁸ A. Krasznahorkay,⁹ M. Krzysiek,⁸ I. Kuti,⁹ A. Lepailleur,² S. Lukyanov,⁶ A. Maj,⁸ V. Maslov,⁶ K. Mazurek,⁸ P. Morfouace,³ J. Mrazek,¹¹ F. Negoita,¹ M. Niikura,³ L. Olivier,³ Y. Penionzhkevich,⁶ L. Perrot,³ C. Petrone,^{1,12} Z. Podolyák,¹³ C. Rigollet,⁷ T. Roger,² F. Rotaru,¹ D. Sohler,⁹ M. Stanoiu,¹ I. Stefan,³ L. Stuhl,⁹ J. C. Thomas,² Z. Vajta,⁹ M. Vandebrouck,¹⁴ and O. Wieland¹⁵

¹Horia Hulubei National Institute for Physics and Nuclear Engineering, P.O. Box MG-6, 077125 Bucharest-Magurele, Romania

²Grand Accélérateur National d'Ions Lourds (GANIL), CEA/DSM - CNRS/IN2P3, B. P. 55027, F-14076 Caen Cedex 5, France

³Université Paris-Saclay, CNRS/IN2P3, IJCLab, 91405 Orsay, France

⁴Université de Strasbourg, CNRS, IPHC UMR 7178, F-67000 Strasbourg, France

⁵Istituto Nazionale di Fisica Nucleare (INFN), Laboratori Nazionali di Legnaro, Legnaro, Italy

⁶FLNR, JINR, 141980 Dubna, Moscow region, Russia

⁷Nuclear Energy group, ESRIG, University of Groningen, 9747AA Groningen, The Netherlands

⁸The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, ul. Radzikowskiego 152, 31-342 Krakow, Poland

⁹Institute of Nuclear Research (Atomki), P.O. Box 51, H-4001 Debrecen, Hungary

¹⁰Centre d'Études Nucléaires de Bordeaux Gradignan-UMR 5797, CNRS/IN2P3, Université de Bordeaux, Chemin du Solarium, BP 120, 33175 Gradignan, France


¹¹Nuclear Physics Institute, AS CR, CZ 25068, Rez, Czech Republic

¹²Faculty of Physics, University of Bucharest, Magurele, 077125, Romania

¹³Department of Physics, University of Surrey, Guildford GU2 7HX, United Kingdom

¹⁴Irfu, CEA, Université Paris-Saclay, F-91191 Gif-sur-Yvette, France

¹⁵INFN Sezione di Milano, via Celoria 16, 20133, Milano, Italy

 (Received 19 May 2020; revised 8 April 2021; accepted 10 September 2021; published 23 September 2021)

The reduced transition probabilities $B(E2; 0_{g.s.}^+ \rightarrow 2_1^+, 2_2^+)$ in ^{70}Zn and the full $B(E2; 0_{g.s.}^+ \rightarrow 2^+)$ strength up to $S_n = 7.79$ MeV in ^{68}Ni have been determined at the LISE/GANIL facility using the Coulomb-excitation technique at intermediate beam energy on a ^{208}Pb target. The γ rays emitted in-flight were detected with an array of 46BaF_2 crystals. The angles of the deflected nuclei were determined in order to disentangle and extract the Coulomb and nuclear contributions to the excitation of the 2^+ states. The measured $B(E2; 0_{g.s.}^+ \rightarrow 2_1^+)$ of $1432(124) e^2 \text{fm}^4$ for ^{70}Zn falls in the lower part of the published values which clustered either around 1600 or above $2000 e^2 \text{fm}^4$, while the $B(E2; 0_{g.s.}^+ \rightarrow 2_2^+)$ of $53(7) e^2 \text{fm}^4$ agrees very well with the two published values. The relatively low $B(E2; 0_{g.s.}^+ \rightarrow 2_1^+)$ of $301(38) e^2 \text{fm}^4$ for ^{68}Ni agrees with previous studies and confirms a local magicity at $Z = 28$, $N = 40$. Combining the results of the low-energy spectra of ^{68}Ni and ^{70}Zn and their shell-model interpretations, it is interesting to notice that four different shapes (spherical, oblate, prolate, and triaxial) are present. Finally, a summed $E2$ strength of only about $150 e^2 \text{fm}^4$ has been found experimentally at high excitation energy, likely due to proton excitations across the $Z = 28$ gap. The experimental distribution of this high-energy $E2$ excitation agrees with shell-model calculations, but its strength is about two times weaker.

DOI: [10.1103/PhysRevC.104.034318](https://doi.org/10.1103/PhysRevC.104.034318)

I. INTRODUCTION

The $^{68}\text{Ni}_{40}$ nucleus is, as $^{14}\text{C}_8$ and $^{34}\text{Si}_{20}$, composed of a doubly closed shell originating from the spin-orbit interaction for protons and from the harmonic-oscillator-like shape of the mean-field potential for neutrons. As for ^{14}C and ^{34}Si , ^{68}Ni lies at the verge of a so-called island of inversion where the cost of promoting pairs of neutrons across the shell gap is lower than the gain through quadrupole and pairing correlations.

The doubly-closed-shell effect can be best viewed by the sudden increase of the energy of the first 2_1^+ state, which goes

in concert with a decrease of the reduced transition probability $B(E2; 0_1^+ \rightarrow 2_1^+)$ to excite it. Adding two protons to these nuclei preserves a significant effect of double magicity in ^{16}O and ^{36}S , and single magicity in ^{70}Zn , whose structure resembles, according to Ref. [1], a vibrator with some single-particle features. On the other hand, when removing two protons from these nuclei, no increase in the 2_1^+ energy is observed and the $B(E2; 0_{g.s.}^+ \rightarrow 2_1^+)$ values are increasing: $^{12}\text{Be}_8$, $^{32}\text{Mg}_{20}$, and $^{66}\text{Fe}_{40}$ nuclei are deformed (see, e.g., Refs. [2–4]). As discussed in Ref. [5], these nuclei lie in regions where the hierarchy between the proton-neutron forces

TABLE I. $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ values (in units of $e^2 \text{fm}^4$) in ^{70}Zn obtained from different experiments. The last column gives the evaluated value.

	Stelson [19]	Singh [20]	Neuhausen [18]	Sorlin [11]	Kenn [21]	Celikovic [22]	Pritychenko [16]
$B(E2)$	1600(140)	2350(250)	2050(190)	1640(280)	1650(190)	1515(125)	1559(80)

in presence plays a crucial role in changing the spacing (and even the ordering) between orbits, thus favoring the breaking of magicity at $N = 8, 20$, and 40 .

The ^{34}Si and ^{68}Ni nuclei carry fingerprints of this proximity to the island of inversion through the presence of a 0_2^+ state below the 2_1^+ state. In these two cases, the determination of the reduced transition probabilities $\rho(E0)$ and $B(E2; 0_2^+ \rightarrow 2_1^+)$ suggests a shape coexistence between spherical 0_1^+ and deformed 0_2^+ configurations (see, e.g., Ref. [6] for ^{34}Si and Refs. [7–9] for ^{68}Ni). Their weak $B(E2; 0_1^+ \rightarrow 2_1^+)$ values [10,11] have been interpreted as due to the fact that the 2_1^+ state mainly comes from neutron excitation(s). As neutrons carry a much smaller effective charge than protons, $e_v \simeq 0.5$ as compared with $e_\pi \simeq 1.5$, the resulting neutron contribution to the $B(E2)$ value, which scales with e^2 , is small compared with a similar excitation involving protons.

In the three nuclei $^{14}_8\text{C}$, $^{34}_{20}\text{Si}$, and $^{68}_{40}\text{C}$, neutron excitations across the shell gap involve a change of parity from the normally occupied to the valence states. Therefore, the structure of the 2_1^+ state must involve 2p2h excitations, with the breaking of a pair. Consequently, the energy of the 2_1^+ state is not directly related to the size of the neutron shell gap but to a more complex energy balance that involves pairing and quadrupole correlations, which reduce its excitation energy as compared with the size of the shell gap. The size of the $N = 40$ shell gap can be more directly inferred from the energy of negative-parity states coming from the (1p1h) excitation across the gap. In ^{68}Ni , the long-lived (0.86 ms) 5^- state at 2.849 MeV [12], formed by $\nu p_{1/2}$ or $\nu f_{5/2}$ excitations to the $\nu g_{9/2}$ orbit, is such a candidate.

Contrary to neutrons, proton excitations across the gap in ^{34}Si and ^{68}Ni occur between orbits of the same parity, e.g., between the $0f_{7/2}$ and $(1p_{3/2}, 0f_{5/2})$ orbitals in ^{68}Ni . Therefore, a proton 2_p^+ state can be constructed already at the 1p1h level and should be existing in the spectra of ^{34}Si and ^{68}Ni . The corresponding $B(E2; 0_1^+ \rightarrow 2_p^+)$ value could be larger than that for the 2_1^+ state, depending on the rigidity of the proton shell gap against ph excitations, the fragmentation of the strength, as well as the proton ph content of the ground state. In the case of ^{34}Si , similar $B(E2)$ values have been predicted for the 2_1^+ and 2_p^+ states [13], but a firm identification of 2_p^+ state is still lacking. In ^{68}Ni , shell-model calculations of Langanke *et al.* [14] have predicted that a large fraction (about twice as large as that of the 2_1^+ state) of the $B(E2)$ strength goes to 2_p^+ states present above 4 MeV.

In the present work, we wish to confirm the low $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ value in ^{68}Ni , found to be 255(60) in Ref. [11] and $280_{-100}^{+120} e^2 \text{fm}^4$ in Ref. [15], and explore if significant $B(E2)$ strength exists at higher excitation

energy. We have also remeasured the $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ and $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+)$ values in ^{70}Zn , for which various experiments gave rather inconsistent $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ values. As shown in Table I, four results cluster around $1600 e^2 \text{fm}^4$, while two others are above $2000 e^2 \text{fm}^4$. The last column gives the evaluated $B(E2)$ values [16].

The puzzling discrepancy between the determined $B(E2)$ values in ^{70}Zn could point to a side-feeding of the 2_1^+ from the 2_2^+ state. If its contribution was large and not subtracted, it would artificially enhance the measured $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$. However, with a branching ratio from the 2_2^+ to the 2_1^+ of 60% [17] (the remaining going to the g.s.), and a low $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+)$ of about $50 e^2 \text{fm}^4$ [1,18], a possible contamination would be only of about $0.6 \times 50 = 30 e^2 \text{fm}^4$. This would therefore not explain the discrepancies on the reported $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ values observed in Table I.

The $B(E2)$ values of ^{70}Zn and ^{68}Ni were obtained in the present work by Coulomb interaction with a ^{208}Pb target, using the same experimental technique as in Ref. [23]. In the following, we describe the production methods of the beams of interest, followed by a brief presentation of the experimental setup and of the methodology to extract the Coulomb and nuclear quadrupole excitation strengths, which is then applied to the two nuclei. The search for measurable $B(E2)$ strength at higher excitation energies in ^{68}Ni is addressed in the last experimental part of this work, followed by a theoretical interpretation of the structure of ^{70}Zn and ^{68}Ni using shell-model calculations.

II. EXPERIMENT

The ^{68}Ni nuclei were produced at GANIL in fragmentation reactions of a 60A MeV $^{70}\text{Zn}^{29+}$ beam, with an average intensity of $\approx 1.2 \mu\text{Ae}$ on a $145 \mu\text{m}$ ^9Be target. They were separated from other reaction products by the LISE3 spectrometer [24], using a wedge-shaped Be degrader of $221 \mu\text{m}$ inserted at its intermediate focal plane. Coulomb excitation was induced on a $200 \text{mg}/\text{cm}^2$ ^{208}Pb target, placed at the image focal plane of the spectrometer. In addition, a $146.4 \text{mg}/\text{cm}^2$ ^{12}C target was used to study more specifically the nuclear contribution. In total, 2.6×10^9 ^{68}Ni nuclei were produced for the ^{208}Pb target study with a purity of 80% at energies of 47.68A and 40.8A MeV before and at mid-target depth, respectively.

The ^{70}Zn primary beam was also used at a reduced intensity with an energy of 51.8A MeV to study the $B(E2)$ values of the 2_1^+ and 2_2^+ states. In total, 1.75×10^8 ^{70}Zn nuclei reached the secondary ^{208}Pb target with a purity of 100% and an energy at mid-target depth of 44.4A MeV.

The incoming beam profile and interaction point on target were determined by using two sets of position-sensitive gas-filled detectors (CATS) [25] located 140 and 75 cm upstream of the Pb target.

Forty-six hexagonal BaF_2 crystals of the Château de Cristal were placed in two hemispheres at a mean distance of 25 cm from the ^{208}Pb target for the detection of the in-flight γ rays, emitted within a 3 ns time window [26] with respect to the passage of the ions through the CATS detectors. The efficiency of this array was first determined using calibration sources up to 1.4 MeV, and cross-checked using GEANT4 simulations. Then, the in-flight efficiency was simulated by taking into account the angular distribution of the emitted $E2$ γ rays, calculated by the DWEIKO code [27] for the different cases under study. We confirmed *a posteriori* that these distributions agree with the experimental ones. A photopeak efficiency of 25(1)% was found at 885 keV, which is the energy of the $2_1^+ \rightarrow 0_1^+$ transition in ^{70}Zn . As this transition has been seen in each individual crystal of BaF_2 , it has also been used to optimize the Doppler-shift correction by comparing the observed γ peak for each crystal, to the known energy at rest. It was assumed that the γ -ray emission occurred at mid-target depth to apply Doppler corrections, which is a good approximation for states with short lifetimes (here shorter than 4 ps). According to the simulations, the intensity of the photopeak signal is still dominant at γ -ray energies below 8 MeV, above which the single-escape peak becomes dominant.

The deflected nuclei were identified by their time-of-flight and energy loss in a double-sided silicon-strip annular detector (DSSSD) located 50 cm downstream of the target. The DSSSD consisted of four quadrants having on the front side 16 annular strips of 1.9 mm width each, and on the back side 24 radial strips of 3.4° pitch, each grouped three by three. With a central hole of 3 cm, this geometry allowed us to detect and identify ions scattered between 1° and 6.5° in the laboratory frame. The efficiencies of the DSSSD to detect nuclei that underwent Coulomb excitation were determined using GEANT4 simulations by shooting the target with beam particles having the measured position and energy profiles. Their angular distributions were calculated with the DWEIKO [27] and ECIS [28] codes corresponding to 2^+ state excitations. They were folded to account for the angular straggling (≈ 0.6 degrees) induced by interactions between the beams and the Pb target. After having removed the few dead strips from the analysis, DSSSD detection efficiencies of 62% and 72% were found for ^{70}Zn and ^{68}Ni , respectively.

The nuclei that passed through the central hole of the DSSSD were identified in an ionization chamber (CHIO) and stopped in a plastic detector, which was surrounded by two HPGe detectors, used to determine the isomeric content of the ^{68}Ni beam. The production of two isomeric states of ^{68}Ni , 5^- with $T_{1/2}$ of 860(50) μs [12] and 0_2^+ with $T_{1/2}$ of 268(12) ns [8], was quantified from the characteristic delayed γ -ray transitions observed in Fig. 1 after the implantation of ^{68}Ni nuclei. The 815 keV transition that connects the 5^- state to the 2^+ state at 2034 keV was used to determine that 31(2)% of the incoming ^{68}Ni were in the 5^- isomeric state. The number of counts in the 511-keV peak was used to infer an isomeric ratio of 2.45(2)% for the 0_2^+ at 1604 keV [7,8], obtained after

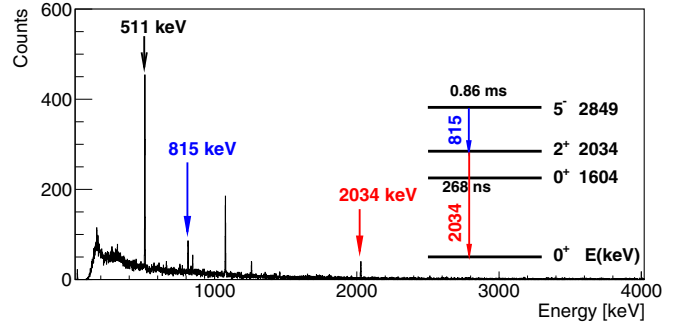


FIG. 1. γ -ray spectrum of one HPGe detector associated with the deexcitation of ^{68}Ni implanted in the plastic detector. The lines corresponding to the decays of the 0_2^+ and 5^- isomers are indicated by arrows.

correcting for the fact that only two-thirds of its decay occurs by pair creation [29]. The other γ -ray transitions observed in Fig. 1 come from the β decay of ^{68}Ni [30].

III. EXPERIMENTAL RESULTS

A. Coulomb excitation of ^{70}Zn

The Coulomb interaction dominates over the nuclear interaction at small deflection angles, where the impact parameter b of the reaction is larger than the range of the nuclear force. Using semiclassical trajectories [31], a minimal impact parameter b_{\min} , above which nuclear excitations can be neglected, has been derived in Refs. [32,33]. It is linked to a maximal scattering angle $\theta_{\text{lab}}^{\max}$, up to which Coulomb excitation dominates, that can be approximated by [34]

$$\theta_{\text{lab}}^{\max} = \frac{2Z_p Z_t e^2}{b_{\min} M_p c^2 \beta^2 \gamma}, \quad (1)$$

where the subscripts “ p ” and “ t ” stand for the projectile and target, respectively. In this equation, the minimal impact parameter b_{\min} is obtained as follows:

$$b_{\min} = R_p + R_t + \delta + \frac{\pi a}{2\gamma} \text{ fm}, \quad (2)$$

with

$$R_{p(t)} = r_0 A_{p(t)}^{1/3} \text{ fm}, \quad (3)$$

where $\delta = 6$ fm, $r_0 = 1.3$ fm, $A_{p(t)}$ is the mass of the projectile and target, respectively, and

$$a = \frac{Z_p Z_t e^2}{M_p c^2 \beta^2}, \quad \gamma = \frac{1}{\sqrt{1 - \beta^2}}. \quad (4)$$

Using the present mean velocity value $\langle \beta \rangle$ of 0.298, it is found that the Coulomb process is dominant for impact parameters larger than $b_{\min} = 20.29$ fm, i.e., for deflection angles lower than $\theta_{\text{lab}}^{\max} = 3.3^\circ$. This is significantly smaller than the grazing angle of the reaction that is here 4.5° . However, as we shall discuss later, at this intermediate energy regime, there is no angular range in which the nuclear part is totally negligible.

The Doppler-corrected and background-subtracted γ -ray spectrum of ^{70}Zn , displayed in Fig. 2, has been obtained

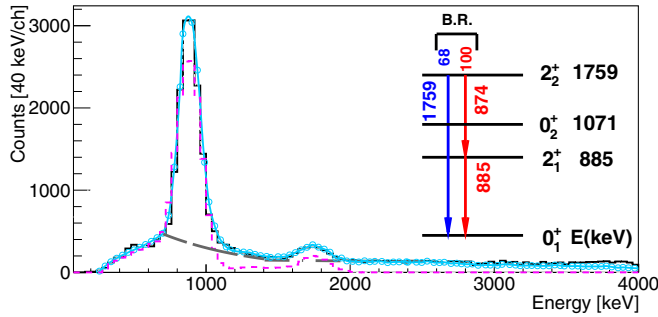


FIG. 2. Doppler-corrected and background-subtracted γ -ray spectrum of ^{70}Zn measured with Château de Cristal (black histogram). The intensities of the two peaks were determined using the simulated response function of the array (dashed magenta line) on top of two exponential functions accounting for the background (long-dashed gray line). The result of the best-fit procedure is illustrated with blue circles.

using the full angular range up to 6.5° . The background spectrum was obtained by collecting radiations out of the prompt time-window, which mostly displays structures arising from the intrinsic radioactivity of the BaF_2 crystals (see, e.g., Ref. [26]). Two γ peaks are seen in Fig. 2, an unresolved doublet at around 880 keV and a single-peak at 1759 keV. The latter transition results from the deexcitation of the 2_2^+ state to the g.s. This level decays directly to the g.s., with a branching ratio of 68/168, and to the 2_1^+ state via a cascade of two γ rays of 874 and 885 keV, with a ratio of 100/168 [17], both ratios having 10% error. Therefore, when extracting the reduced transition probability for the 2_1^+ state, the feeding from the 2_2^+ state has to be subtracted, which was not done in most of the works presented in Table I.

The $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ and $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+)$ values of ^{70}Zn are obtained from the comparison between experimental and simulated differential cross sections, the latter being obtained from the DWEIKO [27] and ECIS [28] codes. The use of these two codes allows us to estimate the error made in the derivation of the $B(E2)$ values with two different methods. As there exists no optical potential available for our target-projectile-energy combination, the parameters derived from the $^{86}\text{Kr} + ^{208}\text{Pb}$ reaction at 43 MeV/A [35] were used in our fits. To see the influence of the choice of potential on the fit of experimental data, two additional potentials $^{40}\text{Ar} + ^{208}\text{Pb}$ at 44 MeV/A [36] and $^{17}\text{O} + ^{208}\text{Pb}$ at 84 MeV/A [37] were used. Very similar values (with less than 4% difference) were found in the determination of the nuclear parameter β_{N} using these three optical potentials. This stability of the extracted nuclear parameter over the choice of optical model potential parameters is due to the fact that only the extreme tail of the potential influences the reaction mechanism.

Since the energy resolution of Château de Cristal (typically 18% at 885 keV) does not allow a discrimination between the 885 and 874 keV transitions emitted in the $2_2^+ \rightarrow 2_1^+ \rightarrow 0_1^+$ cascade, the best approach is to first fit the experimental angular distribution of the 2_2^+ state, obtained by gating on the 1759-keV γ ray, which corresponds to the direct decay to the g.s. Once the Coulomb (β_{C_2}) and nuclear (β_{N_2}) deformations

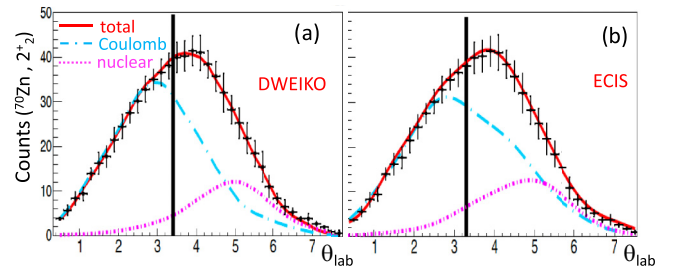


FIG. 3. Comparison between the experimental angular distribution (black crosses) of the deflected ^{70}Zn nuclei in the laboratory frame, gated on the 1759-keV peak of Fig. 2, with the simulated ones (continuous red lines) obtained with the (a) DWEIKO and (b) ECIS codes, leading to $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+)$ values of 62(7) and 55(6) $e^2 \text{fm}^4$, respectively. The Coulomb and nuclear contributions are shown with blue (long-dashed-dotted lines) and magenta (dotted lines), respectively. See text and Table II for the corresponding β_{C} and β_{N} parameters. The black vertical line marks the maximum safe scattering angle of the reaction. θ_{lab} is given in degrees.

of this 2_2^+ state are determined, they can be used to constrain the fit of the angular distribution gated on the (874, 885)-keV doublet in order to determine the β_{C_1} and β_{N_1} values of the 2_1^+ state.

The Coulomb and nuclear processes are present at all scattering angles and interfere. However, as the Coulomb excitation at intermediate energy is usually the dominating contribution below the “safe” scattering angle, the fit of the angular distribution was first done using only β_{C_2} up to 3° in Fig. 3. The β_{N_2} parameter was set in order to reach a qualitative agreement for the high angular range of the spectrum. This way, a first couple of approximated β_{C_2} and β_{N_2} values was derived and used as initial parameters of a χ^2 minimization procedure applied to the full angular distribution. The values of β_{C_2} and β_{N_2} , obtained after convergence of the fit, are listed in Table II. The same procedure has been applied for other cases described later in the text.

The best agreement between the simulated (continuous red line) and the experimental (black crosses) amplitude of the angular distribution, gated on the 1759-keV γ -ray, is obtained for $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+) = 62(7) e^2 \text{fm}^4$ [$\beta_{\text{C}_2} = 0.045(3)$], and a nuclear contribution of β_{N_2} of 0.054(3) with the DWEIKO code. The best fit with ECIS is for $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+) = 55(6) e^2 \text{fm}^4$ [$\beta_{\text{C}_2} = 0.042(3)$] and $\beta_{\text{N}_2} = 0.049(3)$. The resulting simulations of the angular distributions of the deflected ^{70}Zn nuclei for the two codes are shown in Fig. 3.

At this relatively high energy of the scattered ions, the double-step excitation of the 2_2^+ state through the intermediate 2_1^+ state is expected to be small. It has been estimated to be 0.7 mb using the GOSIA code [38], as compared with the presently measured $0_{\text{g.s.}}^+ \rightarrow 2_2^+$ cross section of 7.8(9) mb over the 1° - 3.3° angular range. Subtracting this estimated contribution from the mean value of the two fits, 58(7) $e^2 \text{fm}^4$, leads to 53(7) $e^2 \text{fm}^4$, which is in excellent agreement with the values of 50(13) and 50(10) $e^2 \text{fm}^4$, obtained in Ref. [18] and [1]. The present experimental results are reported in Table II.

The experimental cross section for the 2_1^+ state is found to be 228(20) mb in the angular range between 1° and 3.3° .

TABLE II. Experimental cross sections, σ (in mb), between 1° and $\theta_{\text{lab}}^{\text{max}}$, are reported together with β_C , β_N , and $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2^+)$ (in $e^2 \text{fm}^4$) obtained with the DWEIKO and ECIS codes for ^{70}Zn and ^{68}Ni .

Nucleus	J^π	$\theta_{\text{lab}}^{\text{max}}$ (deg)	σ (mb)	DWEIKO			ECIS			Adopted
				β_C	β_N	$B(E2)$	β_C	β_N	$B(E2)$	$B(E2)$
^{70}Zn	2_1^+	3.3	228(20)	0.215(8)	0.288(12)	1419(120)	0.217(9)	0.241(8)	1447(127)	1432(124)
^{70}Zn	2_2^+	3.3	7.8(9)	0.045(3)	0.054(3)	62(7)	0.042(2)	0.049(3)	55(6)	53(7) ^a
^{68}Ni	2_1^+	3.4	55(7)	0.107(6)	0.102(8)	292(35)	0.110(7)	0.099(7)	312(40)	301(38)

^aThe calculated double-step feeding contribution to the 2_2^+ state has been subtracted from the $B(E2)$ value of 58(7) $e^2 \text{fm}^4$ obtained from the mean results of the two codes, leading to the adopted value of 53(7) $e^2 \text{fm}^4$ (see text for details).

The angular distribution corresponding to the excitation of the 2_1^+ state includes the contamination of the 2_2^+ state that has just been derived earlier and for which we keep the contributions (β_{C_2} , β_{N_2}) as fixed. The fit of the global angular distribution with the DWEIKO code (continuous red line) in Fig. 4(a) includes a Coulomb part with $\beta_{C_1} = 0.215(8)$, which corresponds to a $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ value of 1419(120) $e^2 \text{fm}^4$ and a nuclear contribution of $\beta_{N_1} = 0.288(12)$. When using the ECIS code instead [Fig. 4(b)], a comparable $B(E2)$ value of 1447(127) $e^2 \text{fm}^4$ is found. While the β_{C_1} value is very similar to that obtained with the DWEIKO code, the fitted β_{N_1} value of 0.241 is significantly smaller and not compatible with the value of 0.288 obtained with DWEIKO.

Our adopted $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ of 1432(124) $e^2 \text{fm}^4$ corresponds to the mean value extracted using the DWEIKO or ECIS codes (see Table II). It agrees, within the error bars, with those of Refs. [11,19,21,22] and with the evaluated value of 1525(75) $e^2 \text{fm}^4$ [16]. It confirms that the $B(E2)$ values above 2000 $e^2 \text{fm}^4$ obtained from inelastic scattering of electrons [18] and protons [20] are too large.

B. Coulomb and nuclear excitations of ^{68}Ni

1. Excitation of the 2_1^+ state on the ^{208}Pb target

The reduced transition probability $B(E2)$ to the 2_1^+ state at 2.034 MeV in ^{68}Ni has been obtained using the same procedure as described above for ^{70}Zn . The fraction of ^{68}Ni fragments produced in an isomeric state has been subtracted from the number of incoming ^{68}Ni nuclei. In this assumption,

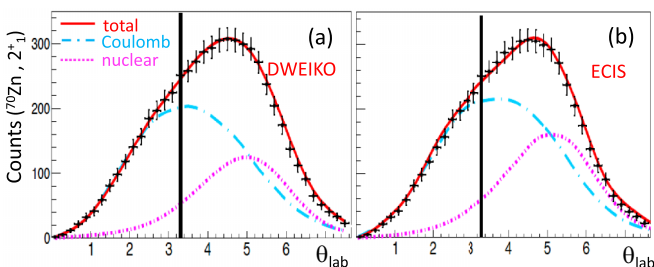


FIG. 4. Same as Fig. 3 but gated on the (874, 885)-keV doublet. (a) The simulated angular distribution (continuous red line) is obtained using $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+) = 1419(120)$ $e^2 \text{fm}^4$ and $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+) = 62(7)$ $e^2 \text{fm}^4$ with the DWEIKO code. (b) Values of $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+) = 1447(127)$ $e^2 \text{fm}^4$ and $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+) = 55(6)$ $e^2 \text{fm}^4$ are obtained with the ECIS code. θ_{lab} is given in degrees.

we neglect the contribution of their excitation to higher excited states (especially for the 5^- isomeric state whose beam fraction is large) that would decay through the 2_1^+ state. This hypothesis is further justified by the fact that we do not see any γ ray in coincidence with the line at 2.034 MeV.

With a mean $\langle\beta\rangle = v/c$ value of 0.286 and $b_{\text{min}} = 20.3$ fm, a value of $\theta_{\text{lab}}^{\text{max}} = 3.4^\circ$ is calculated from Eq. (1). The Doppler-corrected background-subtracted γ -ray spectrum in Fig. 5(a), obtained by including the angular distribution of the ejectile up to 6.5 degrees, clearly shows a 2.034 MeV line corresponding to the excitation of the 2_1^+ state. The integrated

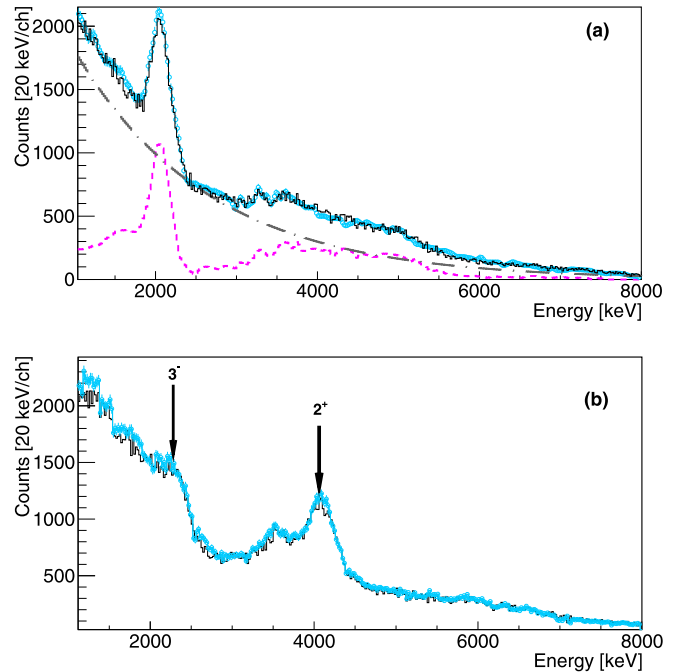


FIG. 5. (a) Doppler-corrected and background-subtracted γ -ray spectrum conditioned by the detection of scattered ^{68}Ni in the DSSSD (black histogram). The number of counts in the 2_1^+ state at 2.034 MeV was extracted by using the simulated response function of the detection system (dashed magenta line). The curve in blue circles corresponds to the global fit, which includes an exponential decay (long-dashed gray line), the excitations of the 3^- and 2_1^+ states in ^{208}Pb , and high-energy contributions discussed in Sec. III B 3. (b) Same as panel (a), but for the non-Doppler-corrected spectrum. The decays of the 3^- and 2_1^+ states in ^{208}Pb are visible at 2.614 and 4.085 MeV (with single-escape), respectively, for which our measured $B(E3)$ and $B(E2)$ values agree with those of the literature.

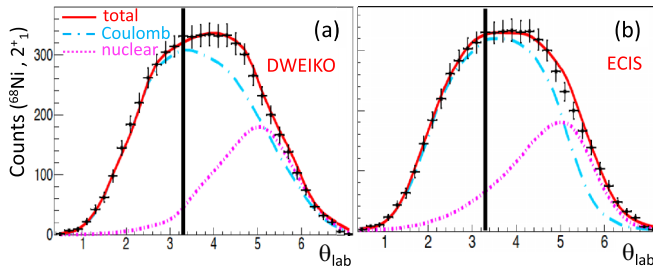


FIG. 6. Same as Fig. 3 but for ^{68}Ni . $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ values of 292(35) and 312(40) $e^2 \text{fm}^4$ are obtained from the comparison between the experimental distribution (black crosses) with the (a) DWEIKO and (b) ECIS codes, respectively. θ_{lab} is given in degrees.

cross section between 1° and 3.4° is 55(7) mb. The complex fit of the high-energy part of the spectrum will be discussed in Sec. III B 3.

Figure 5(b) displays the non-Doppler-corrected spectrum of ^{68}Ni (black histogram), together with the obtained fit (open blue circles). A clear photopeak is seen at about 4 MeV with the related single-escape peak. This double-peak structure corresponds to the excitation of the 2_1^+ state at 4.085 MeV in ^{208}Pb , for which a $B(E2)$ value of 3025(421) $e^2 \text{fm}^4$ is measured, compatible with the adopted value of 3010(160) $e^2 \text{fm}^4$ [16]. In the high-energy tail of the in-flight component of the 2_1^+ state in ^{68}Ni , the excitation of the 3_1^- state in ^{208}Pb is also observed for which a $B(E3; 0_{\text{g.s.}}^+ \rightarrow 3_1^-)$ value of 652 000(9000) $e^2 \text{fm}^6$ was found, in agreement with Ref. [39].

The fit of the angular distribution of the ejectiles, observed in coincidence with the γ ray (red curve in Fig. 6) from the 2_1^+ state, includes Coulomb and nuclear parts whose parameters β_C and β_N are listed in Table II. $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ values of 292(35) and 312(40) $e^2 \text{fm}^4$ are obtained when using the DWEIKO and ECIS codes, respectively. The adopted $B(E2)$ value of 301(38) $e^2 \text{fm}^4$ from these two results agrees, within one sigma, with the $B(E2)$ values of 255(60) and 280(60) $e^2 \text{fm}^4$ obtained in Ref. [11] directly and relative to ^{70}Zn , respectively. It also agrees with the less accurate value of 280_{-100}^{+120} $e^2 \text{fm}^4$ from Ref. [15].

2. Excitation of the 2_1^+ state on Carbon target

Interactions between the ^{68}Ni nuclei and a C target were also studied in order to confirm the small nuclear contribution $\beta_N \simeq 0.1$ derived from the study with the Pb target. As seen in Fig. 7, the experimental Doppler-corrected γ -ray spectrum associated with the C target (solid black histogram) does not show evidence of the 2_1^+ state excitation at 2.034 MeV. This indicates that the β_N value is below the sensitivity of the present measurement. This hypothesis is confirmed by the fact that the spectrum shown in magenta stars of Fig. 7, which displays the result of a GEANT4 simulation including the experimental background contribution and a β_N value of 0.10 derived from the excitation with the ^{208}Pb target, does not exhibit any peak at 2 MeV. As shown in the dashed blue spectrum of Fig. 7, a γ transition at the energy of the 2_1^+ state would have been seen in case of $\beta_N \gtrsim 0.2$.

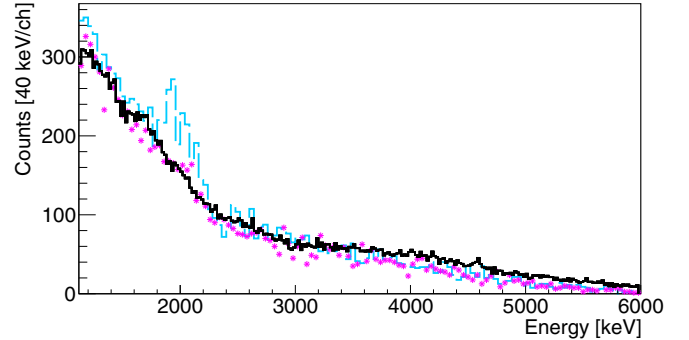


FIG. 7. Comparison between the experimental Doppler-corrected energy spectrum for a ^{68}Ni beam on a C target (solid black histogram) and simulated spectra assuming $\beta_N = 0.1$ (magenta stars) and $\beta_N = 0.2$ (dashed blue histogram). The adopted value of $\beta_N = 0.102(8)$, as deduced from the study with the Pb target, is too small for a peak to be seen when ^{68}Ni scatters off the C target.

3. Excitation of the higher-energy states

The structures that appear between 2.5 and 8 MeV in the Doppler-corrected spectrum of Fig. 5(a) come from some processes that are discussed now.

Figure 8(a) displays the Doppler-corrected ^{68}Ni spectrum. The best fit (open blue circles of Fig. 8) was obtained by introducing in-flight components at 3.3, 5.5, 6.3, 7.5, 8.7, and 9.3 MeV (magenta squares) and taking into account the

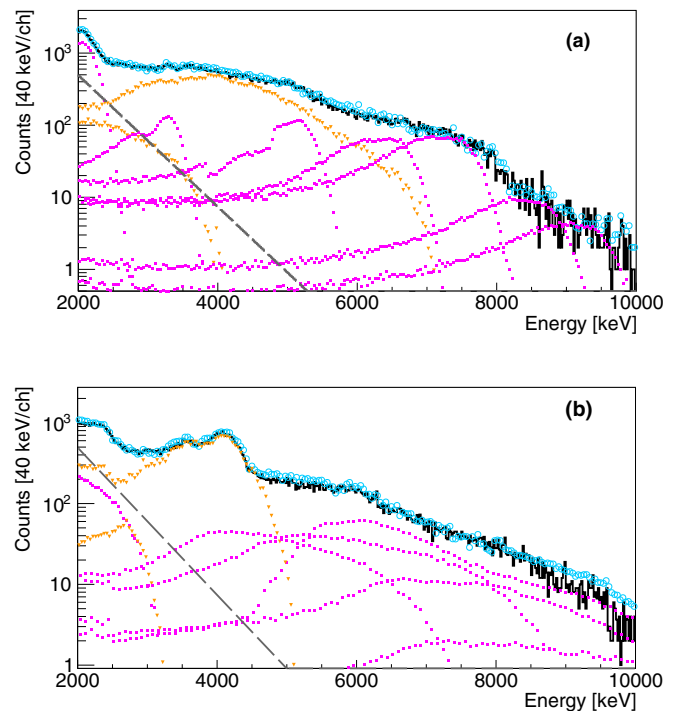


FIG. 8. (a) Doppler-corrected experimental γ -ray spectra of ^{68}Ni (black histogram) compared with GEANT4 simulations (open blue circles) including the excitation of several in-flight γ -ray components (magenta squares) discussed in the text and the excitation of the 2_1^+ and 3_1^- states of the ^{208}Pb target (orange triangles). (b) Same as panel (a) but not Doppler corrected.

excitation of the 2_1^+ and 3_1^- states in the ^{208}Pb target (orange triangles). Note that, as the single-escape component dominates above 8 MeV, the two latter peaks are observed 511 keV lower than the value used in the fit. We do not find evidence of transitions that cascade through the 2_1^+ state. By assuming that the four states below $S_n = 7.8$ MeV belong to $E2$ excitations, strengths of 15(2), 110(15), 42(5), and 16(2) $e^2\text{fm}^4$ are found. As for the states above S_n , they likely correspond to collective $E1$ excitations (likely to be Pigmy dipole resonance modes) whose much shorter lifetime, as compared with $E2$, make their γ decay not completely negligible as compared with neutron decay. Figure 8(b) corresponds to the non Doppler-corrected spectrum fit with the same components as the top spectrum. The consistent agreement of the fitting procedure for the top and bottom spectra gives more confidence in the proposed $E2$ components.

This experimental distribution of $E2$ strength is compared in Fig. 11 to shell-model calculations, whose details will be discussed in Sec. IV B.

IV. DISCUSSION

To interpret the present experimental results, we performed shell-model (SM) calculations for the $E2$ response in ^{70}Zn and ^{68}Ni . These theoretical results are supplemented by deformed Hartree-Fock and beyond-mean-field calculations in order to assess the collective features from a complementary approach. The valence space incorporates the needed degrees of freedom for the description of collectivity and the breaking of $Z = 28$ and $N = 40$ cores. It is composed of the pf shells for protons and $2p_{3/2}1f_{5/2}2p_{1/2}1g_{9/2}2d_{5/2}$ orbitals for neutrons. We used the LNPS effective interaction [40] with recent small adjustments to extend its reliability up to $N = 50$, including $g_{9/2}-d_{5/2}$ particle-hole excitations, with minor consequences at $N = 40$. All electromagnetic properties are calculated using the microscopic effective charges $(e_p, e_n) = (1.31, 0.46)$ from Refs. [41,42].

A. ^{70}Zn

The energies and $E2$ rates are compared with the full shell-model diagonalization in Fig. 9 (middle panel, level scheme labeled as “Shell Model”). Note that the $B(E2 \downarrow)$ values are indicated downwards in Fig. 9. They should be multiplied by 5 for the upwards values, such as for the $0^+ \rightarrow 2^+$ transition given in Table II.

The ^{70}Zn nucleus exhibits a rather complex spectrum with the observation of a low-lying excited 0_2^+ state (at 1.07 MeV excitation energy) and several 2^+ states below 2.0 MeV (at 0.87, 1.76, and 1.96 MeV excitation energy). The agreement is generally good for both the energies and $B(E2)$ values: the calculations produce at the same time a strong $B(E2; 2_1^+ \rightarrow 0_{\text{g.s.}}^+)$ with a value of 314 $e^2\text{fm}^4$ compared with 286 $e^2\text{fm}^4$ experimentally, as well as hindered $B(E2; 2_2^+ \rightarrow 0_{\text{g.s.}}^+)$ and $B(E2; 2_3^+ \rightarrow 0_{\text{g.s.}}^+)$ with values of 1 and 2 $e^2\text{fm}^4$, respectively, compared with the experimental value of 10.6 $e^2\text{fm}^4$ for the $2_3^+ \rightarrow 0_{\text{g.s.}}^+$ transition.

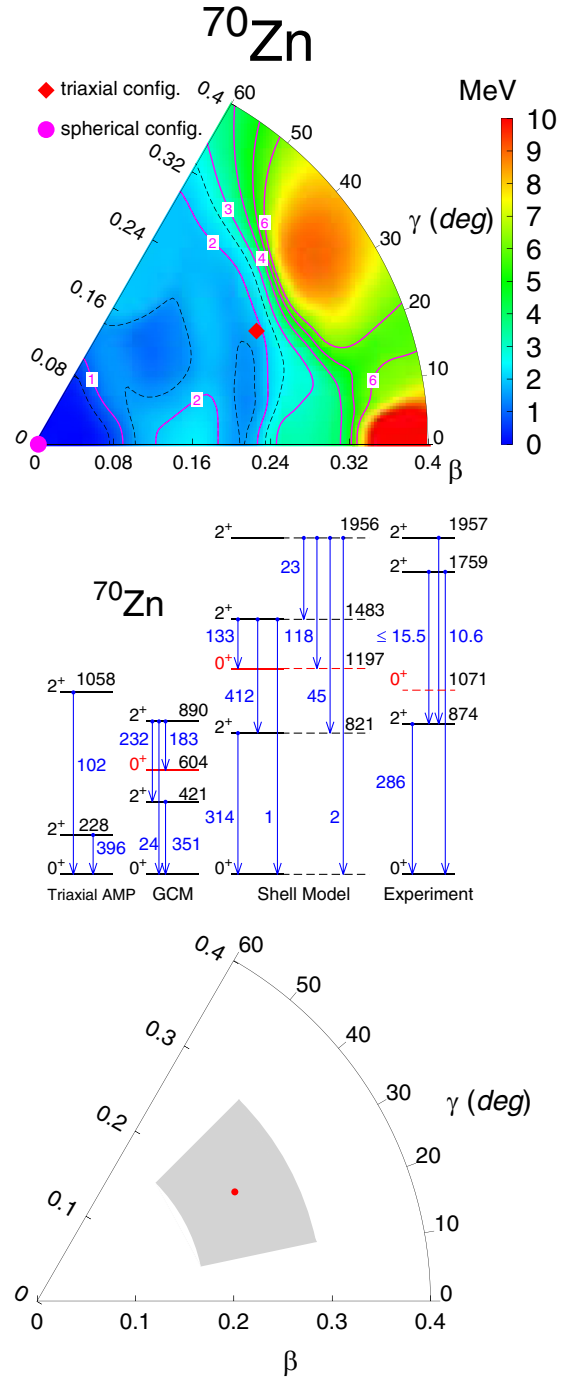


FIG. 9. (a) Potential-energy surface from constrained Hartree-Fock minimization in the shell-model basis for ^{70}Zn . (b) Level scheme and $B(E2 \downarrow)$ values (in MeV and $e^2\text{fm}^4$, respectively) for different theoretical models described in the text and from the present experiment (except for the 2_3^+ state observed in another work and whose $B(E2)$ value is unknown). (c) 1σ contours in the β - γ sextant for the ground-state of ^{70}Zn . The red dot represents the effective β and γ value.

The branchings of the 2_2^+ state, either decaying directly to the ground state, or through a cascade to the 2_1^+ state are known to be of 68/168 and 100/168, respectively (see

inset of Fig. 2). An upper value of the $B(E2; 2_2^+ \rightarrow 2_1^+)$ of $15.5 e^2 \text{fm}^4$ can be derived from this ratio, assuming a negligible $M1$ contribution, and taking our experimental value of $B(E2; 2_2^+ \rightarrow 0_1^+) = 10.6 e^2 \text{fm}^4$. Any $M1$ contribution will further reduce this $B(E2)$ value.

There is a large difference between the calculated decay rates of the $2_2^+ \rightarrow 2_1^+$ and $2_2^+ \rightarrow 0_1^+$ that are 412 and $1 e^2 \text{fm}^4$, respectively. This is at variance with experimental values of ≤ 15.5 and 10.6, respectively. The calculated decay rate of the $2_3^+ \rightarrow 2_1^+$ transition ($45 e^2 \text{fm}^4$) is about ten times smaller than that of the $2_2^+ \rightarrow 2_1^+$ transition. Therefore, from its decay properties, it is likely that the 2_3^+ state obtained in the SM calculation corresponds to the 2_2^+ state observed experimentally. The very small quadrupole moment value of the 2_3^+ state indicates its noncollective character. Because of this, predictions of its decay patterns and partial lifetimes are particularly sensitive to mixing with tiny components of the wave function.

The potential-energy surface (PES) displayed in the upper panel of Fig. 9 is obtained from constrained Hartree-Fock calculations in the shell-model basis. At the mean-field level, the nucleus exhibits a spherical minimum in competition with a triaxial shape slightly higher (≈ 1.6 MeV) in energy. The angular-momentum projection already favors the triaxial shape energetically and beyond mean-field mixing around the triaxial projected minimum stabilizes it.

The evolution of the obtained energy spectrum and associated $B(E2)$ values for various approximations is shown in the middle panel of Fig. 9. From left to right, we present the projected triaxial solution (labeled ‘‘Triaxial AMP’’) extracted from the top panel, this solution correlated through a generator coordinate method (GCM) mixing with 17 states, and the final correlated shell-model results. All three approximations reproduce the magnitude of the yrast $2_1^+ \rightarrow 0_{\text{g.s.}}^+$ $E2$ transition. However, the gradual suppression of the $2_2^+ \rightarrow 0_{\text{g.s.}}^+$ transition occurs as correlations are added, although it is difficult to trace back which ones are the most relevant.

The Kumar invariants method [43] has also been applied to the fully correlated shell-model calculations to extract the (β, γ) parameters in ^{70}Zn and their associated variances. As seen from Fig. 9(c), ^{70}Zn exhibits a peculiar regime with a strong triaxial collectivity. This triaxial degree of freedom is in full agreement with what was obtained from the mean-field analysis discussed above.

B. ^{68}Ni

The ^{68}Ni nucleus has been described almost two decades ago as a mixture between a spherical closed shell and a superfluid nucleus [11]. In addition, although not clearly assigned yet, the occurrence of a superdeformed state at low excitation energy has also been recently debated [7]. The PES shown in Fig. 10 reveals the occurrence of three minima (spherical, oblate, and prolate) within a range of 3–4 MeV excitation energy (0.0, 2.26, and 3.53 MeV, respectively). These minima correspond to the first three low-lying observed 0^+ states obtained in the shell-model calculations at 0.0, 1.4, and 2.4 MeV, which are in very good agreement with the 0.0, 1.6, and 2.5 MeV experimental values. This agreement between

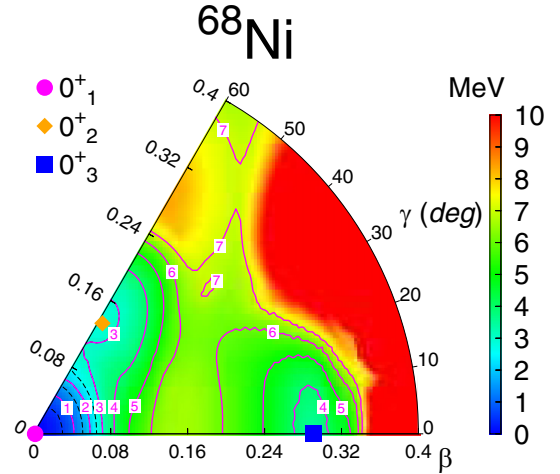


FIG. 10. Potential-energy surface from constrained Hartree-Fock minimization in the shell-model basis for ^{68}Ni . The three predicted 0^+ minima of the PES, spherical, oblate, and prolate, are shown with different symbols.

the full-space diagonalization and the experimental excitation energies is striking as the theoretical spectrum is very sensitive to the details of the residual interaction. Analysis of the obtained wave functions shows configurations which can be interpreted in terms of particle-hole (ph) excitations: the first two 0^+ states appear to be partially mixed, with leading components of $N = 40$ closed shell (CS) and $2p2h$ neutron excitations. As compared with the ground state, the 0_2^+ state has an additional component of proton $1p1h$ with the neutron $2p2h$ coupled to $J = 2$. The 0_3^+ state appears as a $6p6h$ excitation composed of proton $2p2h$ and neutron $4p4h$ components. The Kumar deformation parameters from Ref. [43] agree very well with the present constrained minimization. This strengthens the complementarity between the two descriptions (the deformed Hartree-Fock in the intrinsic frame of ^{68}Ni , the SM in the laboratory).

To generate the $E2$ response from the ^{68}Ni ground state, we computed the $E2$ structure function and show the comparison with the experimental data in Fig. 11. Its $E2$ energy-dependent strength agrees relatively well with the experimental result

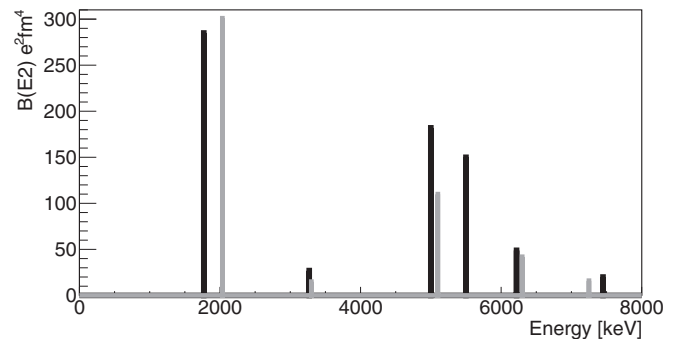


FIG. 11. Comparison between the calculated (black) and experimental (light gray) $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2^+)$ strength distribution in ^{68}Ni , up to $S_n = 7.79$ MeV (see text for explanation).

up to the neutron separation energy value of 7.79 MeV. The largest contribution goes to the 2_1^+ state which has the same structure in terms of ph excitations as the 0_2^+ state. The remaining part of the strength is shared by several levels around 5 MeV, which are characterized by larger proton excitation contributions as well as by larger neutron excitations across the $N = 40$ gap, as compared with the structure of the ground state. These combined features lead to a relatively weak $E2$ strength at high energies excited from the ground state of ^{68}Ni . Experimental values are about two times weaker than calculated.

As discussed above, we think that the remaining observed strength above S_n is more likely due to $E1$ decays coming from the excitation of PDR states. We therefore do not make a comparison with $E2$ calculated strengths in this energy range.

V. CONCLUSIONS

Using the intermediate energy Coulomb-excitation technique, reduced transition probabilities $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ of $1432(124) e^2 \text{fm}^4$ and $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+)$ of $53(7) e^2 \text{fm}^4$ have been determined for ^{70}Zn . The small but nonzero $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_2^+)$ value confirms that ^{70}Zn does not behave like a good vibrational nucleus [1], for which the excitation of the second phonon would have been strictly forbidden in first order. Both the present beyond-mean-field calculations and shell-model diagonalization point to the development of a triaxial shape for the ground-state band of ^{70}Zn . These theoretical descriptions are in agreement with the Zn isotopes systematics recently obtained using beyond-mean-field techniques and the Gogny force in Ref. [44].

As for ^{68}Ni , our adopted $B(E2; 0_{\text{g.s.}}^+ \rightarrow 2_1^+)$ value of $301(38) e^2 \text{fm}^4$ is in very good agreement with the $B(E2)$ values reported in Refs. [11,15] and with the shell-model calculations of the present work. The nuclear contribution to the 2_1^+ state in ^{68}Ni , extracted both from interactions with a Pb target beyond the safe scattering angle and with a C target, has

been found to be very small and comparable to the Coulomb contribution $\beta_N \simeq \beta_C \simeq 0.10$. This confirms the rigidity of ^{68}Ni against both nuclear and Coulomb excitations.

We have identified for the first time a significant $B(E2)$ strength at high energy in $^{68}\text{Ni}_{40}$. It is centered around 5.6 MeV and amounts to about two-thirds that of the 2_1^+ state. This strength can likely be attributed to proton excitations across the $Z = 28$ closed shell, whose experimental contribution is half of that calculated by Langanke *et al.* [14] and by the present shell-model calculations.

A comparison of the $E2$ strength distribution with those of $^{14}_6\text{C}_8$ and $^{34}_{14}\text{Si}_{20}$ nuclei, which also exhibit a combined proton and neutron magic numbers originating from the spin-orbit and harmonic-oscillator potentials, would be interesting. This unfortunately cannot be made as for the two latter nuclei, only the $E2$ strength to the first 2^+ state has been measured so far.

Finally, combining the present experimental results, those of Refs. [7–9] and the present theoretical calculations, it is remarkable to see that four different shapes (spherical, oblate, prolate, and triaxial) are present when considering the low-energy 0^+ states of ^{68}Ni and the ground state of ^{70}Zn . This demonstrates the extreme sensitivity of the nuclear structure in this region to very small changes in degree of freedom (here by increasing slightly the excitation energy in ^{68}Ni or adding two protons to reach ^{70}Zn).

ACKNOWLEDGMENTS

This work was supported by the Romanian National Authority grant for Scientific Research, CNCS - UEFISCDI, Project No. PN-II-RU-TE-2011-3-0051, the ENSAR GA 654002 project, the National Research, Development and Innovation Fund of Hungary, financed under the K18 funding scheme with Project No. K128947 and by the European Regional Development Fund (Contract No. GINOP-2.3.3-15-2016-00034). We are grateful for the technical support provided by the GANIL facility staff. O.S. also wishes to thank D. Weisshaar for his useful comments on the paper.

-
- [1] D. Mucher, G. Gurdal, K.-H. Speidel, G. J. Kumbartzki, N. Benczer-Koller, S. J. Q. Robinson, Y. Y. Sharon, L. Zamick, A. F. Lisetskiy, R. J. Casperson, A. Heinz, B. Krieger, J. Leske, P. Maier-Komor, V. Werner, E. Williams, and R. Winkler, *Phys. Rev. C* **79**, 054310 (2009).
- [2] C. Morse *et al.*, *Phys. Lett. B* **780**, 227 (2018).
- [3] T. Motobayashi *et al.*, *Phys. Lett. B* **346**, 9 (1995).
- [4] W. Rother, A. Dewald, H. Iwasaki, S. M. Lenzi, K. Starosta, D. Bazin, T. Baugher, B. A. Brown, H. L. Crawford, C. Fransen, A. Gade, T. N. Ginter, T. Glasmacher, G. F. Grinyer, M. Hackstein, G. Ilie, J. Jolie, S. McDaniel, D. Miller, P. Petkov *et al.*, *Phys. Rev. Lett.* **106**, 022502 (2011).
- [5] O. Sorlin, *EPJ Web Conf.* **66**, 01016 (2014).
- [6] F. Rotaru, F. Negoita, S. Grevy, J. Mrazek, S. Lukyanov, F. Nowacki, A. Poves, O. Sorlin, C. Borcea, R. Borcea, A. Buta, L. Caceres, S. Calinescu, R. Chevrier, Z. Dombradi, J. M. Daugas, D. Lebhertz, Y. Penionzhkevich, C. Petrone, D. Sohler, M. Stanoiu, and J. C. Thomas, *Phys. Rev. Lett.* **109**, 092503 (2012).
- [7] F. Recchia, C. J. Chiara, R. V. F. Janssens, D. Weisshaar, A. Gade, W. B. Walters, M. Albers, M. Alcorta, V. M. Bader, T. Baugher, D. Bazin, J. S. Berryman, P. F. Bertone, B. A. Brown, C. M. Campbell, M. P. Carpenter, J. Chen, H. L. Crawford, H. M. David, D. T. Doherty *et al.*, *Phys. Rev. C* **88**, 041302(R) (2013).
- [8] S. Suchyta, S. N. Liddick, Y. Tsunoda, T. Otsuka, M. B. Bennett, A. Chemey, M. Honma, N. Larson, C. J. Prokop, S. J. Quinn, N. Shimizu, A. Simon, A. Spyrou, V. Tripathi, Y. Utsuno, and J. M. VonMoss, *Phys. Rev. C* **89**, 021301(R) (2014).
- [9] B. P. Crider *et al.*, *Phys. Lett. B* **763**, 108 (2016).
- [10] R. W. Ibbotson, T. Glasmacher, B. A. Brown, L. Chen, M. J. Chromik, P. D. Cottle, M. Fauerbach, K. W. Kemper, D. J. Morrissey, H. Scheit, and M. Thoennessen, *Phys. Rev. Lett.* **80**, 2081 (1998).
- [11] O. Sorlin, S. Leenhardt, C. Donzaud, J. Duprat, F. Azaiez, F. Nowacki, H. Grawe, Z. Dombradi, F. Amorini, A. Astier,

- D. Baiborodin, M. Belleguic, C. Borcea, C. Bourgeois, D. M. Cullen, Z. Dlouhy, E. Dragulescu, M. Gorska, S. Grévy, D. Guillemaud-Mueller *et al.*, *Phys. Rev. Lett.* **88**, 092501 (2002).
- [12] R. Grzywacz, R. Beraud, C. Borcea, A. Emsallem, M. Glogowski, H. Grawe, D. Guillemaud-Mueller, M. Hjorth-Jensen, M. Houry, M. Lewitowicz, A. C. Mueller, A. Nowak, A. Plochocki, M. Pfitzner, K. Rykaczewski, M. G. Saint-Laurent, J. E. Sauvestre, M. Schaefer, O. Sorlin, J. Szerypo *et al.*, *Phys. Rev. Lett.* **81**, 766 (1998).
- [13] R. Lică *et al.*, *Phys. Rev. C* **100**, 034306 (2019).
- [14] K. Langanke, J. Terasaki, F. Nowacki, D. J. Dean, and W. Nazarewicz, *Phys. Rev. C* **67**, 044314 (2003).
- [15] N. Bree, I. Stefanescu, P. A. Butler, J. Cederkall, T. Davinson, P. Delahaye, J. Eberth, D. Fedorov, V. N. Fedosseev, L. M. Fraile, S. Franchoo, G. Georgiev, K. Gladnishki, M. Huyse, O. Ivanov, J. Iwanicki, J. Jolie, U. Koster, T. Kroll, R. Krucken *et al.*, *Phys. Rev. C* **78**, 047301 (2008).
- [16] B. Pritychenko *et al.*, *At. Data Nucl. Data Tables* **107**, 1 (2016).
- [17] <https://www.nndc.bnl.gov/nudat2>.
- [18] R. Neuhausen *et al.*, *Nucl. Phys. A* **263**, 249 (1976).
- [19] P. H. Stelson and F. K. McGowan, *Nucl. Phys.* **32**, 652 (1962).
- [20] K. P. Singh, D. C. Tayal, and H. S. Hans, *Phys. Rev. C* **58**, 1980 (1998).
- [21] O. Kenn, K. H. Speidel, R. Ernst, S. Schielke, S. Wagner, J. Gerber, P. Maier-Komor, and F. Nowacki, *Phys. Rev. C* **65**, 034308 (2002).
- [22] I. Celikovic, Ph.D. thesis Université de Caen. 2013, <https://tel.archives-ouvertes.fr/tel-00981493>.
- [23] S. Calinescu, L. Caceres, S. Grevy, O. Sorlin, Z. Dombradi, M. Stanoiu, R. Astatatyan, C. Borcea, R. Borcea, M. Bowry, W. Catford, E. Clement, S. Franchoo, R. Garcia, R. Gillibert, I. H. Guerin, I. Kuti, S. Lukyanov, A. Lepailleur, V. Maslov *et al.*, *Phys. Rev. C* **93**, 044333 (2016).
- [24] R. Anne *et al.*, *Nucl. Instrum. Methods Phys. Res., Sect. A* **257**, 215 (1987).
- [25] S. Ottini-Hustache *et al.*, *Nucl. Instrum. Methods Phys. Res., Sect. A* **431**, 476 (1999).
- [26] S. Leenhardt *et al.*, *Eur. Phys. J. A* **14**, 1 (2002).
- [27] C. Bertulani, C. M. Campbell, and T. Glasmacher, *Comput. Phys. Commun.* **152**, 317 (2003).
- [28] J. Raynal, <https://www.oecd-neo.org/science/om200/raynal.pdf>.
- [29] M. Bernas *et al.*, *J. Phys., Lett.* **45**, L-851 (1994).
- [30] S. Franchoo, M. Huyse, K. Kruglov, Y. Kudryavtsev, W. F. Mueller, R. Raabe, I. Reusen, P. Van Duppen, J. Van Roosbroeck, L. Vermeeren, A. Wöhr, K. L. Kratz, B. Pfeiffer, and W. B. Walters, *Phys. Rev. Lett.* **81**, 3100 (1998).
- [31] A. Winther and K. Alder, *Nucl. Phys. A* **319**, 518 (1979).
- [32] W. W. Wilcke *et al.*, *At. Data Nucl. Data Tables* **25**, 391 (1980).
- [33] F. W. N. de Boer *et al.*, *Z. Phys. A* **325**, 457 (1986).
- [34] J. D. Jackson, *Classical Electrodynamics* (Wiley, New York, 1974), p. 620.
- [35] P. Roussel-Chomaz *et al.*, *Phys. Lett. B* **209**, 187 (1988).
- [36] T. Suomijarvi *et al.*, *Nucl. Phys. A* **509**, 369 (1990).
- [37] R. Liguori Neto *et al.*, *Nucl. Phys. A* **560**, 733 (1993).
- [38] M. Zielińska *et al.*, *Eur. Phys. J. A* **52**, 99 (2016).
- [39] D. Goutte, J. B. Bellicard, J. M. Cavedon, B. Frois, M. Huet, P. Leconte, P. X. Ho, S. Platchkov, J. Heisenberg, J. Lichtenstadt, C. N. Papanicolas, and I. Sick, *Phys. Rev. Lett.* **45**, 1618 (1980).
- [40] S. M. Lenzi, F. Nowacki, A. Poves, and K. Sieja, *Phys. Rev. C* **82**, 054301 (2010).
- [41] M. Dufour and A. P. Zuker, *Phys. Rev. C* **54**, 1641 (1996).
- [42] H. L. Crawford, R. M. Clark, P. Fallon, A. O. Macchiavelli, T. Baugher, D. Bazin, C. W. Beausang, J. S. Berryman, D. L. Bleuel, C. M. Campbell, M. Cromaz, G. de Angelis, A. Gade, R. O. Hughes, I. Y. Lee, S. M. Lenzi, F. Nowacki, S. Paschalis, M. Petri, A. Poves *et al.*, *Phys. Rev. Lett.* **110**, 242701 (2013).
- [43] A. Poves, F. Nowacki, and Y. Alhassid, *Phys. Rev. C* **101**, 054307 (2020).
- [44] M. Rocchini *et al.*, *Phys. Rev. C* **103**, 014311 (2021).