EUROPEAN ORGANIZATION FOR PARTICLE PHYSICS

LEPEWWG/96-02
ALEPH 96-107 PHYSIC 96-98
DELPHI 96-121 PHYS 631
L3 Note 1975
OPAL Technical Note TN 399
SLD Physics Note 52
30 July 1996

A Combination of Preliminary LEP and SLD Electroweak Measurements and Constraints on the Standard Model

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Internal Note

Prepared from Contributions of the LEP and SLD Experiments to the 1996 summer conferences.

Abstract

This note presents a combination of published and preliminary electroweak results from the four LEP collaborations and the SLD collaboration which were prepared for the 1996 summer conferences. Averages of the results concerning electroweak physics are presented. They are derived from the measurements of hadronic and leptonic cross sections, the leptonic forward-backward asymmetries, the τ polarization asymmetries, the $b\overline{b}$ and $c\overline{c}$ partial widths and forward-backward asymmetries and the $q\overline{q}$ charge asymmetry. The main differences with respect to the results presented at the 95 summer conferences are the updated measurements of the mass and width of the Z, including the results of the 1995 LEP energy scan. In addition, there are updated measurements of the $b\overline{b}$ and $c\overline{c}$ partial widths, the forward-backward asymmetry for b and c quarks, the τ polarization asymmetries and of the $q\overline{q}$ charge asymmetry. The results are compared to precise electroweak measurements from other experiments. The parameters of the Standard Model are evaluated, first using the combined LEP electroweak measurements, and then using the full set of precise electroweak results.

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1 Introduction

In a recent note [1] the four LEP experiments have presented parameters derived from the Z resonance using published and preliminary results based on data recorded until the end of 1994. The preliminary results were contributions to the 1995 International Europhysics Conference on High-Energy Physics, 27/7 to 2/8/1995, Brussels, Belgium and the 17th International Symposium on Lepton-Photon Interactions, 10-15/8/1995, Beijing, China.

Since then several additional preliminary analyses have become available. In addition, there are preliminary results from the 1995 Z scan. Results from the the high energy (130–140 GeV) run at the end of 1995 are presented elsewhere [2].

The data consist of the hadronic and leptonic cross sections, the leptonic forward-backward asymmetries, the τ polarization asymmetries, the $\mathrm{b\overline{b}}$ and $\mathrm{c\overline{c}}$ partial widths and forward-backward asymmetries and the $\mathrm{q\overline{q}}$ charge asymmetry. Many technical aspects of their combination have already been described in References 3, 4 and references therein. It should be stressed, however, that several measurements included in the current combination are still preliminary.

This note is organized in the following manner. In Section 2 the results on the Z line shape and leptonic forward-backward asymmetries are presented, while Section 3 contains the measurements of the τ polarization. Section 4 describes the parameters resulting from heavy flavor analyses, Section 5 new results for the inclusive hadronic charge asymmetry. Section 6 is devoted to the interpretation of the results. In Sections 6.1, 6.2 and 6.3 several LEP electroweak measurements are combined to determine the effective neutral current coupling constants and to give a value of the effective electroweak mixing angle. We also quote values for these parameters when the left-right and left-right forward-backward asymmetries from SLD [5] are included. The determination of the number of light neutrino species is discussed in Section 6.4. In Section 6.5 the LEP data and also data from SLD [5–7], from neutrino interactions [8–10] and from measurements of the mass of the W boson [11–14] and the top quark [15–17] are used to constrain the parameters of the Standard Model.

2 Z Lineshape and Lepton Forward-Backward Asymmetries

The results presented here are based on the data taken during the energy scans in 1990 and 1991 with centre-of-mass energies, $\sqrt{s}$, in a range $|\sqrt{s} - m_{\mathrm{Z}}| < 3$ GeV, on the data collected at the Z peak in 1992 and on preliminary analyses of the energy scans in 1993 and the peak running in 1994. All the experiments include preliminary results on the lepton forward-backward asymmetries using data from the 1995 energy scan where each experiment collected approximately 40 pb^{-1} at and near the Z peak. In addition, three experiments include cross-section data from the 1995 energy scan.

The total statistics and the systematic errors on the individual analyses of the four LEP collaborations are given in Tables 1 and 2. Details of the individual analyses can be found in References 18–21.

The measurement of the LEP energies, and the associated uncertainties, are important in the determination of the mass and width of the Z. In the previous note [1] the treatment of the LEP energies used was that described in Reference 22. Since then, the 1995 scan has been performed. The instrumentation of LEP was improved by employing NMR devices in two LEP dipole magnets in octants 4 and 8. Furthermore, in four physics fills resonant depolarization measurements were made at both the beginning and end of fills. There were also two machine development fills in which resonant depolarization measurements were made over periods of several hours. From all of these

measurements it is concluded that there is a significant rise in the LEP energy during the course of a fill. Such a rise term was included in the 1993 analysis, but the magnitude of the rise observed in 1995 was considerably larger than that estimated for the 1993 scan.

Preliminary results for the 1995 LEP energies are now available [23]. In the light of the 1995 energy rise studies, the energy determinations for the 1993 scan, and the 1994 peak data (for which the same model was used as for the 1993 scan) have been revised. Studies are still in progress, and the 1993 and 1994 energy values and errors are preliminary.

The LEP energy error assignments are given in the form of an error matrix with the six entries, corresponding to the peak and two off peak-points for each of the years 1993 and 1995. In addition, an error was evaluated for the 1994 peak data.

Because of the correlations in the LEP energy error matrix between the 1993 and 1995 scans, the estimation of those components of the LEP energy errors which are common between experiments is more difficult than for previous combinations of LEP data. This is further complicated by the fact that, at the present time, only three LEP experiments have cross-section data available from the 1995 scan. The procedure adopted to estimate the LEP energy errors for the combinations of the 5 and 9 parameter fits described below was to perform a global fit to the hadronic cross-section data for all experiments for the years 1993, 1994 and 1995, taking into account both the LEP errors and correlations as well as those on the cross-sections. From these studies the common LEP errors relevant for the 5 and 9 parameter fits were determined. This procedure is exact apart from the neglect of data from earlier years and from the leptonic channels.

For the averaging of results the LEP experiments provide a standard set of 9 parameters describing the information contained in hadronic and leptonic cross sections and leptonic forward-backward asymmetries [3, 24]. These parameters have initial-state QED corrections, as well as t-channel and s/t-interference contributions in the case of the e^+e^- final state, removed. They are convenient for fitting and averaging since they have small correlations. The parameters are:

- The mass and total width of the Z boson, where the definition is based on the Breit-Wigner denominator $(s - m_Z^2 + is\Gamma_Z/m_Z)$ [25].
- The hadronic pole cross section of Z exchange:
$$\sigma_h^0 \equiv \frac{12\pi}{m_Z^2}\frac{\Gamma_{ee}\Gamma_{had}}{\Gamma_Z^2}\,.$$
Here Γ_{ee} and Γ_{had} are the partial widths of the Z for decays into electrons and hadrons.
- The ratios:
$$R_e \equiv \Gamma_{had}/\Gamma_{ee} \quad R_\mu \equiv \Gamma_{had}/\Gamma_{\mu\mu} \quad R_\tau \equiv \Gamma_{had}/\Gamma_{\tau\tau}. \tag{1}$$
Here $\Gamma_{\mu\mu}$ and $\Gamma_{\tau\tau}$ are the partial widths of the Z for the decays $Z \to \mu^+\mu^-$ and $Z \to \tau^+\tau^-$. Even under the assumption of lepton universality a small difference of 0.2% is expected between the values for R_e and R_μ, and the value for R_τ, owing to mass corrections to $\Gamma_{\tau\tau}$.
- The pole asymmetries, $A_{FB}^{0,e}$, $A_{FB}^{0,\mu}$ and $A_{FB}^{0,\tau}$ for the processes $e^+e^- \to e^+e^-$, $e^+e^- \to \mu^+\mu^-$ and $e^+e^- \to \tau^+\tau^-$. In terms of the effective vector and axial-vector neutral current couplings of fermions, g_{Vf} and g_{Af}, the pole asymmetries are expressed as:[1]
$$A_{FB}^{0,f} \equiv \frac{3}{4}\mathcal{A}_e\mathcal{A}_f \tag{2}$$

[1] In the definition of $A_{FB}^{0,f}$, effects coming from γ exchange, γ/Z interference, as well as real and imaginary parts of the photon vacuum polarization are not included. They are accounted for explicitly in the fitting formulae used by the experiments, and are fixed to their Standard Model values.

with:

$$\mathcal{A}_{\mathrm{f}} \equiv \frac{2 g_{V\mathrm{f}} g_{A\mathrm{f}}}{g_{V\mathrm{f}}^2 + g_{A\mathrm{f}}^2} \,. \tag{3}$$

This set of 9 parameters does not describe the Z production and decay completely, because it does not include the interference of the Z exchange with the γ exchange. This contribution is investigated in a separate note [26]. For the results presented in this Section, the γ-exchange contributions and the γZ interference terms are fixed to their Standard Model values.[2]

The four sets of 9 parameters provided by the LEP experiments are presented in Table 3. The covariance matrix of these parameters is as described in our previous paper [24]. It is constructed from the covariance matrices of the individual LEP experiments and common systematic errors. These common errors arise from the theoretical uncertainty in the luminosity normalization affecting the hadronic pole cross section, $\Delta\sigma_{\mathrm{h}}^0/\sigma_{\mathrm{h}}^0 = 0.11\%$, from the uncertainty of the LEP centre-of-mass energy spread of about 1 MeV [23], resulting in $\Delta\Gamma_{\mathrm{Z}} \approx 0.2$ MeV, and from the uncertainty in the LEP energy calibration. The latter uncertainty causes errors of $\Delta m_{\mathrm{Z}} \approx 1.5$ MeV, $\Delta\Gamma_{\mathrm{Z}} \approx 1.7$ MeV, and $\Delta A_{\mathrm{FB}}^{0,\ell} \approx 0.0005$ for each lepton species ($\ell = \mathrm{e}, \mu, \tau$). It should be noted that the error from the LEP energy beam spread has been reduced considerably from previous determinations. Full correlation between $A_{\mathrm{FB}}^{0,\mu}$ and $A_{\mathrm{FB}}^{0,\tau}$ and full anti-correlation between $A_{\mathrm{FB}}^{0,\mathrm{e}}$ and $A_{\mathrm{FB}}^{0,\mu}$ or $A_{\mathrm{FB}}^{0,\tau}$ is used. This anti-correlation for $A_{\mathrm{FB}}^{0,\mathrm{e}}$ is an approximation of the effect of the t-channel contribution for a typical LEP experimental acceptance for the $\mathrm{e^+e^-}$ final state. The combined parameter set and its correlation matrix are given in Tables 4 and 5.

If lepton universality is assumed, the set of 9 parameters given above is reduced to a set of 5 parameters. R_ℓ is defined as $R_\ell \equiv \Gamma_{\mathrm{had}}/\Gamma_{\ell\ell}$, where $\Gamma_{\ell\ell}$ refers to the partial Z width for the decay into a pair of massless charged leptons.

The data of each of the four LEP experiments are consistent with lepton universality (the difference in χ^2 over the difference in d.o.f. with and without the assumption of lepton universality are 6/4, 4/4, 4/4 and 5/4 for ALEPH, DELPHI, L3 and OPAL, respectively). Based on this assumption Table 6 provides the five parameters m_{Z}, Γ_{Z}, σ_{h}^0, R_ℓ and $A_{\mathrm{FB}}^{0,\ell}$ for the individual LEP experiments. The four experiments all use the above definition of $\Gamma_{\ell\ell}$. Tables 7 and 8 provide these five parameters and the corresponding correlation matrix for the combined result of the four LEP experiments. Figure 1 shows, for each lepton species and for the combination assuming lepton universality, the resulting 68% probability contours in the R_ℓ-$A_{\mathrm{FB}}^{0,\ell}$ plane.

For completeness the partial decay widths of the Z boson are listed in Table 9.

[2] If instead the γZ interference terms are entirely determined from LEP cross-section data (including the 130-140 GeV data), the total error on the LEP average of m_{Z} increases from 2.2 MeV to 4.8 MeV [26].

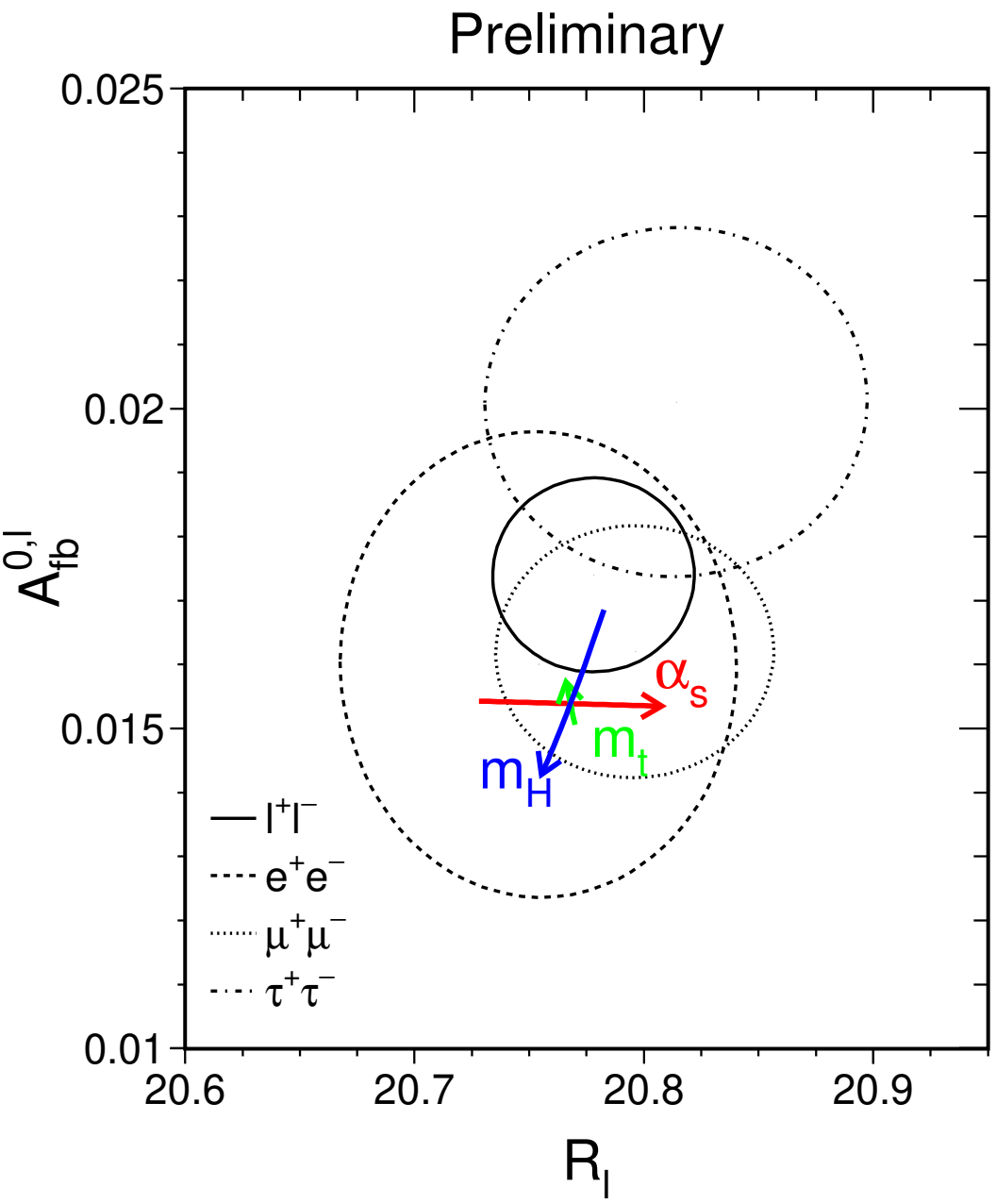

Figure 1: Contours of 68% probability in the R_ℓ-$A_{\rm FB}^{0,\ell}$ plane. The Standard Model prediction for $m_{\rm Z} = 91.1863$ GeV, $m_{\rm t} = 175$ GeV, $m_{\rm H} = 300$ GeV, and $\alpha_s = 0.123$ is also shown. The lines with arrows correspond to the variation of the Standard Model prediction when $m_{\rm t}$, $m_{\rm H}$ or $\alpha_s(m_{\rm Z}^2)$ are varied in the intervals $m_{\rm t} = 175 \pm 6$ GeV, $m_{\rm H} = 300^{+700}_{-240}$ GeV, and $\alpha_s(m_{\rm Z}^2) = 0.123 \pm 0.006$, respectively. The arrows point in the direction of increasing values of $m_{\rm t}$, $m_{\rm H}$ and α_s.

		ALEPH	DELPHI	L3	OPAL	LEP
$q\bar{q}$	'90-'91	451	357	416	454	1678
	'92	680	697	678	733	2788
	'93 prel.	640	677	646	646	2609
	'94 prel.	1654	1241	1307	1524	5726
	'95 prel.	739	584	311	–	1634
	total	4164	3556	3358	3357	14435
$\ell^+\ell^-$	'90-'91	55	36	40	58	189
	'92	82	70	58	88	298
	'93 prel.	78	74	64	82	298
	'94 prel.	190	129	127	184	630
	'95 prel.	80	67	28	42	217
	total	485	376	317	454	1632

Table 1: The LEP statistics in units of 10^3 events used for the analysis of the Z line shape and lepton forward-backward asymmetries.

	ALEPH			DELPHI			L3			OPAL		
	'93 prel.	'94 prel.	'95 prel.	'93 prel.	'94 prel.	'95 prel.	'93 prel.	'94 prel.	'95 prel.	'93 prel.	'94 prel.	'95 prel.
$\mathcal{L}^{\mathrm{exp.}(b)}$	0.087%	0.073%	0.092%	0.24%	0.09%	0.09%	0.10%	0.078%	0.128%	0.076%	0.079%	[a]
σ_{had}	0.073%	0.073%	0.076%	0.10%	0.10%	0.10%	0.052%	0.051%	0.10%	0.15%	0.16%	[a]
σ_{e}	0.50%	0.48%	0.47%	0.44%	0.50%	0.60%	0.30%	0.23%	1.0%	0.23%	0.24%	[a]
σ_{μ}	0.25%	0.26%	0.25%	0.28%	0.30%	[a]	0.31%	0.31%	1.0%	0.16%	0.15%	[a]
σ_{τ}	0.34%	0.32%	0.39%	0.80%	0.60%	[a]	0.67%	0.65%	0.60%	0.43%	0.46%	[a]
$A_{\mathrm{FB}}^{\mathrm{e}}$	0.0031	0.0031	0.0028	0.0025	0.0022	0.0025	0.003	0.003	0.01	0.0016	0.0016	0.002
A_{FB}^{μ}	0.0005	0.0005	0.0005	0.0010	0.0015	0.0015	0.0008	0.0008	0.005	0.001	0.001	0.001
A_{FB}^{τ}	0.0009	0.0007	0.0009	0.0020	0.0020	0.0020	0.003	0.003	0.003	0.002	0.002	0.002

Table 2: The experimental systematic errors for the analysis of the Z line shape and lepton forward-backward asymmetries at the Z peak. The errors quoted do not include the common uncertainty due to the LEP energy calibration. The treatment of correlations between the errors for different years is described in References 18–21.

[a] No preliminary result quoted yet.

[b] In addition, there is a theoretical error for the calculation of the small angle Bhabha cross section of 0.11%, which has been treated as common to all experiments. The ALEPH theoretical error is 0.16% and the DELPHI theoretical error for 1993 0.17%.

	ALEPH	DELPHI	L3	OPAL
$m_{\rm Z}$(GeV)	91.1873±0.0030	91.1859±0.0028	91.1883±0.0029	91.1824±0.0039
$\Gamma_{\rm Z}$(GeV)	2.4950±0.0047	2.4896±0.0042	2.4996±0.0043	2.4956±0.0053
$\sigma^0_{\rm h}$(nb)	41.576±0.083	41.566±0.079	41.411±0.074	41.53±0.09
R_e	20.64±0.09	20.93±0.14	20.78±0.11	20.82±0.14
R_μ	20.88±0.07	20.70±0.09	20.84±0.10	20.79±0.07
R_τ	20.78±0.08	20.78±0.15	20.75±0.14	20.99±0.12
$A^{0,\rm e}_{\rm FB}$	0.0187±0.0039	0.0179±0.0051	0.0148±0.0063	0.0104±0.0052
$A^{0,\mu}_{\rm FB}$	0.0179±0.0025	0.0153±0.0026	0.0176±0.0035	0.0146±0.0025
$A^{0,\tau}_{\rm FB}$	0.0196±0.0028	0.0223±0.0039	0.0233±0.0049	0.0178±0.0034
χ^2/d.o.f.	195/217	174/157	142/159	$12/6^{(a)}$

Table 3: Line shape and asymmetry parameters from 9-parameter fits to the data of the four LEP experiments.
$^{(a)}$This parameter set has been obtained from a parameter transformation applied to the 15 parameters of the OPAL fit [21], which treats the γZ interference terms for leptons as additional free parameters. The extra parameters for the γZ interference terms have been fixed to their Standard Model values in the transformation. The χ^2/d.o.f. for the 15-parameter fit to the data is 87/132.

Parameter	Average Value
$m_{\rm Z}$(GeV)	91.1863±0.0020
$\Gamma_{\rm Z}$(GeV)	2.4946±0.0027
$\sigma^0_{\rm h}$(nb)	41.508±0.056
$R_{\rm e}$	20.754±0.057
R_μ	20.796±0.040
R_τ	20.814±0.055
$A^{0,\rm e}_{\rm FB}$	0.0160±0.0024
$A^{0,\mu}_{\rm FB}$	0.0162±0.0013
$A^{0,\tau}_{\rm FB}$	0.0201±0.0018

Table 4: Average line shape and asymmetry parameters from the data of the four LEP experiments given in Table 3, without the assumption of lepton universality. The χ^2/d.o.f. of the average is 22/27.

	$m_{\rm Z}$	$\Gamma_{\rm Z}$	$\sigma^0_{\rm h}$	$R_{\rm e}$	R_μ	R_τ	$A^{0,\rm e}_{\rm FB}$	$A^{0,\mu}_{\rm FB}$	$A^{0,\tau}_{\rm FB}$
$m_{\rm Z}$	1.00	0.09	−0.01	0.01	−0.02	−0.01	0.02	0.06	0.04
$\Gamma_{\rm Z}$	0.09	1.00	−0.14	0.00	−0.01	0.00	0.00	0.00	0.00
$\sigma^0_{\rm h}$	−0.01	−0.14	1.00	0.07	0.12	0.08	0.01	0.00	0.00
$R_{\rm e}$	0.01	0.00	0.07	1.00	0.05	0.04	−0.01	0.00	0.00
R_μ	−0.02	−0.01	0.12	0.05	1.00	0.05	−0.01	0.01	0.00
R_τ	−0.01	0.00	0.08	0.04	0.05	1.00	0.00	0.00	0.01
$A^{0,\rm e}_{\rm FB}$	0.02	0.00	0.01	−0.01	−0.01	0.00	1.00	0.01	0.01
$A^{0,\mu}_{\rm FB}$	0.06	0.00	0.00	0.00	0.01	0.00	0.01	1.00	0.01
$A^{0,\tau}_{\rm FB}$	0.04	0.00	0.00	0.00	0.00	0.01	0.01	0.01	1.00

Table 5: The correlation matrix for the set of parameters given in Table 4.

	ALEPH	DELPHI	L3	OPAL
$m_{\rm Z}$(GeV)	91.1874±0.0030	91.1859±0.0028	91.1883±0.0029	91.1822±0.0039
$\Gamma_{\rm Z}$(GeV)	2.4948±0.0047	2.4896±0.0042	2.4996±0.0043	2.4955±0.0053
$\sigma^0_{\rm h}$(nb)	41.578±0.083	41.566±0.079	41.411±0.074	41.53±0.09
R_ℓ	20.766±0.049	20.754±0.068	20.788±0.066	20.83±0.06
$A^{0,\ell}_{\rm FB}$	0.0187±0.0017	0.0175±0.0020	0.0187±0.0026	0.0150±0.0019
χ^2/d.o.f.	200/221	178/161	144/163	$15/10^{(a)}$

Table 6: Line shape and asymmetry parameters from 5-parameter fits to the data of the four LEP experiments, assuming lepton universality. R_ℓ is defined as $R_\ell \equiv \Gamma_{\rm had}/\Gamma_{\ell\ell}$, where $\Gamma_{\ell\ell}$ refers to the partial Z width for the decay into a pair of massless charged leptons.
$^{(a)}$This parameter set has been obtained by a parameter transformation applied to the 15 parameters of the OPAL fit.

Parameter	Average Value
$m_{\rm Z}$(GeV)	91.1863±0.0020
$\Gamma_{\rm Z}$(GeV)	2.4946±0.0027
$\sigma^0_{\rm h}$(nb)	41.508±0.056
R_ℓ	20.778±0.029
$A^{0,\ell}_{\rm FB}$	0.0174±0.0010

Table 7: Average line shape and asymmetry parameters from the results of the four LEP experiments given in Table 6, assuming lepton universality. R_ℓ is defined as $R_\ell \equiv \Gamma_{\rm had}/\Gamma_{\ell\ell}$, where $\Gamma_{\ell\ell}$ refers to the partial Z width for the decay into a pair of massless charged leptons. The χ^2/d.o.f. of the average is 38/31.

	$m_{\rm Z}$	$\Gamma_{\rm Z}$	$\sigma^0_{\rm h}$	R_ℓ	$A^{0,\ell}_{\rm FB}$
$m_{\rm Z}$	1.00	0.09	−0.01	−0.01	0.08
$\Gamma_{\rm Z}$	0.09	1.00	−0.14	−0.01	0.00
$\sigma^0_{\rm h}$	−0.01	−0.14	1.00	0.15	0.01
R_ℓ	−0.01	−0.01	0.15	1.00	0.01
$A^{0,\ell}_{\rm FB}$	0.08	0.00	0.01	0.01	1.00

Table 8: The correlation matrix for the set of parameters given in Table 7.

Without Lepton Universality:		
$\Gamma_{\rm ee}$	(MeV)	83.96±0.15
$\Gamma_{\mu\mu}$	(MeV)	83.79±0.22
$\Gamma_{\tau\tau}$	(MeV)	83.72±0.26

With Lepton Universality:		
$\Gamma_{\ell\ell}$	(MeV)	83.91±0.11
$\Gamma_{\rm had}$	(MeV)	1743.6±2.5
$\Gamma_{\rm inv}$	(MeV)	499.5±2.0

Table 9: Partial decay widths of the Z boson, derived from the results of the 9-parameter (Tables 4 and 5) and the 5-parameter fit (Tables 7 and 8). In the case of lepton universality, $\Gamma_{\ell\ell}$ refers to the partial Z width for the decay into a pair of massless charged leptons.

3 The τ Polarization

The τ polarization, $\mathcal{P}_\tau$, is determined by an indirect measurement of the longitudinal polarization of τ pairs produced in Z decays. It is defined as:

$$\mathcal{P}_\tau \equiv \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}, \tag{4}$$

where σ_R and σ_L are the τ-pair cross sections for the production of a right-handed and left-handed τ^-, respectively. The angular distribution of $\mathcal{P}_\tau$ as a function of the angle θ between the e^- and the τ^-, for $\sqrt{s} = m_Z$, is given by:

$$\mathcal{P}_\tau(\cos\theta) = -\frac{\mathcal{A}_\tau(1+\cos^2\theta) + 2\mathcal{A}_e\cos\theta}{1+\cos^2\theta + 2\mathcal{A}_\tau\mathcal{A}_e\cos\theta}, \tag{5}$$

with $\mathcal{A}_e$ and $\mathcal{A}_\tau$ as defined in Equation (3). Equation (5) neglects corrections for the effects of γ exchange, γZ interference and electromagnetic radiative corrections for initial- and final-state radiation. These effects are, however, taken into account in the experimental analyses. In particular, these corrections account for the $\sqrt{s}$ dependence of the tau polarization, $\mathcal{P}_\tau(\cos\theta)$, which is important since the off-peak data are included in the event samples for all experiments. When averaged over all production angles $\mathcal{P}_\tau$ is a measurement of $\mathcal{A}_\tau$. As a function of $\cos\theta$, $\mathcal{P}_\tau(\cos\theta)$ provides nearly independent determinations of both $\mathcal{A}_\tau$ and $\mathcal{A}_e$, thus allowing a test of the universality of the couplings of the Z to e and τ.

Each experiment makes separate $\mathcal{P}_\tau$ measurements using the five τ decay modes $e\nu\overline{\nu}$, $\mu\nu\overline{\nu}$, $\pi\nu$, $\rho\nu$ and $a_1\nu$ [27–30]. The $\rho\nu$ and $\pi\nu$ are the most sensitive channels, contributing weights of about 40% each in the average. In addition, DELPHI has used an inclusive hadronic analysis. The combination is made of the results from each experiment already averaged over the τ decay modes.

A discussion of the effects of possible common systematic errors between the experiments can be found in [3]. Further study of the uncertainties from the effects of radiative corrections for the ρ decay mode of the τ is desirable.

3.1 Results

Tables 10 and 11 show the most recent results for $\mathcal{A}_\tau$ and $\mathcal{A}_e$ obtained by the four experiments [27–30] and their combination. No common systematics are included in these averages. The statistical correlation between the extracted values of $\mathcal{A}_\tau$ and $\mathcal{A}_e$ is small ($\leq 5\%$), and is neglected.

The average values for $\mathcal{A}_\tau$ and $\mathcal{A}_e$:

$$\mathcal{A}_\tau = 0.1401 \pm 0.0067 \tag{6}$$
$$\mathcal{A}_e = 0.1382 \pm 0.0076, \tag{7}$$

are compatible, as is expected from lepton universality. Assuming $e - \tau$ universality, the values for $\mathcal{A}_\tau$ and $\mathcal{A}_e$ can be combined. This combination is performed neglecting any possible common systematic error between $\mathcal{A}_\tau$ and $\mathcal{A}_e$ within a given experiment. Such errors are estimated to be small, but warrant further study. The combined result of $\mathcal{A}_\tau$ and $\mathcal{A}_e$ gives:

$$\mathcal{A}_\ell = 0.1393 \pm 0.0050. \tag{8}$$

ALEPH	('90 - '92), final	$0.136 \pm 0.012 \pm 0.009$
DELPHI	('90 - '94), prel.	$0.138 \pm 0.009 \pm 0.008$
L3	('90 - '94), prel.	$0.152 \pm 0.010 \pm 0.009$
OPAL	('90 - '94), final	$0.134 \pm 0.009 \pm 0.010$
LEP Average		0.1401 ± 0.0067

Table 10: LEP results for $\mathcal{A}_\tau$. The χ^2/d.o.f. for the average is 1.1/3. The first error is statistical and the second systematic. In the LEP average, statistical and systematic error are combined in quadrature. The systematic component of the error, obtained by combining the individual systematic errors, is ± 0.0045.

ALEPH	('90 - '92), final	$0.129 \pm 0.016 \pm 0.005$
DELPHI	('90 - '94), prel.	$0.140 \pm 0.013 \pm 0.003$
L3	('90 - '94), prel.	$0.156 \pm 0.016 \pm 0.005$
OPAL	('90 - '94), prel.	$0.129 \pm 0.014 \pm 0.005$
LEP Average		0.1382 ± 0.0076

Table 11: LEP results for $\mathcal{A}_e$. The χ^2/d.o.f. for the average is 1.8/3. The first error is statistical and the second systematic. In the LEP average, statistical and systematic error are combined in quadrature. The systematic component of the error, obtained by combining the individual systematic errors, is ± 0.0021.

4 Results from b and c Quarks

The relevant quantities measured in the heavy quark sector at LEP are:

- The ratios of the b and c quark partial widths of the Z to its total hadronic partial width: $R_{\mathrm{b}} \equiv \Gamma_{\mathrm{b\bar{b}}}/\Gamma_{\mathrm{had}}$ and $R_{\mathrm{c}} \equiv \Gamma_{\mathrm{c\bar{c}}}/\Gamma_{\mathrm{had}}$.
- The forward-backward asymmetries, $A_{\mathrm{FB}}^{\mathrm{b\bar{b}}}$ and $A_{\mathrm{FB}}^{\mathrm{c\bar{c}}}$.
- The semileptonic branching ratios, BR(b $\rightarrow \ell$) and BR(b $\rightarrow$ c $\rightarrow \bar{\ell}$), and the average $\mathrm{B^0\overline{B^0}}$ mixing parameter, $\overline{\chi}$. These are often determined at the same time as the widths or asymmetries in multi-parameter fits to lepton tag samples. They are included in the combination procedure to take into account their correlations with the other parameters measured in the same fit.
- The probability that a c-quark fragments into a $\mathrm{D^{*+}}$, a $\mathrm{D^+}$, a $\mathrm{D_s}$ or a charmed baryon[3]. These quantities are determined now with good accuracy by the LEP experiments. The interpretation of the $\mathrm{D^*}$ rate in terms of R_c and the determination of the charm background in the lifetime tag R_{b} measurements can now be done without assumptions on the energy dependence of the D-meson production rates.

There are several motivations for the averaging procedure. Several analyses measure more than one parameter simultaneously, for example the lepton fits. Some of the measurements of electroweak parameters depend explicitly on the values of other parameters, for example R_{b} depends on R_{c}. The common tagging and analysis techniques lead to common sources of systematic uncertainty, in particular for the double-tag measurements of R_{b}. The starting point for the combination is to ensure that all the analyses use a common set of assumptions for input parameters which give rise to systematic uncertainties. The set of parameters used up to the conferences in winter 1996 together with a full description of the averaging procedure have been published in [4]. These parameters have been updated recently [31] to take into account more recent measurements as well as different needs of the new analyses. The correlations and interdependences of the input measurements are then taken into account in a χ^2 minimization which results in the combined electroweak parameters and their correlation matrix. The only significant correlation between any of the resulting electroweak parameters turns out to be between R_{b} and R_{c}. This is discussed at the end of this section.

In a first fit the asymmetry measurements on peak, above peak and below peak were combined at each centre-of-mass energy. The results of this fit, including the SLD results, are given in the appendix. The dependence of the average asymmetries on centre-of-mass energy agrees with the prediction of the Standard Model. To derive the pole asymmetries, $A_{\mathrm{FB}}^{0,\mathrm{q}}$, from the quark asymmetries, all the off-peak asymmetry measurements were then corrected to the peak energy before combining. Only results from this second fit are quoted here. There are therefore 11 parameters in total to be determined in the combination procedure: the two partial widths, two asymmetries, two semileptonic branching ratios, the average mixing parameter and the probabilities for c quark to fragment into a $\mathrm{D^{*+}}$, a $\mathrm{D^+}$ a $\mathrm{D_s}$, or a charmed baryon.

Recently the SLD collaboration has presented precise measurements of R_{b} [6] and of the left-right forward-backward asymmetry for b and c quarks [7]. Since the precision and the dominant sources

[3] Actually the product P(c $\rightarrow \mathrm{D^{*+}}$) $\times$ BR($\mathrm{D^{*+}} \rightarrow \pi^+\mathrm{D^0}$) is fitted since this quantity is needed and measured by the LEP experiments.
The probability that a c-quark fragments into a $\mathrm{D^0}$ is calculated from the constraint that the probabilities for the weakly decaying charmed hadrons add up to one

of systematic uncertainty are similar at LEP and SLD it is useful to produce combined LEP+SLD averages. The left-right forward-backward asymmetries are, in contrast to the unpolarized forward-backward asymmetries, only sensitive to the final state couplings ($\mathcal{A}_{\rm b}$ and $\mathcal{A}_{\rm c}$). They are treated in the averaging procedure as physically independent quantities. However the methods to measure the polarized and unpolarized asymmetries are very similar, so $\mathcal{A}_{\rm b}$ and $\mathcal{A}_{\rm c}$ are included in the averaging procedure in order to estimate the correlation between the SLD and the LEP asymmetries, resulting in a 13-parameter fit.

4.1 Summary of measurements and averaging procedure

The measurements of $R_{\rm b}$ and $R_{\rm c}$ fall into two categories. In the first, called a single-tag measurement, a method to select b or c events is devised, and the number of tagged events is counted. This number must then be corrected for backgrounds from other flavors and for the tagging efficiency to calculate the true fraction of hadronic Z decays of that flavor. The dominant systematic errors come from understanding the branching ratios and detection efficiencies which give the overall tagging efficiency. For the second technique, called a double-tag measurement, the event is divided into two hemispheres. For an $R_{\rm b}$ measurement, writing the number of tagged single hemispheres as N_t, the number of events with both hemispheres tagged as N_{tt}, then for a total of $N_{\rm had}$ hadronic Z decays:

$$\begin{aligned} \frac{N_t}{2N_{\rm had}} &= \varepsilon_{\rm b}R_{\rm b} + \varepsilon_{\rm c}R_{\rm c} + \varepsilon_{\rm uds}(1 - R_{\rm b} - R_{\rm c}), \\ \frac{N_{tt}}{N_{\rm had}} &= \mathcal{C}_{\rm b}\varepsilon_{\rm b}^2 R_{\rm b} + \varepsilon_{\rm c}^2 R_{\rm c} + \varepsilon_{\rm uds}^2(1 - R_{\rm b} - R_{\rm c}), \end{aligned}$$

where $\varepsilon_{\rm b}$, $\varepsilon_{\rm c}$ and $\varepsilon_{\rm uds}$ are the tagging efficiencies per hemisphere for b, c and light-quark events, and $\mathcal{C}_{\rm b} \neq 1$ accounts for the fact that the tagging efficiencies between the hemispheres may be correlated. In practice, $\varepsilon_{\rm b} \gg \varepsilon_{\rm c} \gg \varepsilon_{\rm uds}$, $\mathcal{C}_{\rm b} \approx 1$, and the correlations for the other flavors are neglected. These equations can be solved to give $R_{\rm b}$ and $\varepsilon_{\rm b}$, which neglecting the c and uds backgrounds and the correlations are approximately given by:

$$\begin{aligned} \varepsilon_{\rm b} &\approx 2N_{tt}/N_t, \\ R_{\rm b} &\approx N_t^2/(4N_{tt}N_{\rm had}). \end{aligned}$$

The double-tagging method has the advantage that the tagging efficiency is derived directly from the data, reducing the systematic error of the measurement. The residual background of other flavors in the sample, and the evaluation of the correlation between the tagging efficiencies in the two hemispheres of the event are the main sources of systematic uncertainty in such an analysis.

The measurements included are:

- Lepton fits from all four LEP experiments [32–37]. Each fit measures several parameters chosen from: $R_{\rm b}$, $R_{\rm c}$, $A_{\rm FB}^{\rm b\bar{b}}$, $A_{\rm FB}^{\rm c\bar{c}}$, BR(b $\to \ell$) and BR(b $\to$ c $\to \bar{\ell}$), and $\overline{\chi}$. The measured parameters are correlated. These analyses use hadronic events with one or more leptons in the final state. The semileptonic branching ratios can therefore be measured by a double-tagging technique if the lepton identification efficiency is known. The dominant sources of systematic error for the lepton fits arise from the lepton identification, from other semileptonic branching ratios and from the modeling of the semileptonic decay. In addition to the single/double lepton fits ALEPH has measured BR(b $\to \ell$) and BR(b $\to$ c $\to \bar{\ell}$) in a lifetime tagged sample and $R_{\rm c}$ from low energy electrons assuming BR(c $\to \ell$).
- Event-shape tag for $R_{\rm b}$ from L3 (single tag) [38].

- Lifetime (and lepton) double tag measurements for R_b from ALEPH, DELPHI, L3 , OPAL and SLD [6, 39–42]. These are the most precise determinations of R_b, and dominate the combined result. The features of the double-tag technique were discussed above. These results depend explicitly on the assumed value of R_c. The ALEPH measurement has been improved for this conference by the use of five different tags, two of them being sensitive to uds and c quarks. By measuring all double tag combinations only the uds and c efficiency for the purest b-tag and the hemisphere correlations need to be taken from simulation. In the ALEPH and SLD measurements the charm rejection has been enhanced by using the invariant mass information.

- Measurements of $A_{FB}^{b\bar{b}}$ based on lifetime tagged events with a hemisphere charge measurement from ALEPH , DELPHI and OPAL where the mean b-hemisphere charge is derived from the charge distributions themselves [35, 43, 44]. These contribute roughly the same weight to the combined result as the lepton fits. Note that the quoted ALEPH result is a Standard Model fit to various charge properties both on- and off-peak. This has been converted to a determination assuming independent b partial widths by the authors, and is treated as an effective measurement of the peak $A_{FB}^{b\bar{b}}$ in the averages.

- Analyses with $D/D^{*\pm}$ mesons to measure R_c from ALEPH, DELPHI and OPAL [33,45,46]. All measurements are constructed in a way that no assumptions on the energy dependence of charm fragmentation are necessary. The available measurements can be divided into four groups:

  - inclusive/exclusive double tag (ALEPH, DELPHI, OPAL): In a first step $D^{*\pm}$ mesons are reconstructed in several decay channels and their production rate is measured, which depends on the product $R_c \times P(c \to D^{*+}) \times BR(D^{*+} \to \pi^+ D^0)$. This sample of clean $c\bar{c}$ (and $b\bar{b}$) events is then used to measure $P(c \to D^{*+}) \times BR(D^{*+} \to \pi^+ D^0)$ using a slow pion tag in the opposite hemisphere. The ALEPH measurement is doing the unfolding of R_c internally in the analysis so that no explicit $P(c \to D^{*+}) \times BR(D^{*+} \to \pi^+ D^0)$ is available. However the principle of the method is identical to the one of DELPHI and OPAL.
  - inclusive single/double tag (DELPHI): This measurement measures the single and double tag rate using a slow pion tag. It takes advantage of the much higher efficiency of the inclusive slow pion tag compared to the exclusive reconstruction, however the high background has to be taken into account in the systematics.
  - exclusive double tag (ALEPH): This analysis uses exclusively reconstructed D^{*+}, D^0 and D^+ mesons in different decay channels. It has lower statistics but much better purity than the inclusive analyses.
  - Reconstruction of all weakly decaying D states (DELPHI, OPAL): These analyses make the assumption that the production rates of D^0, D^+, D_s and Λ_c saturate the fragmentation of $c\bar{c}$, since every charm quark ends up in a ground state hadron before the weak decay. This is a single tag measurement, but is insensitive to uncertainties in the charm hadronisation, relying only on the decay branching ratios of the charm hadrons.

- Analyses with a $D^{*\pm}$ meson to measure $A_{FB}^{c\bar{c}}$ from ALEPH [47] or $A_{FB}^{c\bar{c}}$ and $A_{FB}^{b\bar{b}}$ from DELPHI and OPAL [35, 48].

- Measurements of $\mathcal{A}_b$ and $\mathcal{A}_c$ from SLD [7]. These results use lepton, kaon, $D^{*\pm}$ and lifetime plus hemisphere charge tags, with similar sources of systematic error as the LEP asymmetry measurements.

These measurements are presented by the LEP and SLD collaborations in a consistent manner for the purpose of combination [4]. The tables prepared by the experiments include a detailed breakdown of the systematic error of each measurement and its dependence on other electroweak parameters.

Where necessary, the experiments apply small corrections to their results in order to use agreed values and ranges for the input parameters used to calculate systematic errors. The measurements, corrected where necessary, are summarized in the Appendix in Tables 24-33, where the statistical and systematic errors are quoted separately. The correlated systematic entries are from sources shared with one or more other results in the table and are derived from the full breakdown of common systematic uncertainties. The uncorrelated systematic entries are from the remaining sources.

The χ^2 minimization procedure used to derive the values of the heavy-flavor electroweak parameters was published in [4]. The explicit dependences of each measurement on the other parameters is taken into account, for example the dependence of the value of R_b on the assumed value of R_c as described below. The full statistical and systematic covariance matrix for all the measurements is calculated. The correlation matrices relating several measurements made in the same analyses are used, in particular for the multi-parameter lepton fits. Additional correlations arising from common sources of systematic uncertainty are estimated from a detailed breakdown of systematic errors.

Since c-quark events form the main background in the R_b analyses, the value of R_b depends on the value of R_c. If R_b and R_c are measured in the same analysis, this is reflected in the correlation matrix for the results. Otherwise the normalization of the charm contribution is not fixed by the data, and the measurement of R_b depends on the assumed value of R_c, which may be written as:

$$R_b = R_b^{meas} + a(R_c)\frac{(R_c - R_c^{used})}{R_c}. \tag{9}$$

In this expression, R_b^{meas} is the result of the analysis assuming a value of $R_c = R_c^{used}$. The values of R_c^{used} and the coefficients $a(R_c)$ are given in Table 24 where appropriate. The dependences of all other measurements on other electroweak parameters are treated in the same way, with coefficients $a(x)$ describing the dependence on parameter x.

4.2 Treatment of the LEP Asymmetry Measurements

For the 11-parameter fit described above, the peak and off-peak asymmetries were corrected to $\sqrt{s} = 91.26$ GeV, using the predicted dependence from ZFITTER [49]. The slope of the asymmetry around m_Z depends only on the axial coupling and the charge of the initial and final state fermions and is thus independent of the value of the asymmetry itself.

After calculating the overall averages, the quark pole asymmetries, $A_{FB}^{0,q}$, were derived by applying the corrections described below. The input measured asymmetries are all corrected to full acceptance. To relate the pole asymmetries to these numbers a few corrections that are summarized in Table 12 have to be applied:

- Energy shift to $\sqrt{s} = m_Z$ and QED corrections: The corrections due to the energy shift and initial state radiation have been calculated using ZFITTER [49].
- γ exchange, γZ interference: These very small corrections have again been calculated using ZFITTER.

The input measured asymmetries use the thrust axis as an estimate of the quark direction before gluon radiation. The lifetime/jet-charge measurements of asymmetries take into account QCD effects as an inherent part of the analysis, but the measured asymmetries for the analyses using a lepton or D^* tag need to be corrected for the effects of QCD on the event thrust axis direction. The QCD

corrections, using the thrust axis to define θ for the event, have been calculated to first order, including mass corrections [50]. Second order QCD correction have been calculated [51] using the true b quark direction and without mass correction. For example from these calculations a correction of $A_{\rm FB}^{QCD} = A_{\rm FB}^{no\,QCD}(0.969 \pm 0.0025)$ has be estimated [31] for the b asymmetries. In recent studies [31] it appeared that the experimental event selection and signal extraction introduce an important bias to the QCD corrections. In the case of the DELPHI measurement using lepton, the QCD correction to the b (c) quark asymmetry has been estimated to be reduced by 50% (70%) from its theoretical expectation. For the moment not all LEP analysis have estimated the effect of this bias on their measured asymmetry. For this reason, whenever an evaluation of the effect was not available, the measurements were corrected using the DELPHI estimation with inflated error to take into account possible variation in the experimental bias. In this case the correction factor used were 0.984 ± 0.008 for the b and 0.99 ± 0.01 for the c asymmetries.

Source	$\delta A_{\rm FB}^{\rm b}$	$\delta A_{\rm FB}^{\rm c}$
$\sqrt{s} = m_{\rm Z}$	–0.0013	–0.0034
QED corrections	+0.0041	+0.0104
γ, γZ	–0.0003	–0.0008
Total	+0.0025	+0.0062

Table 12: Corrections to be applied to the quark asymmetries. The corrections are to be understood as $A_{\rm FB}^0 = A_{\rm FB}^{meas} + \sum_i (\delta A_{\rm FB})_i$

4.3 Results

4.3.1 Results of the 11-parameter fit to LEP data

Using the full averaging procedure gives the following combined results for the electroweak parameters:

$$\begin{aligned} R_{\rm b}^0 &= 0.2179 \pm 0.0012 \\ R_{\rm c}^0 &= 0.1715 \pm 0.0056 \\ A_{\rm FB}^{0,\rm b} &= 0.0979 \pm 0.0023 \\ A_{\rm FB}^{0,\rm c} &= 0.0733 \pm 0.0049\,, \end{aligned} \tag{10}$$

where all corrections to the asymmetries and partial widths have been already applied. The χ^2/dof is $50/(81-11)$. The corresponding correlation matrix is given in Table 13. If $R_{\rm c}^0$ is fixed to its Standard-Model prediction of 0.172, then the value of $R_{\rm b}^0$ is:

$$R_{\rm b}^0(R_{\rm c}^0 = 0.172) = 0.2179 \pm 0.0011\,.$$

4.3.2 Results of the 13-parameter fit to LEP and SLD data

Including the SLD results on $R_{\rm b}$, $\mathcal{A}_{\rm b}$ and $\mathcal{A}_{\rm c}$ into the fit the following results are obtained:

	R_b^0	R_c^0	$A_{FB}^{0,b}$	$A_{FB}^{0,c}$
R_b^0	1.00	–0.23	0.00	0.01
R_c^0	–0.23	1.00	0.04	-0.07
$A_{FB}^{0,b}$	0.00	0.04	1.00	0.10
$A_{FB}^{0,c}$	0.01	-0.07	0.10	1.00

Table 13: The correlation matrix for the electroweak parameters from the 11-parameter fit

$$\begin{aligned} R_b^0 &= 0.2178 \pm 0.0011 \\ R_c^0 &= 0.1715 \pm 0.0056 \\ A_{FB}^{0,b} &= 0.0979 \pm 0.0023 \\ A_{FB}^{0,c} &= 0.0735 \pm 0.0048 \\ \mathcal{A}_b &= 0.863 \pm 0.049 \\ \mathcal{A}_c &= 0.625 \pm 0.084\,, \end{aligned} \tag{11}$$

with a χ^2/dof of $51/(87-13)$. The corresponding correlation matrix is given in Table 14. In deriving these results the parameters $\mathcal{A}_b$ and $\mathcal{A}_c$ have been treated as independent of the forward-backward asymmetries $A_{FB}^{b\bar{b}}(\mathrm{pk})$ and $A_{FB}^{c\bar{c}}(\mathrm{pk})$. The fit results of $\mathcal{A}_b$ and $\mathcal{A}_c$ take into account the LEP and SLD average values of R_b and R_c, whereas the input numbers assumed the Standard Model values. This accounts for the fit results being lower than the input measurements.

	R_b^0	R_c^0	$A_{FB}^{0,b}$	$A_{FB}^{0,c}$	$\mathcal{A}_b$	$\mathcal{A}_c$
R_b^0	1.00	–0.23	0.00	0.00	–0.03	0.01
R_c^0	–0.23	1.00	0.04	–0.06	0.05	–0.07
$A_{FB}^{0,b}$	0.00	0.04	1.00	0.10	0.04	0.02
$A_{FB}^{0,c}$	0.00	–0.06	0.10	1.00	0.01	0.10
$\mathcal{A}_b$	–0.03	0.05	0.04	0.01	1.00	0.12
$\mathcal{A}_c$	0.01	–0.07	0.02	0.10	0.12	1.00

Table 14: The correlation matrix for the electroweak parameters from the 13-parameter fit

If R_c^0 is fixed to its Standard-Model prediction of 0.172, then the value of R_b^0 is:

$$R_b^0(R_c^0 = 0.172) \;=\; 0.2178 \pm 0.0011\,.$$

The result of the full fit to the LEP/SLC results including the off peak asymmetries can be found in the appendix.

4.4 Comments on the changes in R_b and R_c

Compared to the results available last summer some changes occurred in the central values of R_b and R_c.

For R_c a significant change is coming from the fact that the low energy constraint on $P(c \to D^{*+})$ is no longer used. The low energy number $P(c \to D^{*+}) \times BR(D^{*+} \to \pi^+ D^0) \;=\; 0.178 \pm 0.013$ is

completely consistent with the LEP measurement of 0.163 ± 0.007 however due to the large correlation of -60% R_c is pulled to 0.1637 when $P(c \to D^{*+}) \times BR(D^{*+} \to \pi^+ D^0)$ is fixed to 0.178. In addition some analyses have been improved since last year and a lot of new data have been added. For the combined R_c result the largest errors are statistical (0.0037) and systematics internal to the experiments (0.0029). From the error sources common to the experiments the only relevant sources are the branching ratios $BR(c \to \ell)$ (0.0022) and $BR(D_s \to \phi\pi)$ (0.0011) for which the model independent CLEO measurement [52] is used. The branching ratio $BR(D^0 \to K\pi)$ contributes only 0.0002 to the total error on R_c.

The change in R_b is basically due to the inclusion of new data. The new input parameters have the tendency to lower R_b as well mainly due to the inclusion of the measured gluon splitting rate $g \to c\bar{c}$ but the effect is only of the order 0.0003. The new data presented this summer are:

- ALEPH has presented a new analysis of their 1992 to 1995 statistics with a new very pure tag. This supercedes their published 1992 analysis.
- DELPHI has updated their number with 1994 data,
- L3 has presented for the first time a lifetime tag measurement,

Since some of the R_b measurements depend on the charmed hadron production fractions which are also fitted parameters now it is no longer possible simply to fit the R_b measurements alone. To test the compatibility of the R_b measurements of the LEP/SLD experiments the following procedure has been adopted: in a first step the LEP charm measurements have been combined to obtain a best estimate of the charm production fractions with R_c fixed to its Standard Model value. In a second step the precise single parameter R_b measurements have been fitted together with the result of the first fit. The result of this fit with R_c fixed to 0.172 was $R_b = 0.2174 \pm 0.0012$ with $\chi^2/NDF = 5.1/5$ showing agreement between the different experiments. In addition the new ALEPH and DELPHI results are consistent with their older published numbers.

For the combined R_b with R_c fixed the dominant error sources are again statistics (0.00067) and internal effects (0.00053). The dominant common effects are the inclusive branching ratio $D \to K^0 X$ (0.00022), The charged D decay multiplicity (0.00029), QCD related effects to the hemisphere correlations (0.00031) and the gluon splitting to b and c quark pairs (0.00044).

4.5 Comments on the correlation between R_b and R_c

As noted before, the only large correlation between any of the combined electroweak parameters in Equations 10 and 11 is between R_b and R_c. It is useful to examine how the correlations and interdependencies of the input measurements lead to this correlation between the combined parameters.

The measurements of R_c are almost independent of an assumed value for R_b. For the lepton fit analyses, there is a very small net correlation between the two partial widths. There is a negative statistical correlation from events moving between the b and c contributions to the lepton sample. The overall systematic correlation is positive, dominated by effects such as the subtraction of hadronic background to the lepton sample, or the lepton identification efficiency which affect the b and c contributions in the same way. The net correlation from statistical and systematic effects is small.

The D and D^* meson measurements of R_c are constructed to be independent of an assumed value for R_b, because they make a statistical separation of the b and c contributions based on other event

properties. The overall normalization of the b contribution is therefore fixed by the data. These results depend on the features of b-hadron decay, but not on an assumed value of R_{b}.

The precise double-tag measurements of R_{b} dominate the combined R_{b} result. The value of R_{b} from these measurements depend explicitly on the value of R_{c} assumed, and this dependence combined with the current error on R_{c} is almost entirely responsible for the size of the correlation coefficient between R_{b} and R_{c}.

5 The Hadronic Charge Asymmetry $\langle Q_{FB} \rangle$

The LEP experiments ALEPH [53–55], DELPHI [56, 57], and OPAL [58, 59] have provided measurements of the hadronic charge asymmetry based on the mean difference in jet charges measured in the forward and backward event hemispheres, $\langle Q_{FB} \rangle$. DELPHI has also provided a related measurement of the total charge asymmetry by making a charge assignment on an event-by-event basis and performing a likelihood fit [56]. The experimental values quoted for the average forward-backward charge difference $\langle Q_{FB} \rangle$, cannot be directly compared as some of them include detector dependent effects such as acceptances and efficiencies. We therefore use the effective electroweak mixing angle, $\sin^2\theta^{lept}_{eff}$, defined in Section 6.3, as a means of combining the experimental results summarized in Table 15.

Experiment		$\sin^2\theta^{lept}_{eff}$
ALEPH	90-94, final	$0.2322 \pm 0.0008 \pm 0.0011$
DELPHI	91-94, prel.	$0.2311 \pm 0.0010 \pm 0.0014$
OPAL	91-94, prel.	$0.2326 \pm 0.0012 \pm 0.0013$
Average		0.2320 ± 0.0010

Table 15: Summary of the determination of $\sin^2\theta^{lept}_{eff}$ from inclusive hadronic charge asymmetries at LEP. For each experiment, the first error is statistical and the second systematic. The latter is dominated by fragmentation and decay modeling uncertainties.

The dominant source of systematic error arises from the modeling of the charge flow in the fragmentation process for each flavor. All experiments measure the required charge properties for $Z \to b\bar{b}$ events from the data. ALEPH also determine the charm charge properties from the data. The fragmentation model implemented in the JETSET Monte-Carlo program [60] is used by all experiments as reference; the one of the HERWIG Monte-Carlo program [61] is used for comparison. The JETSET fragmentation parameters are varied to estimate the systematic errors. The central values chosen by the experiments for these parameters are, however, not the same. The degree of correlation between the fragmentation uncertainties for the different experiments requires further investigation. We treat the smaller of the two fragmentation errors in any pair of results as common to both. The present average of $\sin^2\theta^{lept}_{eff}$ from $\langle Q_{FB} \rangle$ and its associated error are not very sensitive to the treatment of common uncertainties.

The ambiguities due to QCD corrections may cause changes in the derived value of $\sin^2\theta^{lept}_{eff}$. These are, however, well below the fragmentation uncertainties and experimental errors. The effect of fully correlating the estimated systematic uncertainties from this source between the experiments on the average and its error is negligible.

There is also some correlation between these results and those for $A^{b\bar{b}}_{FB}$ using jet charges. The dominant source of correlation is again through uncertainties in the fragmentation and decay models used. The typical correlation between the derived values of $\sin^2\theta^{lept}_{eff}$ between the $\langle Q_{FB} \rangle$ and the $A^{b\bar{b}}_{FB}$ jet charge measurement has been estimated to be between 20% and 25%. This leads to only a small change in the relative weights for the $A^{b\bar{b}}_{FB}$ and $\langle Q_{FB} \rangle$ results when averaging their $\sin^2\theta^{lept}_{eff}$ values (Section 6.3). Furthermore, the jet charge method contributes at most half of the weight of the $A^{b\bar{b}}_{FB}$ measurement. Thus, the correlation between $\langle Q_{FB} \rangle$ and $A^{b\bar{b}}_{FB}$ from jet charge will have little impact on the overall Standard Model fit, and is neglected at present.

6 Interpretation of Results

6.1 The Coupling Parameters $\mathcal{A}_f$

The coupling parameters $\mathcal{A}_f$ are defined in terms of the effective vector and axial-vector neutral current couplings of fermions (Equation (3)). The LEP measurements of the forward-backward asymmetries of charged leptons (Section 2) and b and c quarks (Section 4) determine the products $A_{FB}^{0,f} = \frac{3}{4}\mathcal{A}_e\mathcal{A}_f$ (Equation (2)). The LEP measurements of the τ polarization (Section 3), $\mathcal{P}_\tau(\cos\theta)$, determine $\mathcal{A}_\tau$ and $\mathcal{A}_e$ separately (Equation (5)). The SLD collaboration measures the left-right asymmetry, A_{LR} [62,63], which determines the same quantity, $\mathcal{A}_e$, as the τ polarization, with minimal model dependence. Both measurements have small systematic errors. The SLD measurements of the left-right forward-backward asymmetries for b and c quarks [7] are direct determinations of $\mathcal{A}_b$ and $\mathcal{A}_c$.

Table 16 shows the results on the leptonic coupling parameter $\mathcal{A}_\ell$ and their combination assuming lepton universality. Table 17 shows the results on the quark coupling parameters $\mathcal{A}_b$ and $\mathcal{A}_c$ derived from LEP measurements alone (Equation 10) and from the combination of LEP and SLD measurements (Equation 11).

	$\mathcal{A}_\ell$	Cumulative Average	χ^2/d.o.f.
$A_{FB}^{0,\ell}$	0.1523 ± 0.0044		
$\mathcal{P}_\tau(\cos\theta)$	0.1393 ± 0.0050	0.1466 ± 0.0033	3.8/1
A_{LR} (SLD)	0.1542 ± 0.0037	0.1500 ± 0.0025	6.1/2

Table 16: Comparison of the determinations of the leptonic coupling parameter $\mathcal{A}_\ell$ assuming lepton universality. The second column lists the $\mathcal{A}_\ell$ values derived from the quantities listed in the first column. The third column contains the cumulative averages of these $\mathcal{A}_\ell$ results. The averages are derived assuming no correlations between the measurements. The χ^2 per degree of freedom for the cumulative averages is given in the last column.

LEP	using $\mathcal{A}_\ell = 0.1466 \pm 0.0033$
$\mathcal{A}_b$	0.890 ± 0.029
$\mathcal{A}_c$	0.667 ± 0.047
LEP+SLD	**using $\mathcal{A}_\ell = 0.1500 \pm 0.0025$**
$\mathcal{A}_b$	0.867 ± 0.022
$\mathcal{A}_c$	0.646 ± 0.040

Table 17: Determinations of the quark coupling parameters $\mathcal{A}_b$ and $\mathcal{A}_c$ from LEP data alone (using the LEP average for $\mathcal{A}_\ell$) and from LEP+SLD data (using the LEP+SLD average for $\mathcal{A}_\ell$) assuming lepton universality.

6.2 The Effective Vector and Axial-Vector Coupling Constants

The partial widths of the Z into leptons and the lepton forward-backward asymmetries (Section 2), the τ polarization and the τ polarization asymmetry (Section 3) can be combined to determine the effective vector and axial-vector couplings for e, μ and τ. The asymmetries (Equations (2) and (5)) determine the ratio $g_{V\ell}/g_{A\ell}$ (Equation (3)), while the sum of the squares of the couplings is derived from the leptonic partial widths:

$$\Gamma_{\ell\ell} = \frac{G_{\rm F} m_{\rm Z}^3}{6\pi\sqrt{2}}(g_{V\ell}^2 + g_{A\ell}^2)(1+\delta_\ell^{QED}), \tag{12}$$

where $\delta_\ell^{QED} = 3q_\ell^2\alpha(m_{\rm Z}^2)/(4\pi)$ accounts for final state photonic corrections. Corrections due to lepton masses, neglected in Equation 12, are taken into account for the results presented below.

The averaged results for the effective lepton couplings are given in Table 18. Figure 2 shows the 68% probability contours in the $g_{A\ell}$-$g_{V\ell}$ plane. The signs of $g_{A\ell}$ and $g_{V\ell}$ are based on the convention $g_{A{\rm e}} < 0$. With this convention the signs of the couplings of all charged leptons follow from LEP data alone. For comparison, the $g_{V\ell}$-$g_{A\ell}$ relation following from the measurement of $A_{\rm LR}$ from SLD [62,63] is indicated as a band in the $g_{A\ell}$-$g_{V\ell}$-plane of Figure 2. It is consistent with the LEP data. The information on the leptonic couplings from LEP can therefore be combined with the $A_{\rm LR}$ measurement of SLD. The results for this combination are given in the right column of Table 18. The measured ratios of the e, μ and τ couplings provide a test of lepton universality and are also given in Table 18.

	Without Lepton Universality:	
	LEP	LEP+SLD
$g_{V{\rm e}}$	-0.0368 ± 0.0015	-0.03828 ± 0.00079
$g_{V\mu}$	-0.0372 ± 0.0034	-0.0358 ± 0.0030
$g_{V\tau}$	-0.0369 ± 0.0016	-0.0367 ± 0.0016
$g_{A{\rm e}}$	-0.50130 ± 0.00046	-0.50119 ± 0.00045
$g_{A\mu}$	-0.50076 ± 0.00069	-0.50086 ± 0.00068
$g_{A\tau}$	-0.50116 ± 0.00079	-0.50117 ± 0.00079
	Ratios of couplings:	
	LEP	LEP+SLD
$g_{V\mu}/g_{V{\rm e}}$	1.01 ± 0.11	0.935 ± 0.085
$g_{V\tau}/g_{V{\rm e}}$	1.001 ± 0.062	0.959 ± 0.046
$g_{A\mu}/g_{A{\rm e}}$	0.9989 ± 0.0018	0.9993 ± 0.0017
$g_{A\tau}/g_{A{\rm e}}$	0.9997 ± 0.0019	1.0000 ± 0.0019
	With Lepton Universality:	
	LEP	LEP+SLD
$g_{V\ell}$	-0.03688 ± 0.00085	-0.03776 ± 0.00062
$g_{A\ell}$	-0.50115 ± 0.00034	-0.50108 ± 0.00034
g_ν	$+0.5009 \pm 0.0010$	$+0.5009 \pm 0.0010$

Table 18: Results for the effective vector and axial-vector couplings derived from the combined LEP data without and with the assumption of lepton universality. For the right column the SLD measurement of $A_{\rm LR}$ is also included.

The neutrino couplings to the Z can be derived from the measured value of its invisible width, $\Gamma_{\rm inv}$, attributing it exclusively to the decay into three identical neutrino generations ($\Gamma_{\rm inv} = 3\Gamma_{\nu\nu}$) and assuming $g_{A\nu} \equiv g_{V\nu} \equiv g_\nu$. The relative sign of g_ν is chosen to be in agreement with neutrino scattering data [64], resulting in $g_\nu = +0.5009 \pm 0.0010$.

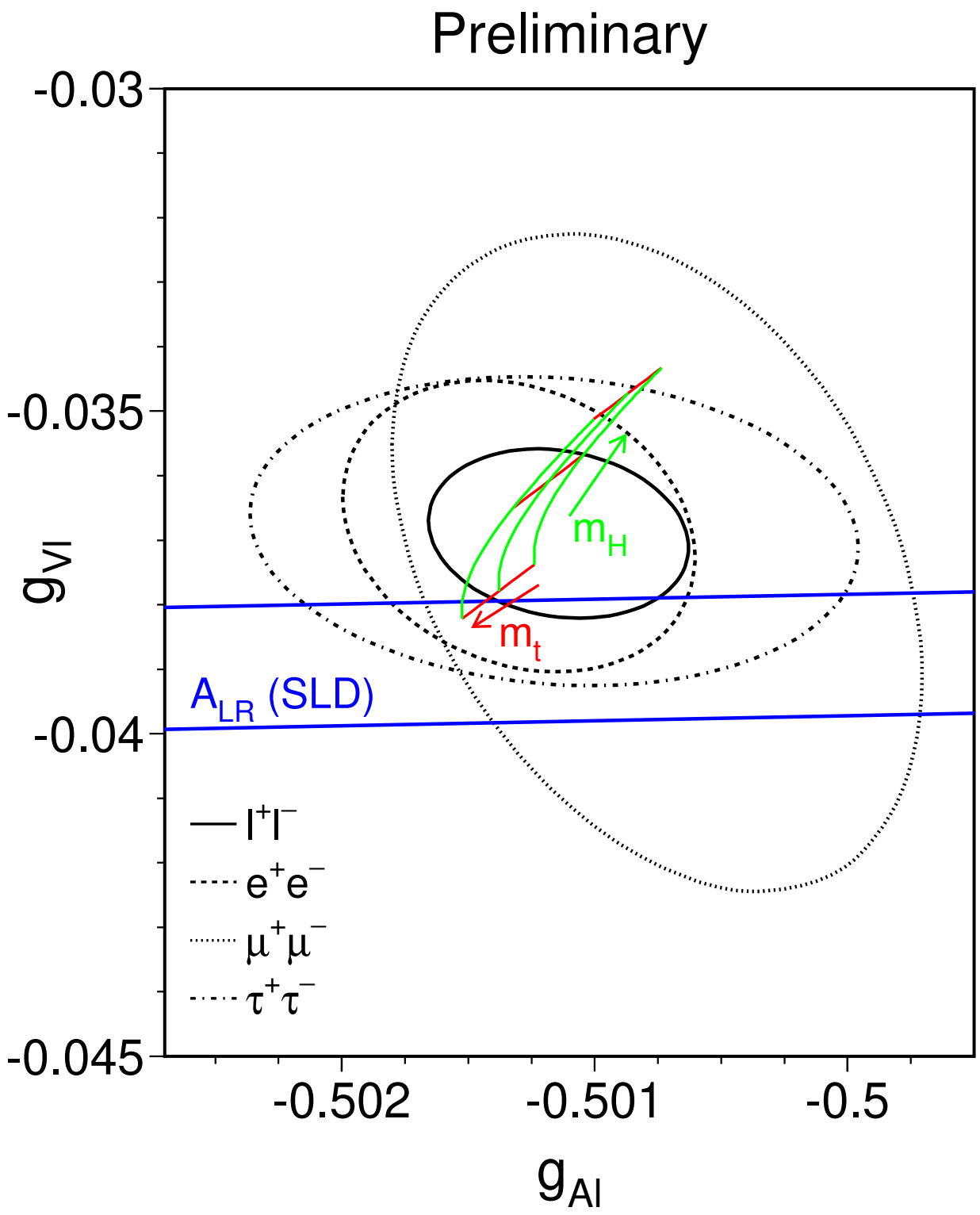

Figure 2: Contours of 68% probability in the $g_{V\ell}$-$g_{A\ell}$ plane from LEP measurements. The solid contour results from a fit assuming lepton universality. Also shown is the one standard deviation band resulting from the $A_{\rm LR}$ measurement of SLD. The grid corresponds to the Standard Model prediction for $m_{\rm t} = 175 \pm 9$ GeV and $m_{\rm H} = 300^{+700}_{-240}$ GeV. The arrows point in the direction of increasing values of $m_{\rm t}$ and $m_{\rm H}$.

6.3 The Effective Electroweak Mixing Angle $\sin^2\theta^{\rm lept}_{\rm eff}$

The asymmetry measurements from LEP can be combined into a single observable, the effective electroweak mixing angle, $\sin^2\theta^{\rm lept}_{\rm eff}$, defined as:

$$\sin^2\theta^{\rm lept}_{\rm eff} \equiv \frac{1}{4}(1 - g_{V\ell}/g_{A\ell})\,, \tag{13}$$

without making any strong model-specific assumptions.

For a combined average of $\sin^2\theta^{\rm lept}_{\rm eff}$ from $A^{0,\ell}_{\rm FB}$, $\mathcal{A}_\tau$ and $\mathcal{A}_{\rm e}$ only the assumption of lepton universality, already inherent in the definition of $\sin^2\theta^{\rm lept}_{\rm eff}$, is needed. In practice no further assumption is involved if the quark forward-backward asymmetries, $A^{\rm b\bar{b}}_{\rm FB}$ and $A^{\rm c\bar{c}}_{\rm FB}$, are included in this average, as these asymmetries have a reduced sensitivity to corrections particular to the hadronic vertex. The results of these determinations of $\sin^2\theta^{\rm lept}_{\rm eff}$ and their combination are shown in Table 19. Also the comparison with the measurement of the left-right asymmetry, $A_{\rm LR}$, from SLD [62,63] is given. Compared to the results presented in our previous note [1], the χ^2 of the average of all determinations has increased by 5. The most significant change in central value is for $A^{\rm b\bar{b}}_{\rm FB}$, although it is consistent within errors of the previous result. The errors on most results have decreased considerably.

	$\sin^2\theta^{\rm lept}_{\rm eff}$	Average by Group of Observations	Cumulative Average	χ^2/d.o.f.
$A^{0,\ell}_{\rm FB}$	0.23085 ± 0.00056			
$\mathcal{A}_\tau$	0.23240 ± 0.00085			
$\mathcal{A}_{\rm e}$	0.23264 ± 0.00096	0.23157 ± 0.00042	0.23157 ± 0.00042	3.9/2
$A^{0,\rm b}_{\rm FB}$	0.23246 ± 0.00041			
$A^{0,\rm c}_{\rm FB}$	0.23155 ± 0.00112	0.23236 ± 0.00038	0.23200 ± 0.00028	6.3/4
$\langle Q_{\rm FB}\rangle$	0.2320 ± 0.0010	0.2320 ± 0.0010	0.23200 ± 0.00027	6.3/5
$A_{\rm LR}$ (SLD)	0.23061 ± 0.00047	0.23061 ± 0.00047	0.23165 ± 0.00024	12.8/6

Table 19: Comparison of several determinations of $\sin^2\theta^{\rm lept}_{\rm eff}$ from asymmetries. Averages are obtained as weighted averages assuming no correlations. The second column lists the $\sin^2\theta^{\rm lept}_{\rm eff}$ values derived from the quantities listed in the first column. The third column contains the averages of these numbers by groups of observations, where the groups are separated by the horizontal lines. The last column shows the cumulative averages. The χ^2 per degree of freedom for the cumulative averages is also given.

6.4 Number of Neutrino Species

An important aspect of our measurement concerns the information related to Z decays into invisible channels. Using the results of Tables 7 and 8, the ratio of the Z decay width into invisible particles and the leptonic decay width is determined:

$$\Gamma_{\rm inv}/\Gamma_{\ell\ell} \;=\; 5.952 \pm 0.023\,.$$

The Standard Model value for the ratio of the partial widths to neutrinos and charged leptons is:

$$(\Gamma_{\nu\nu}/\Gamma_{\ell\ell})_{\rm SM} \;=\; 1.991 \pm 0.001\,.$$

The central value is evaluated for $m_{\rm Z} = 91.1863$ GeV, $m_{\rm t} = 175$ GeV, $m_{\rm H} = 300$ GeV and the error quoted accounts for a variation of $m_{\rm t}$ in the range $m_{\rm t} = 175 \pm 6$ GeV and a variation of $m_{\rm H}$ in the range 60 GeV $\leq m_{\rm H} \leq$ 1000 GeV.

The number of light neutrino species is given by the ratio of the two expressions listed above:

$$N_\nu \;\; = \;\; 2.989 \pm 0.012 \, .$$

6.5 Constraints on the Standard Model

The precise electroweak measurements performed at LEP can be used to check the validity of the Standard Model and, within its framework, to infer valuable information about its fundamental parameters. The accuracy of the measurements makes them sensitive to the top-quark mass, $m_{\rm t}$, and to the mass of the Higgs boson, $m_{\rm H}$, through loop corrections. The leading $m_{\rm t}$ dependence is quadratic and allows a determination of $m_{\rm t}$. The main dependence on $m_{\rm H}$ is logarithmic and therefore, with the present data accuracy, the constraints on $m_{\rm H}$ are still weak.

The various LEP measurements are summarized in Table 20a and are presented in Figures 3, 4, and 5, together with their Standard Model prediction as a function of $m_{\rm t}$. The bands in the Standard Model predictions reflect the linear sum of the expected variations in each quantity due to a change of the strong coupling constant $\alpha_s(m_{\rm Z}^2) = 0.123 \pm 0.006$ [65] and $m_{\rm H}$ in the interval 60 GeV $\leq m_{\rm H} \leq$ 1000 GeV for $m_{\rm Z} = 91.1863$ GeV.

Detailed studies of the theoretical uncertainties in the Standard Model predictions due to missing higher-order electroweak corrections and their interplay with QCD corrections are carried out in the working group on 'Precision calculations for the Z resonance' [66]. Theoretical uncertainties are evaluated by comparing different but, within our present knowledge, equivalent treatments of aspects such as resummation techniques, momentum transfer scales for vertex corrections and factorization schemes. The impact of these intrinsic theoretical uncertainties on $m_{\rm t}$ and $\alpha_s(m_{\rm Z}^2)$ has been estimated by repeating the Standard Model fits in this Section using several combinations of options which were implemented in the electroweak libraries used [67] for the study performed in [66]. As a result the maximal variations of the central values of the fitted parameters correspond to an additional theoretical error of less than 1 GeV on $m_{\rm t}$, less than 0.001 on $\alpha_s(m_{\rm Z}^2)$ and 0.1 on $\log(m_{\rm H})$. The error on $\alpha_s(m_{\rm Z}^2)$ covers missing higher-order electroweak corrections and uncertainties in the interplay of electroweak and QCD corrections but not the effect of missing higher-order QCD corrections. A discussion of theoretical uncertainties in the determination of α_s can be found in [66,68]. These theoretical errors have been neglected for the results presented in Tables 21 and 22. Although the theoretical error on $\log(m_{\rm H})$ is still smaller than the experimental error, it is significantly larger than the theoretical error on $m_{\rm t}$ or $\alpha_s(m_{\rm Z}^2)$. More studies on the effect would be welcome.

At present the impact of theoretical uncertainties on the determination of $m_{\rm t}$ from precise electroweak measurements is small compared to the error due to the uncertainty in the value of the fine structure constant $\alpha(m_{\rm Z}^2)$. The uncertainty in $\alpha(m_{\rm Z}^2)$ arises from the contribution of light quarks to the photon vacuum polarization. Recently there have been several reevaluations of $\alpha(m_{\rm Z}^2)$ [69–72]. For the results presented in this Section, a value of $\alpha(m_{\rm Z}^2) = 1/(128.896 \pm 0.090)$ [71] is used. This uncertainty causes an error of 0.00023 on the Standard Model prediction of $\sin^2\theta^{\rm lept}_{\rm eff}$, and an error of 4 GeV on $m_{\rm t}$ (for fixed $m_{\rm H}$), which are included in the results listed in Table 21. The effect on the Standard Model prediction for $\Gamma_{\ell\ell}$ is negligible. The $\alpha_s(m_{\rm Z}^2)$ values for the Standard Model fits presented in this Section are stable against a variation of $\alpha(m_{\rm Z}^2)$ in the interval quoted. For the fits with the Higgs mass left free (see Table 22, the error is 1 GeV on $m_{\rm t}$ and 0.2 on $\log(m_{\rm H})$, which are also included in the results.

	Measurement with Total Error	Standard Model	Pull
a) LEP			
line-shape and lepton asymmetries:			
m_Z [GeV]	91.1863 ± 0.0020	91.1861	0.1
Γ_Z [GeV]	2.4946 ± 0.0027	2.4960	-0.5
σ_h^0 [nb]	41.508 ± 0.056	41.465	0.8
R_ℓ	20.778 ± 0.029	20.757	0.7
$A_{FB}^{0,\ell}$	0.0174 ± 0.0010	0.0159	1.4
+ correlation matrix Table 8			
τ polarization:			
$\mathcal{A}_\tau$	0.1401 ± 0.0067	0.1458	-0.9
$\mathcal{A}_e$	0.1382 ± 0.0076	0.1458	-1.0
b and c quark results:			
R_b [a]	0.2179 ± 0.0012	0.2158	1.8
R_c [a]	0.1715 ± 0.0056	0.1723	-0.1
$A_{FB}^{0,b}$ [a]	0.0979 ± 0.0023	0.1022	-1.8
$A_{FB}^{0,c}$ [a]	0.0733 ± 0.0049	0.0730	0.1
+ correlation matrix Table 13			
$q\bar{q}$ charge asymmetry:			
$\sin^2\theta_{eff}^{lept}$ ($\langle Q_{FB} \rangle$)	0.2320 ± 0.0010	0.23167	0.3
b) SLD			
$\sin^2\theta_{eff}^{lept}$ (A_{LR} [62,63])	0.23061 ± 0.00047	0.23167	-2.2
R_b [6] [a]	0.2149 ± 0.0038	0.2158	-0.2
$\mathcal{A}_b$ [7]	0.863 ± 0.049	0.935	-1.4
$\mathcal{A}_c$ [7]	0.625 ± 0.084	0.667	-0.5
c) $p\bar{p}$ and νN			
m_W [GeV] ($p\bar{p}$ [73])	80.356 ± 0.125	80.353	0.3
$1 - m_W^2/m_Z^2$ (νN [8–10])	0.2244 ± 0.0042	0.2235	0.2
m_t [GeV] ($p\bar{p}$ [15–17])	175 ± 6	172	0.5

Table 20: Summary of measurements included in the combined analysis of Standard Model parameters. Section a) summarizes LEP averages, Section b) SLD results for $\sin^2\theta_{eff}^{lept}$ from the measurement of the left-right polarization asymmetry, for R_b and for $\mathcal{A}_b$ and $\mathcal{A}_c$ from polarized forward-backward asymmetries and Section c) electroweak precision measurements from $p\bar{p}$ colliders and νN scattering. The Standard Model results in column 2 and the pulls (difference between measurement and fit in units of the total measurement error) in column 3 are derived from the Standard Model fit including all data (Table 22, column 2) with the Higgs mass treated as a free parameter.

[a] For fits which combine LEP and SLD heavy flavor measurements we use as input the heavy flavor results given in Equation (11) and their correlation matrix in Table 14 in Section 4 of this note.

	LEP	LEP + SLD	LEP + SLD + $p\bar{p}$ and νN data
m_t (GeV)	$171 \pm 8 \,^{+17}_{-19}$	$177 \,^{+7}_{-8} \,^{+17}_{-19}$	$177 \pm 7 \,^{+16}_{-19}$
$\alpha_s(m_Z^2)$	$0.122 \pm 0.003 \pm 0.002$	$0.121 \pm 0.003 \pm 0.002$	$0.121 \pm 0.003 \pm 0.002$
χ^2/d.o.f.	10/9	20/12	20/14
$\sin^2\theta_{\rm eff}^{\rm lept}$	$0.23209 \pm 0.00024 \,^{+0.00007}_{-0.00016}$	$0.23179 \pm 0.00022 \,^{+0.00006}_{-0.00013}$	$0.23179 \pm 0.00020 \,^{+0.00006}_{-0.00014}$
$1 - m_W^2/m_Z^2$	$0.2247 \pm 0.0009 \,^{+0.0003}_{-0.0002}$	$0.2238 \pm 0.0008 \,^{+0.0004}_{-0.0002}$	$0.2238 \pm 0.0008 \,^{+0.0003}_{-0.0002}$
m_W (GeV)	$80.292 \pm 0.048 \,^{+0.010}_{-0.018}$	$80.337 \pm 0.041 \,^{+0.010}_{-0.021}$	$80.338 \pm 0.040 \,^{+0.009}_{-0.018}$

Table 21: Results of fits to LEP and other electroweak precision data for m_t and $\alpha_s(m_Z^2)$. No external constraint on $\alpha_s(m_Z^2)$ has been imposed. The first column presents the results obtained using LEP data only (Table 20a). The second column gives the result when the SLD measurements of the left-right asymmetry and electroweak heavy flavor results (Table 20b) are also added. In the third column also the combined data from $p\bar{p}$ colliders and νN experiments (Table 20c) are included. The central values and the first errors quoted refer to $m_H = 300$ GeV. The second errors correspond to the variation of the central value when varying m_H in the interval 60 GeV $\leq m_H \leq$ 1000 GeV. The bottom part of the table lists derived results for $\sin^2\theta_{\rm eff}^{\rm lept}$, $1 - m_W^2/m_Z^2$ and m_W.

Table 21 shows the constraints obtained on m_t and $\alpha_s(m_Z^2)$ when fitting the measurements in Table 20 to up-to-date Standard Model calculations [67]. The fits have been repeated for $m_H = 60$, 300 and 1000 GeV and the difference in the fitted parameters is quoted as the second uncertainty. The results obtained using only LEP data (Table 20a), as well as those obtained by including preliminary results from the SLD collaboration (Table 20b) for the left-right asymmetry, A_{LR} [62,63], and the measurements with heavy flavors (R_b [6], $\mathcal{A}_b$ and $\mathcal{A}_c$ [7]) are shown in Table 21. The right-most column of Table 21 gives the Standard Model constraints obtained by including in addition the published measurements of m_W from UA2 [11], CDF [12,13], a preliminary result for m_W from DØ [14][4], and the measurements of the neutrino neutral to charged current ratios from CDHS [8], CHARM [9] and CCFR [74] (Table 20c).

The χ^2/d.o.f. values for all these fits have probabilities ranging from 6% to 33%. In our previous report [1], the measurements of R_b and R_c contributed a χ^2 of approximately 15 for all the Standard Model fits in Table 21. For the new data set reported on here, this is no longer the case (see Table 20).

Figure 6 shows a comparison of the leptonic partial width from LEP (Table 9) and the effective electroweak mixing angle from asymmetries measured at LEP and SLD (Table 19), with the Standard Model. Good agreement with the Standard Model prediction is observed. The star shows the prediction if, among the electroweak radiative corrections only the photon vacuum polarization is included, showing evidence that LEP/SLD data are sensitive to genuine electroweak corrections. Note that the error on $\alpha(m_Z^2)$ (shown as the length of the arrow attached to the star) is as large as the error on $\sin^2\theta_{\rm eff}^{\rm lept}$ from LEP/SLD.

The value of $\alpha_s(m_Z^2)$ derived from an analysis of electroweak precision tests within the Standard Model framework depends essentially on R_ℓ, Γ_Z and σ_h^0. The result is in very good agreement with that obtained from event-shape measurements at LEP ($\alpha_s(m_Z^2) = 0.123 \pm 0.006$ [65]) and of similar

[4]See Reference 73 for a combination of these m_W measurements.

precision. The strong coupling constant can also be determined from the parameter R_ℓ alone. For $m_Z = 91.1863$ GeV, and imposing $m_t = 175 \pm 6$ GeV as a constraint, $\alpha_s = 0.124 \pm 0.004 \pm 0.002$ is obtained, where the second error accounts for the variation of the result when varying m_H in the range 60 GeV $\leq m_H \leq$ 1000 GeV.

In Figure 7 the fitted result for R_b with R_c fixed to its Standard Model value is plotted versus $\sin^2\theta^{lept}_{eff}$. If one assumes the Standard Model dependence of the partial widths on $\sin^2\theta^{lept}_{eff}$ for the light quarks and the c quark, and takes $\alpha_s(m_Z^2) = 0.123 \pm 0.006$, R_ℓ imposes a constraint on the two variables. A good agreement is seen for these 3 experimentally independent measurements, showing the consistency of the LEP data.

The fitted value of m_t is in excellent agreement with the top mass value $m_t = 175 \pm 6$ GeV reported [15–17] by the CDF and DØ collaborations. Note, however, that there is a large correlation between the top mass and the Higgs mass. This can be easily seen in the large variation (36 GeV) of the top mass when changing the Higgs mass between 60 and 1000 GeV. This large correlation is due to the fact that most of the observables listed in Table 20 are sensitive to both m_t and m_H. With the direct measurement of m_t that is now available, it should be possible to constrain m_H.

To constrain m_H, we first perform a fit to the LEP data alone as in Table 21, but fitting as well the Higgs mass. The result is shown in Table 22, column 1. This fit shows that the LEP data prefer a light top quark and a light Higgs, albeit with very large errors. The strongly asymmetric errors on m_H are due to the fact that to first order, the radiative corrections in the Standard Model are proportional to $\log(m_H)$. The correlation between the top quark mass and the Higgs mass is 0.78. It should be noted that the correlation would be even larger if the R_b measurement is not used, as R_b is insensitive to m_H. We then perform a second fit to all data, including the TEVATRON top mass result, which is shown in column 2 of Table 22. As can be expected, both m_t and m_H increase. The correlation between m_t and m_H is reduced somewhat to 0.64. This can also be seen in Figure 8 which shows the contours in m_t and m_H for these two fits.

	LEP	LEP +SLD + $p\overline{p}$ and νN data + m_t
m_t [GeV]	155^{+18}_{-13}	172 ± 6
m_H [GeV]	86^{+202}_{-51}	149^{+148}_{-82}
$\log(m_H)$	$1.93^{+0.52}_{-0.39}$	$2.17^{+0.30}_{-0.35}$
$\alpha_s(m_Z^2)$	0.121 ± 0.003	0.120 ± 0.003
χ^2/d.o.f.	9/8	19/14

Table 22: Results of the fits to LEP data alone and to all data including the to quark mass determination. As the sensitivity to m_H is logarithmic, both the m_H as well as $\log(m_H)$ are quoted.

In Figure 9 the observed value of $\Delta\chi^2 \equiv \chi^2 - \chi^2_{min}$ as a function of m_H is plotted for the fit shown including the CDF/DØ m_t measurement. The shaded band shows the additional error due to the missing higher order corrections. Taking this error into account yields the 95% confidence level upper limit on m_H of 550 GeV.

7 Conclusions

The combination of the many precise electroweak results yields stringent constraints on the Standard Model. All LEP measurements agree well with the predictions. Including all measurements, the data show some sensitivity to the Higgs mass.

The LEP experiments wish to stress that this report reflects a preliminary status at the time of the 1996 summer conferences. A definitive statement on these results has to wait for publication by each collaboration.

Acknowledgments

We would like to thank the SL Division of CERN for the efficient operation of the LEP accelerator, the precise information on the absolute energy scale and their close cooperation with the four experiments. We would also like to thank members of the CDF, DØ and SLD Collaborations for making results available to us in advance of the conferences and for useful discussions concerning their combination.

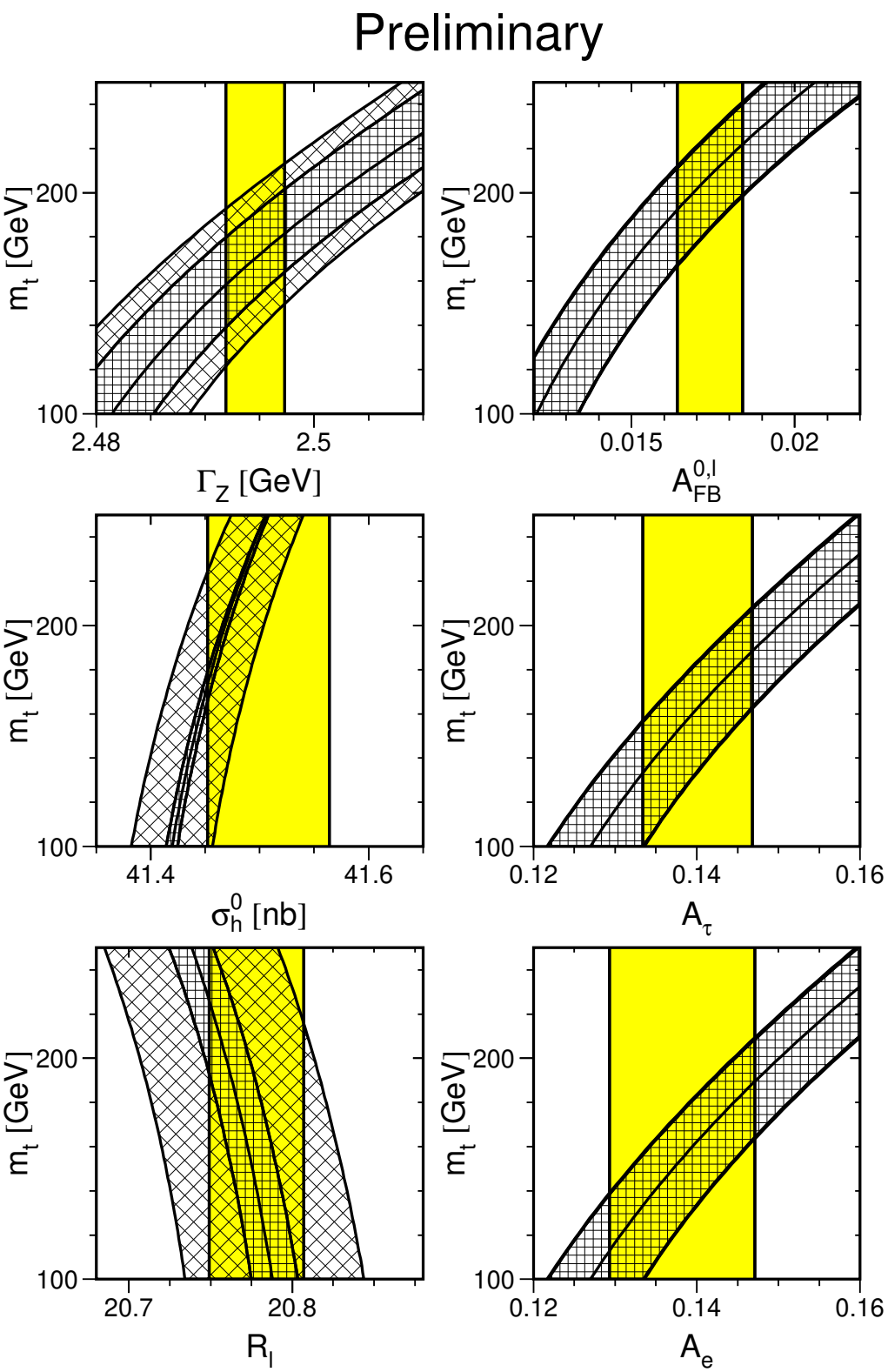

Figure 3: Comparison of LEP measurements with the Standard Model prediction as a function of m_t. The cross-hatch pattern parallel to the axes indicates the variation of the Standard Model prediction with m_H spanning the interval 60 GeV $\leq m_H \leq$ 1000 GeV, and the diagonal cross-hatch pattern corresponds to a variation of $\alpha_s(m_Z^2)$ within the interval $\alpha_s(m_Z^2) = 0.123 \pm 0.006$. The total width of the band corresponds to the linear sum of both uncertainties. The experimental errors on the parameters are indicated as vertical bands.

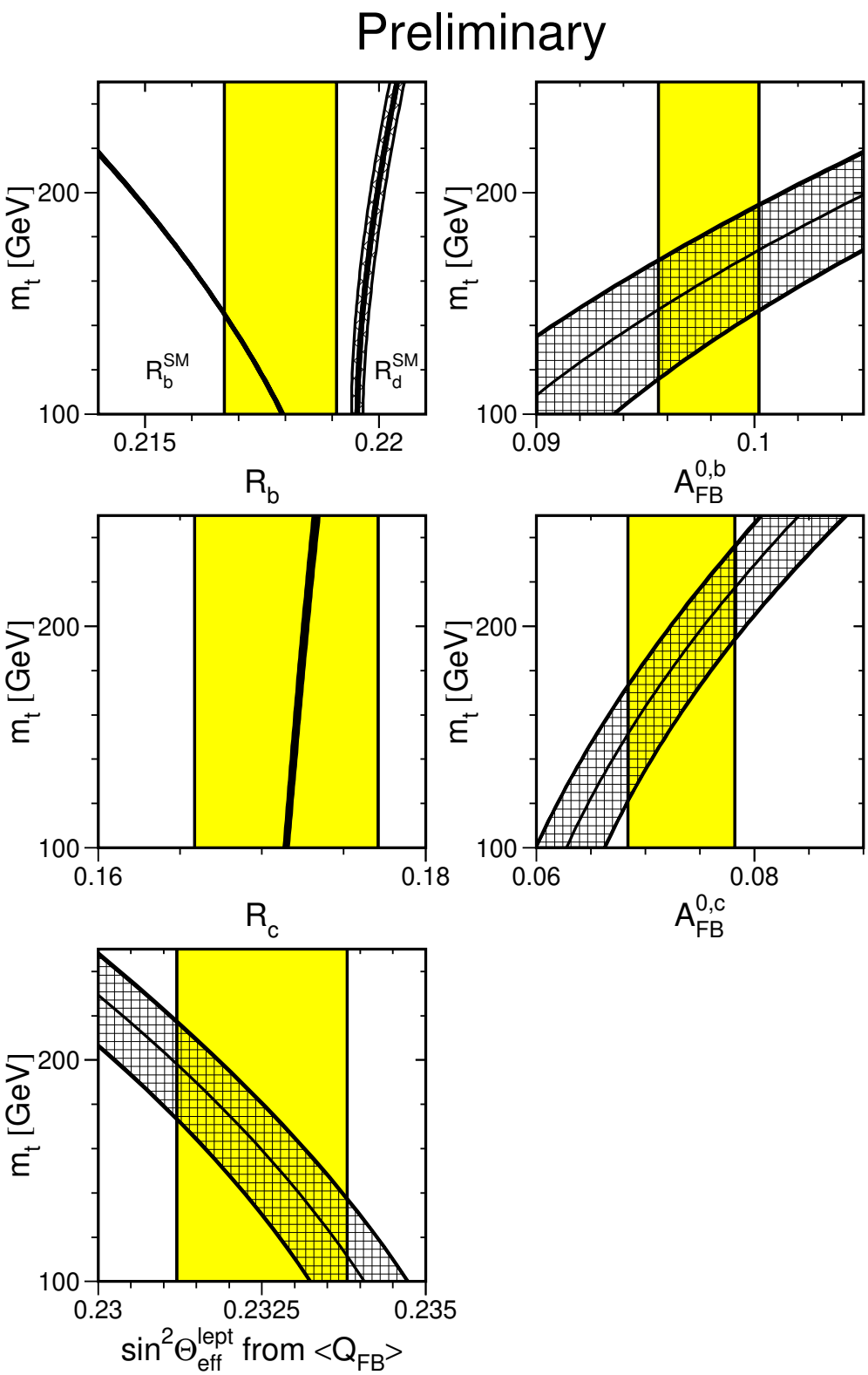

Figure 4: Comparison of LEP measurements with the Standard Model prediction as a function of m_{t} (*c.f.* Figure 3). For ratios of hadronic partial widths the variations with m_{H} and $\alpha_s(m_{\mathrm{Z}}^2)$ nearly cancel. For the comparison of R_{b} with the Standard Model the value of R_{c} has been fixed to its Standard Model prediction. To illustrate the impact of special vertex corrections to R_{b} the Standard Model prediction for R_{d} is also indicated.

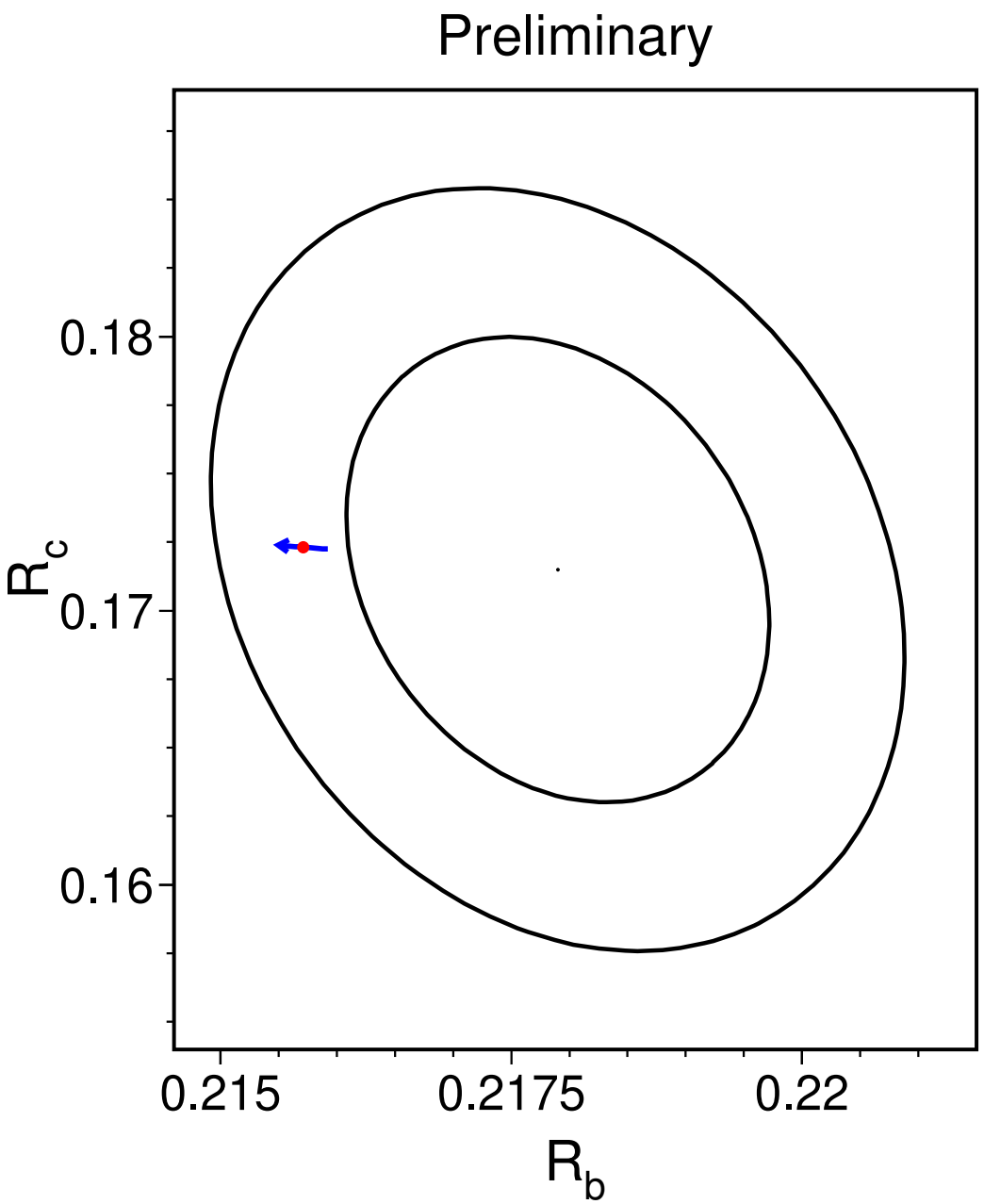

Figure 5: Contours in the R_b-R_c plane derived from LEP data, corresponding to 68% and 95% confidence levels assuming Gaussian systematic errors. The Standard Model prediction for m_t = 175 $\pm$ 6 GeV is also shown. The arrow points in the direction of increasing values of m_t.

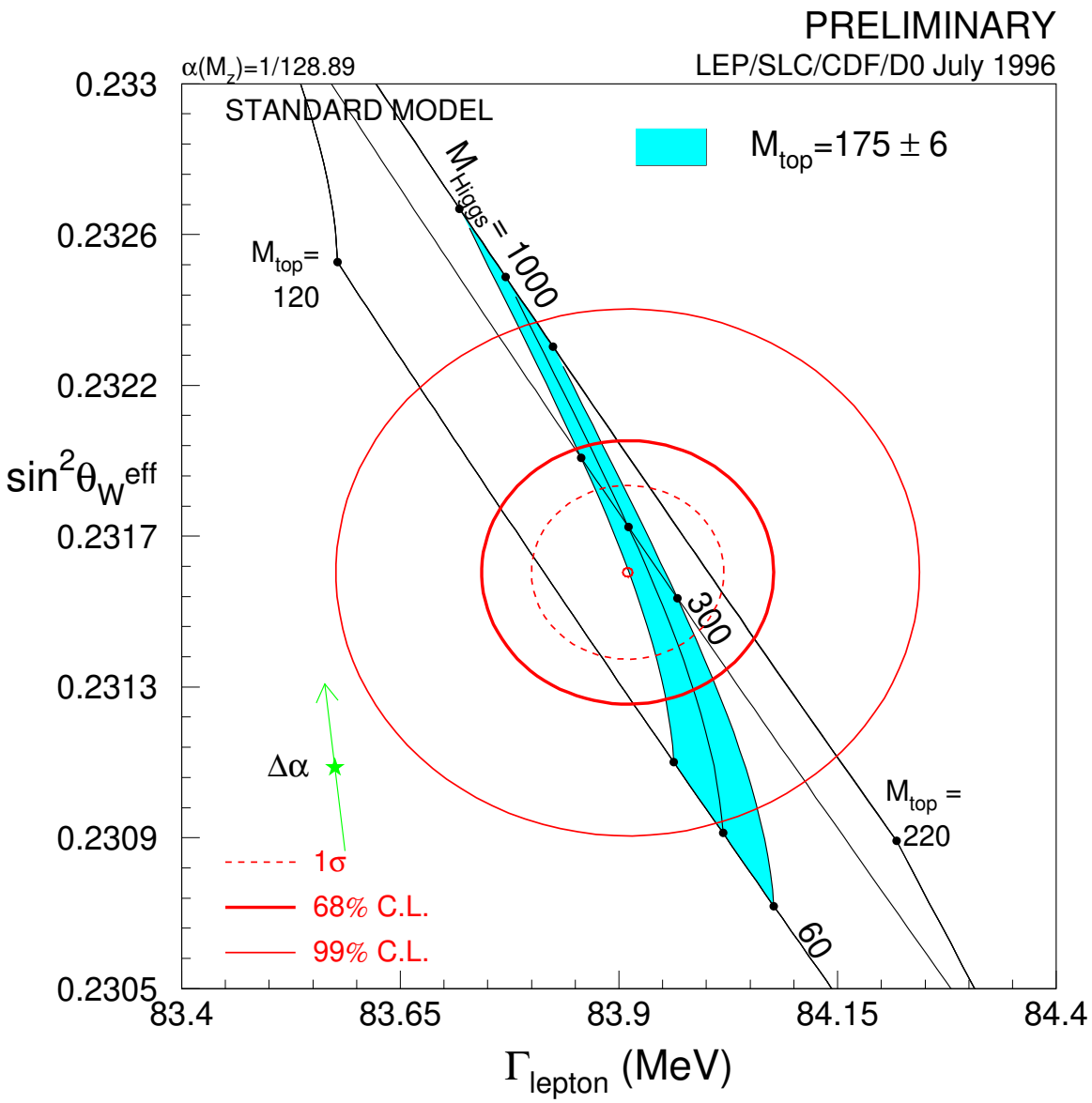

Figure 6: The LEP/SLD measurements of $\sin^2\theta^{\rm lept}_{\rm eff}$ (Table 19) and $\Gamma_{\ell\ell}$ (Table 9) and the Standard Model prediction. The star shows the predictions if among the electroweak radiative corrections only the photon vacuum polarization is included. The corresponding arrow shows variation of this prediction if $\alpha(m_{\rm Z}^2)$ is changing by one standard deviation. This variation gives an additional uncertainty to the Standard Model prediction shown in the figure.

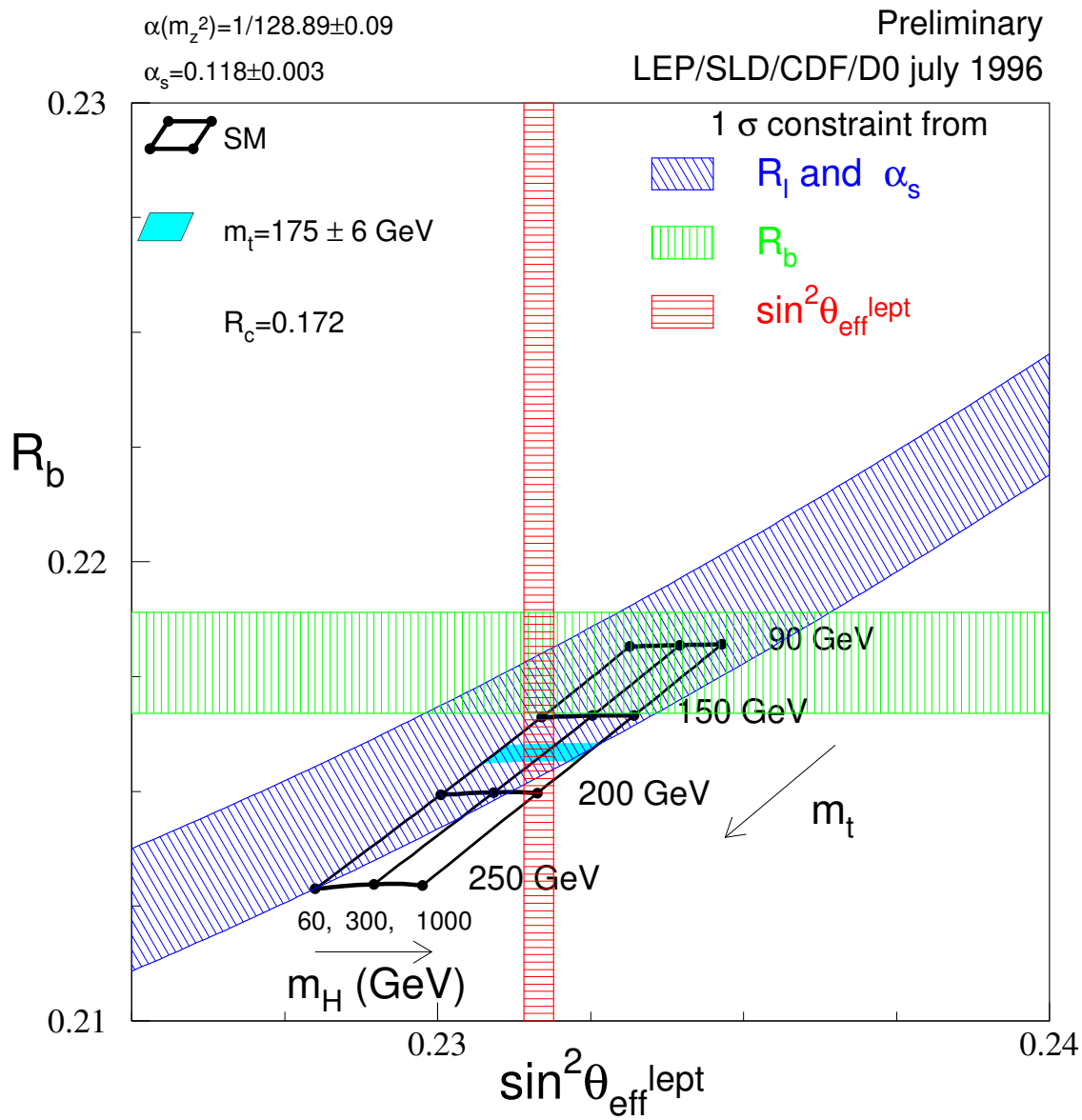

Figure 7: The LEP/SLD measurements of $\sin^2\theta_{\rm eff}^{\rm lept}$ (Table 19) and $R_{\rm b}(R_{\rm c} = 0.172)$ and the Standard Model prediction. Also shown is the constraint resulting from the measurement of R_ℓ on these variables, assuming $\alpha_s(m_{\rm Z}^2) = 0.123 \pm 0.006$, as well as the Standard Model dependence of light-quark partial widths on $\sin^2\theta_{\rm eff}^{\rm lept}$. The Standard Model value for $R_{\rm c}$ is assumed.

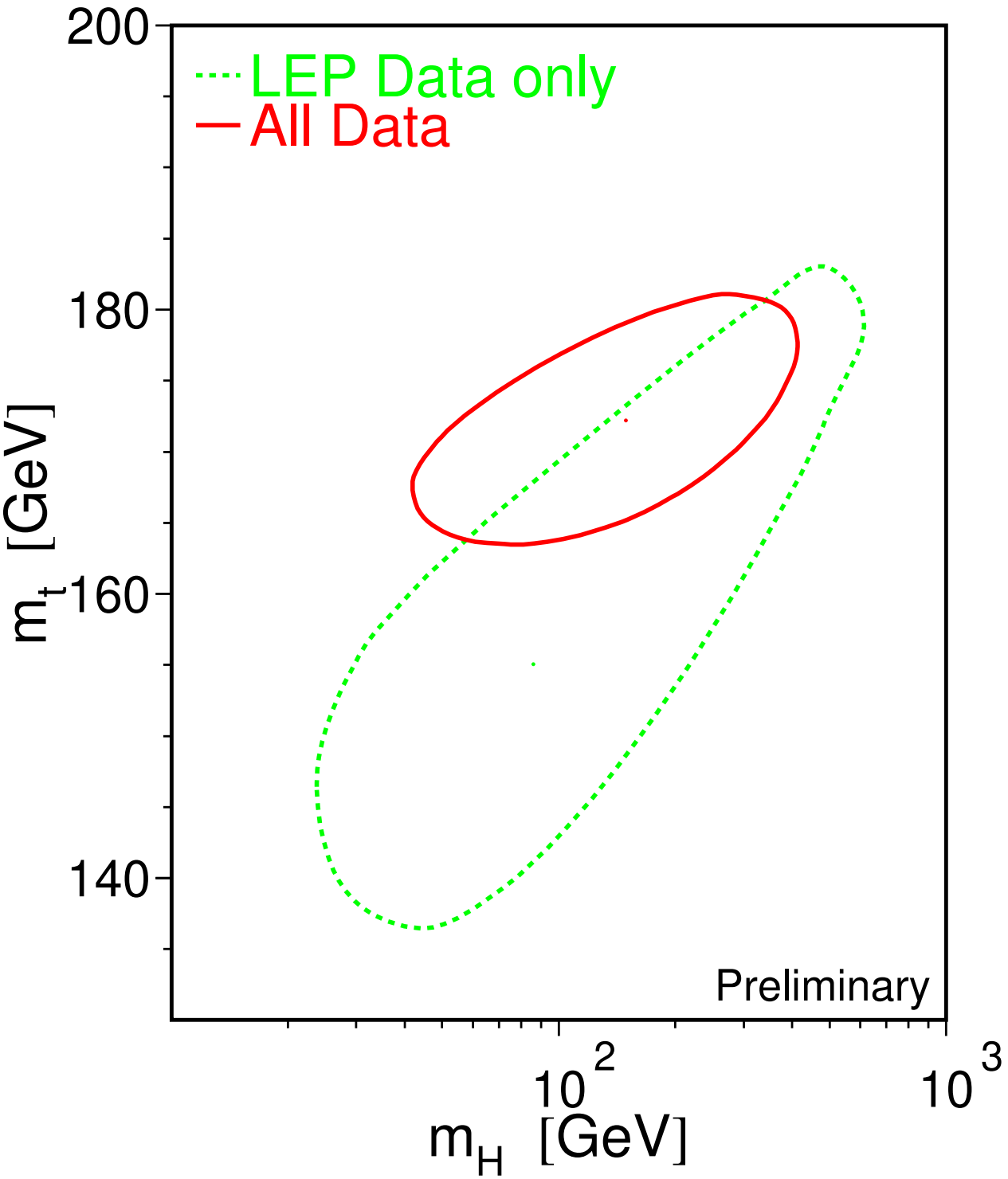

Figure 8: The 68% confidence level contours in m_t and m_H for the fits to LEP data only (dashed curve) and to all data including the CDF/DØ m_t measurement (solid curve).

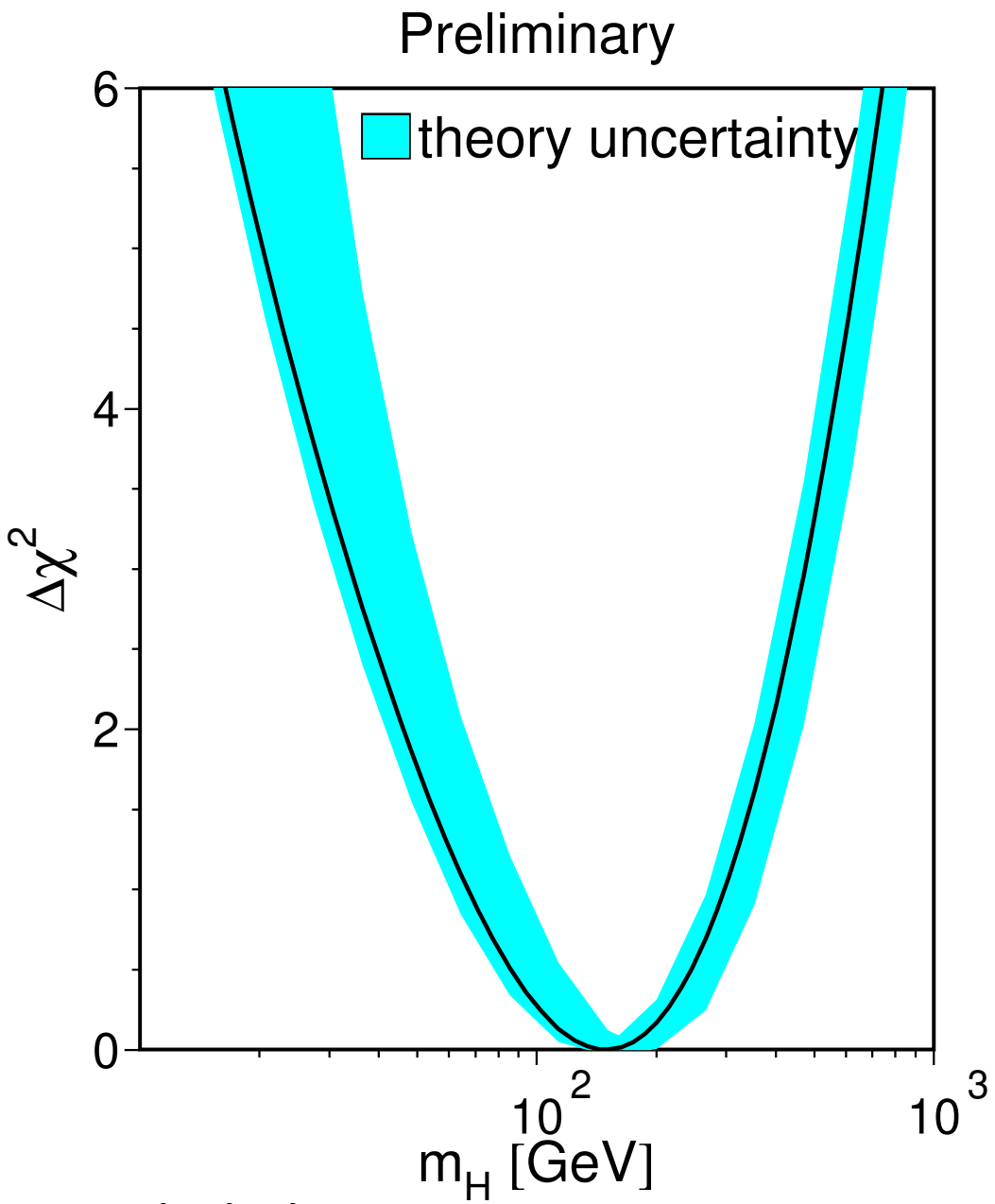

Figure 9: $\Delta\chi^2 = \chi^2 - \chi^2_{min}$ vs $m_{\rm H}$ curves. Continuous line: using all data (last column of Table 22); the band represents the theoretical error due to missing higher order corrections.

Appendix

Heavy Flavor fit including off peak asymmetries

The full 17 parameter fit to the LEP and SLD data gave the following results:

$$
\begin{aligned}
R_{\mathrm{b}} &= 0.2175 \pm 0.0011 \\
R_{\mathrm{c}} &= 0.1717 \pm 0.0056 \\
A_{\mathrm{FB}}^{\mathrm{b\bar{b}}}(-2) &= 0.051 \pm 0.011 \\
A_{\mathrm{FB}}^{\mathrm{c\bar{c}}}(-2) &= -0.038 \pm 0.019 \\
A_{\mathrm{FB}}^{\mathrm{b\bar{b}}}(\mathrm{pk}) &= 0.0961 \pm 0.0024 \\
A_{\mathrm{FB}}^{\mathrm{c\bar{c}}}(\mathrm{pk}) &= 0.0674 \pm 0.0050 \\
A_{\mathrm{FB}}^{\mathrm{b\bar{b}}}(+2) &= 0.110 \pm 0.009 \\
A_{\mathrm{FB}}^{\mathrm{c\bar{c}}}(+2) &= 0.138 \pm 0.016 \\
\mathcal{A}_{\mathrm{b}} &= 0.862 \pm 0.049 \\
\mathcal{A}_{\mathrm{c}} &= 0.627 \pm 0.085 \\
\mathrm{BR}(\mathrm{b} \to \ell) &= 0.1122 \pm 0.0021 \\
\mathrm{BR}(\mathrm{b} \to \mathrm{c} \to \bar{\ell}) &= 0.0803 \pm 0.0034 \\
\overline{\chi} &= 0.1217 \pm 0.0046 \\
f(\mathrm{D}^{+}) &= 0.222 \pm 0.021 \\
f(\mathrm{D_s}) &= 0.116 \pm 0.028 \\
f(\mathrm{c_{baryon}}) &= 0.082 \pm 0.022 \\
\mathrm{P}(\mathrm{c} \to \mathrm{D}^{*+}) \times \mathrm{BR}(\mathrm{D}^{*+} \to \pi^{+}\mathrm{D}^{0}) &= 0.1626 \pm 0.0066
\end{aligned}
$$

with a χ^2/dof of $49/(87-17)$. The corresponding correlation matrix is given in Table 23. The energy for the peak-2, peak and peak+2 results are respectively 89.55 GeV, 91.26 GeVand 92.94 GeV.

	1) R_b	2) R_c	3) $A_{FB}^{b\bar{b}}$ (−2)	4) $A_{FB}^{c\bar{c}}$ (−2)	5) $A_{FB}^{b\bar{b}}$ (pk)	6) $A_{FB}^{c\bar{c}}$ (pk)	7) $A_{FB}^{b\bar{b}}$ (+2)	8) $A_{FB}^{c\bar{c}}$ (+2)	9) $\mathcal{A}_b$	10) $\mathcal{A}_c$	11) BR (1)	12) BR (2)	13) $\overline{\chi}$	14) $f(D^+)$	15) $f(D_s)$	16) $f(c_{bar.})$	17) PcDst
1)	1.00	–0.23	0.00	–0.01	0.00	0.01	–0.01	0.00	–0.03	0.01	–0.13	–0.02	–0.01	–0.14	–0.11	0.13	0.16
2)	–0.23	1.00	0.00	0.01	0.04	–0.06	0.01	–0.04	0.05	–0.07	0.06	0.12	0.00	–0.09	0.14	0.12	–0.60
3)	0.00	0.00	1.00	0.13	0.03	0.01	0.01	0.00	0.01	0.01	0.03	–0.03	0.09	0.00	0.00	0.00	0.00
4)	–0.01	0.01	0.13	1.00	0.02	0.02	0.01	0.00	0.00	0.00	0.01	–0.02	0.01	0.00	0.00	0.00	0.00
5)	0.00	0.04	0.03	0.02	1.00	0.10	0.07	–0.01	0.04	0.02	0.03	–0.14	0.24	0.00	0.00	0.00	–0.03
6)	0.01	–0.06	0.01	0.02	0.10	1.00	0.00	0.14	0.01	0.10	0.12	–0.21	0.12	0.00	–0.01	0.00	0.05
7)	–0.01	0.01	0.01	0.01	0.07	0.00	1.00	0.13	0.02	0.00	0.00	–0.04	0.11	0.00	0.00	0.00	0.00
8)	0.00	–0.04	0.00	0.00	–0.01	0.14	0.13	1.00	0.01	0.05	0.05	–0.08	0.04	–0.01	–0.01	0.01	0.02
9)	–0.03	0.05	0.01	0.00	0.04	0.01	0.02	0.01	1.00	0.12	0.02	–0.05	0.11	–0.01	0.01	0.01	–0.04
10)	0.01	–0.07	0.01	0.00	0.02	0.10	0.00	0.05	0.12	1.00	0.08	–0.23	0.11	0.00	–0.01	–0.01	0.04
11)	–0.13	0.06	0.03	0.01	0.03	0.12	0.00	0.05	0.02	0.08	1.00	–0.16	0.30	0.01	0.02	–0.01	–0.04
12)	–0.02	0.12	–0.03	–0.02	–0.14	–0.21	–0.04	–0.08	–0.05	–0.23	–0.16	1.00	–0.34	0.00	0.02	0.00	–0.07
13)	–0.01	0.00	0.09	0.01	0.24	0.12	0.11	0.04	0.11	0.11	0.30	–0.34	1.00	0.00	0.00	–0.01	0.00
14)	–0.14	–0.09	0.00	0.00	0.00	0.00	0.00	–0.01	–0.01	0.00	0.01	0.00	0.00	1.00	–0.30	–0.19	0.07
15)	–0.11	0.14	0.00	0.00	0.00	–0.01	0.00	–0.01	0.01	–0.01	0.02	0.02	0.00	–0.30	1.00	–0.26	–0.07
16)	0.13	0.12	0.00	0.00	0.00	0.00	0.00	0.01	0.01	–0.01	–0.01	0.00	–0.01	–0.19	–0.26	1.00	–0.05
17)	0.16	–0.60	0.00	0.00	–0.03	0.05	0.00	0.02	–0.04	0.04	–0.04	–0.07	0.00	0.07	–0.07	–0.05	1.00

Table 23: The correlation matrix for the set of the 17 heavy flavor parameters. BR(1) and BR(2) denote BR($b \to \ell$) and BR($b \to c \to \bar{\ell}$) respectively, PcDst denotes P($c \to D^{*+}$) $\times$ BR($D^{*+} \to \pi^+ D^0$).

The Measurements used in the Heavy Flavor Averages

In the following tables, preliminary results are indicated by the symbol †. The values of centre-of-mass energy are given where relevant. In each table, the result used as input to the average procedure is given followed by the statistical error, the correlated and uncorrelated systematic errors, the total systematic error, and any dependence on other electroweak parameters. In the case of the asymmetries the QCD corrected result moved to a common energy (89.55 GeV, 91.26 GeVand 92.94 GeVrespectively for peak-2, peak and peak+2 results) is quoted as *corrected* asymmetry. The asymmetries quoted with a ‡are not QCD corrected.

Contributions to the correlated systematic error quoted here are from any sources of error shared with one or more other results from different experiment in the same table, and the uncorrelated errors from the remaining sources. In the case of $\mathcal{A}_c$ and $\mathcal{A}_b$ from SLD the quoted correlated systematic error has contribution from any source shared with one or more other measurements from LEP experiment. Constants such as $a(x)$ denote the dependence on the assumed value of $x^{\rm used}$, which is also given.

	ALEPH		DELPHI		L3			OPAL		SLD
Tagging	92-95† multiple [39]	90-91 lepton [32]	90-94† multiple [40]	91-92 lepton [34]	lifetime [41]	91 shape [38]	90-91† lepton [36]	92-94† multiple [42]	90-91 lepton [37]	93-95† lifetime [6]
R_b	0.2158	0.2162	0.2202	0.2146	0.2185	0.2220	0.2193	0.2190	0.2240	0.2149
Statistical	0.0009	0.0062	0.0014	0.0089	0.0028	0.0030	0.0081	0.0014	0.0110	0.0032
Uncorrelated	0.0008	0.0040	0.0011	0.0063	0.0027	0.0053	0.0047	0.0012	0.0045	0.0018
Correlated	0.0008	0.0031	0.0013	0.0020	0.0018	0.0036	0.0021	0.0016	0.0045	0.0011
Total Systematic	0.0011	0.0050	0.0017	0.0066	0.0032	0.0064	0.0051	0.0020	0.0063	0.0021
$a(R_c)$	-0.0033		-0.0172		-0.0251	-0.0209	-0.0232	-0.0184	-0.0132	-0.0093
$R_c^{\rm used}$	0.1710		0.1720		0.1710	0.1710	0.1710	0.1710	0.1710	0.1710
$a(\mathrm{BR}(b \to \ell))$						-0.0210				
$\mathrm{BR}(b \to \ell)^{\rm used}$ [%]						10.50				
$a(\mathrm{BR}(b \to c \to \bar{\ell}))$							0.0342			
$\mathrm{BR}(b \to c \to \bar{\ell})^{\rm used}$ [%]							7.90			
$a(f(D^+))$	-0.0011		-0.0045		-0.0066			-0.0043		-0.0009
$f(D^+)^{\rm used}$	0.2310		0.2330		0.2310			0.2310		0.2310
$a(f(D_s))$	-0.0002		-0.0014		0.0003			-0.0005		-0.0003
$f(D_s)^{\rm used}$	0.1100		0.1020		0.1100			0.1100		0.1100
$a(f(c_{\rm baryon}))$	0.0002		0.0006		0.0009			0.0010		0.0004
$f(c_{\rm baryon})^{\rm used}$	0.0630		0.0650		0.0630			0.0630		0.0630

Table 24: The measurements of R_b.

	ALEPH			DELPHI		OPAL	
Tagging	90-91 lepton [32]	91-95† $D^{*\pm}$ [33]	92-95† lepton [33]	91-94† c count + $D^{*\pm}$ [45]	91-92 lepton [34]	90-95† $D^{*\pm}$ [46]	90-94 c count [46]
R_c	0.167	0.172	0.165	0.169	0.164	0.182	0.167
Statistical	0.005	0.010	0.007	0.008	0.008	0.011	0.011
Uncorrelated	0.015	0.009	0.005	0.007	0.017	0.014	0.011
Correlated	0.011	0.005	0.009	0.004	0.0114	0.006	0.005
Total Systematic	0.019	0.010	0.011	0.008	0.020	0.015	0.012

Table 25: The measurements of R_c.

	ALEPH			DELPHI		L3	OPAL		
Tagging	90-95 lepton [32]	90-95 lepton [32]	90-95 lepton [32]	91-94† lepton [35]	91-94† $D^{*\pm}$ [35]	90-93† lepton [36]	91-94 jet [44]	90-95† lepton [37]	90-94† $D^{*\pm}$ [48]
$\sqrt{s}$ (GeV)	88.380	89.380	90.210	89.430	89.540	89.560	89.520	89.490	89.490
$A_{FB}^{b\bar{b}}(-2)$	-3.50	5.44	9.05	6.38	$1.96^{\ddagger}_{\ddagger}$	$7.02^{\ddagger}_{\ddagger}$	6.30	$3.48^{\ddagger}_{\ddagger}$	$-8.60^{\ddagger}_{\ddagger}$
$A_{FB}^{b\bar{b}}(-2)$Corrected	-0.64	5.85	7.52	6.67	2.01	7.11	6.37	3.68	-8.59
Statistical	11.41	2.01	6.06	3.80	10.48	3.50	3.40	1.73	10.80
Uncorrelated	0.10	0.04	0.07	0.16	1.13	0.36	0.21	0.16	2.65
Correlated	0.17	0.05	0.11	0.11	0.11	0.14	0.03	0.04	1.23
Total Systematic	0.20	0.06	0.13	0.19	1.14	0.38	0.21	0.16	2.93
$a(R_b)$	1.2178	-0.2436	-0.9742	-0.7233		-1.6416	-1.8000	-0.1000	
R_b^{used}	0.2192	0.2192	0.2192	0.2170		0.2160	0.2160	0.2155	
$a(R_c)$	-0.8977	0	0.5985	0.1221		0.9126	0.0100	0.1000	
R_c^{used}	0.1710	0.1710	0.1710	0.1710		0.1690	0.1730	0.1720	
$a(A_{FB}^{c\bar{c}}(-2))$	-0.1823	-0.1787	-0.2514			-0.2723	-0.1980		
$A_{FB}^{c\bar{c}}(-2)^{used}$	-2.35	-2.34	-3.35			-2.96	-2.98		
$a(BR(b \to \ell))$	1.1340	-0.2160	-0.8100	-0.9706		-0.9282		0.3406	
$BR(b \to \ell)^{used}$ [%]	11.34	11.34	11.34	11.00		10.50		10.90	
$a(BR(b \to c \to \bar{\ell}))$	0.0629	-0.0943	-0.2201	0.1580		-0.0658		-0.5298	
$BR(b \to c \to \bar{\ell})^{used}$ [%]	7.86	7.86	7.86	7.90		7.90		8.30	
$a(\overline{\chi})$	14.3630	4.3804	9.3371	2.0533					
$\overline{\chi}^{used}$	0.12680	0.12680	0.12680	0.12100					

Table 26: The measurements of $A_{FB}^{b\bar{b}}(-2)$ (in units of 10^{-2}).

	ALEPH		DELPHI			L3	OPAL		
Tagging	90-95 lepton [32]	91-95† jet [43]	91-94† lepton [35]	91-94† $D^{*\pm}$ [35]	91-94† jet [35]	90-93† lepton [36]	91-94† jet [44]	90-95† lepton [37]	90-94† $D^{*\pm}$ [48]
$\sqrt{s}$ (GeV)	91.210	91.187	91.230	91.230	91.230	91.270	91.250	91.240	91.240
$A_{FB}^{b\bar{b}}(pk)$	9.86	9.27	10.76	$7.13^{\ddagger}_{\ddagger}$	9.88	$10.28^{\ddagger}_{\ddagger}$	9.73	$8.81^{\ddagger}_{\ddagger}$	$9.40^{\ddagger}_{\ddagger}$
$A_{FB}^{b\bar{b}}(pk)$Corrected	9.96	9.41	10.82	7.30	9.94	10.42	9.75	8.99	9.59
Statistical	0.45	0.39	0.76	2.50	0.72	1.00	0.67	0.44	2.70
Uncorrelated	0.10	0.26	0.21	1.27	0.36	0.37	0.31	0.14	2.19
Correlated	0.17	0.23	0.23	0.13	0.12	0.15	0.25	0.15	0.47
Total Systematic	0.20	0.34	0.31	1.28	0.38	0.39	0.39	0.20	2.24
$a(R_b)$	-1.4613	-0.1840	-2.8933		-0.6000	-1.6416	-10.0900	-0.7000	
R_b^{used}	0.2192	0.2178	0.2170		0.2210	0.2160	0.2160	0.2155	
$a(R_c)$	1.0260		1.2214		0.2400	0.9126	0.1100	0.6000	
R_c^{used}	0.1710		0.1710		0.1710	0.1690	0.1730	0.1720	
$a(A_{FB}^{c\bar{c}}(pk))$	0.4965					0.5752	0.7156		
$A_{FB}^{c\bar{c}}(pk)^{used}$	6.41					6.25	6.15		
$a(BR(b \to \ell))$	-1.2960		-3.5588			-0.6468		-0.3406	
$BR(b \to \ell)^{used}$ [%]	11.34		11.00			10.50		10.90	
$a(BR(b \to c \to \bar{\ell}))$	-0.1886		0.4740			-0.2831		-0.3532	
$BR(b \to c \to \bar{\ell})^{used}$ [%]	7.86		7.90			7.90		8.30	
$a(\overline{\chi})$	3.3429		3.4467						
$\overline{\chi}^{used}$	0.12680		0.12100						

Table 27: The measurements of $A_{FB}^{b\bar{b}}(pk)$ (in units of 10^{-2}).

	ALEPH			DELPHI		L3	OPAL		
Tagging	90-95 lepton [32]	90-95 lepton [32]	90-95 lepton [32]	91-94† lepton [35]	91-94† $D^{*\pm}$ [35]	90-93† lepton [36]	91-94 jet [44]	90-95† lepton [37]	90-94† $D^{*\pm}$ [48]
$\sqrt{s}$ (GeV)	92.050	92.940	93.900	93.017	92.940	92.930	92.940	92.950	92.950
$A_{FB}^{b\bar{b}}(+2)$	3.91	10.54	8.98	15.24	5.63‡	11.02‡	17.30	10.34‡	-2.10‡
$A_{FB}^{b\bar{b}}(+2)$Corrected	4.99	10.54	8.14	15.16	5.72	11.20	17.30	10.49	-2.14
Statistical	5.17	1.60	7.68	3.60	9.55	2.90	2.90	1.43	9.00
Uncorrelated	0.08	0.15	0.12	0.49	1.64	0.37	0.70	0.25	2.44
Correlated	0.14	0.24	0.24	0.41	0.13	0.15	0.06	0.28	1.26
Total Systematic	0.16	0.29	0.27	0.64	1.65	0.40	0.71	0.37	2.74
$a(R_b)$	-0.9742	-1.9484	-1.8267	-2.8933		-1.6416	-18.3000	-0.8000	
R_b^{used}	0.2192	0.2192	0.2192	0.2170		0.2160	0.2160	0.2155	
$a(R_c)$	0.9191	1.4749	1.5817	-0.9771		0.9126	0.1600	0.8000	
R_c^{used}	0.1710	0.1710	0.1710	0.1710		0.1690	0.1730	0.1720	
$a(A_{FB}^{c\bar{c}}(+2))$	0.9650	0.9708	-0.5035			1.1156	0.6800		
$A_{FB}^{c\bar{c}}(+2)^{used}$	12.48	12.51	-0.65			12.13	12.00		
$a(BR(b \to \ell))$	-0.8640	-1.7280	-1.6200	-3.2353		-0.6300		-1.3625	
$BR(b \to \ell)^{used}$ [%]	11.34	11.34	11.34	11.00		10.50		10.90	
$a(BR(b \to c \to \bar{\ell}))$	-0.1258	-0.2515	-0.2515	0.4740		-0.3028		0.7064	
$BR(b \to c \to \bar{\ell})^{used}$ [%]	7.86	7.86	7.86	7.90		7.90		8.30	
$a(\overline{\chi})$	3.5504	6.5475	9.6137	4.7667					
$\overline{\chi}^{used}$	0.12680	0.12680	0.12680	0.12100					

Table 28: The measurements of $A_{FB}^{b\bar{b}}(+2)$ (in units of 10^{-2}).

	ALEPH	DELPHI	OPAL	
Tagging	91-94† $D^{*\pm}$ [47]	91-94† $D^{*\pm}$ [35]	90-95† lepton [37]	91-94† $D^{*\pm}$ [48]
$\sqrt{s}$ (GeV)	89.400	89.540	89.490	89.490
$A_{FB}^{c\bar{c}}(-2)$	-4.90‡‡	0.20‡‡	-6.81‡‡	3.90‡‡
$A_{FB}^{c\bar{c}}(-2)$Corrected	-4.05	0.26	-6.52	4.30
Statistical	7.60	5.19	2.44	5.10
Uncorrelated	0.85	0.55	0.43	0.81
Correlated	0.06	0.07	0.21	0.31
Total Systematic	0.85	0.56	0.48	0.86
$a(R_b)$			-3.4000	
R_b^{used}			0.2155	
$a(R_c)$			3.2000	
R_c^{used}			0.1720	
$a(A_{FB}^{b\bar{b}}(-2))$	0.2295			
$A_{FB}^{b\bar{b}}(-2)^{used}$	-1.34			
$a(BR(b \to \ell))$			-1.7031	
$BR(b \to \ell)^{used}$ [%]			10.90	
$a(BR(b \to c \to \bar{\ell}))$			-1.4128	
$BR(b \to c \to \bar{\ell})^{used}$ [%]			8.30	

Table 29: The measurements of $A_{FB}^{c\bar{c}}(-2)$ (in units of 10^{-2}).

	ALEPH		DELPHI		L3	OPAL	
Tagging	90-91 lepton [32]	91-94† $D^{*\pm}$ [47]	91-94† lepton [35]	91-94† $D^{*\pm}$ [35]	90-91 lepton [36]	90-95† lepton [37]	90-94† $D^{*\pm}$ [48]
$\sqrt{s}$ (GeV)	91.260	91.200	91.230	91.230	91.240	91.240	91.240
$A_{FB}^{c\bar{c}}$(pk)	9.10‡‡	6.40‡‡	8.42	7.52‡‡	7.84‡‡	5.82‡‡	6.30‡‡
$A_{FB}^{c\bar{c}}$(pk)Corrected	9.19	6.76	8.57	7.74	8.02	5.97	6.46
Statistical	2.00	1.30	1.39	1.21	3.70	0.59	1.20
Uncorrelated	1.56	0.20	0.91	0.55	2.42	0.37	0.53
Correlated	1.05	0.18	0.75	0.12	0.60	0.38	0.16
Total Systematic	1.88	0.27	1.18	0.57	2.50	0.53	0.56
$a(R_b)$			3.6167		4.3200	4.1000	
R_b^{used}			0.2170		0.2160	0.2155	
$a(R_c)$			-6.3514		-6.7600	-3.8000	
R_c^{used}			0.1710		0.1690	0.1720	
$a(A_{FB}^{b\bar{b}}$(pk))		-1.5110			6.4274		
$A_{FB}^{b\bar{b}}$(pk)^{used}		8.81			8.84		
$a(BR(b \to \ell))$			4.8529		3.5007	5.1094	
$BR(b \to \ell)^{used}$ [%]			11.00		10.50	10.90	
$a(BR(b \to c \to \bar{\ell}))$			-3.7920		-3.2917	-1.7660	
$BR(b \to c \to \bar{\ell})^{used}$ [%]			7.90		7.90	8.30	

Table 30: The measurements of $A_{FB}^{c\bar{c}}$(pk) from D^* meson and lepton-tag analyses (in units of 10^{-2}).

	ALEPH	DELPHI	OPAL	
Tagging	91-94† $D^{*\pm}$ [47]	91-94† $D^{*\pm}$ [35]	90-95† lepton [37]	90-94† $D^{*\pm}$ [48]
$\sqrt{s}$ (GeV)	93.000	92.940	92.950	92.950
$A_{FB}^{c\bar{c}}(+2)$	$10.90^{\ddagger}_{\ddagger}$	$7.97^{\ddagger}_{\ddagger}$	$15.43^{\ddagger}_{\ddagger}$	$15.80^{\ddagger}_{\ddagger}$
$A_{FB}^{c\bar{c}}(+2)$Corrected	10.85	8.05	15.56	15.93
Statistical	6.10	4.55	2.02	4.10
Uncorrelated	0.71	0.55	0.90	1.06
Correlated	0.28	0.17	0.38	0.27
Total Systematic	0.77	0.58	0.97	1.09
$a(R_b)$			9.6000	
R_b^{used}			0.2155	
$a(R_c)$			-8.9000	
R_c^{used}			0.1720	
$a(A_{FB}^{b\bar{b}}(+2))$	-2.0639			
$A_{FB}^{b\bar{b}}(+2)^{used}$	12.04			
$a(BR(b \to \ell))$			9.5375	
$BR(b \to \ell)^{used}$ [%]			10.90	
$a(BR(b \to c \to \bar{\ell}))$			-1.5894	
$BR(b \to c \to \bar{\ell})^{used}$ [%]			8.30	

Table 31: The measurements of $A_{FB}^{c\bar{c}}(+2)$ (in units of 10^{-2}).

	SLD		
Tagging	93-95† lepton [7]	93-95† jet [7]	94-95† $K^{\pm}$ [7]
$\sqrt{s}$ (GeV)	91.260	91.260	91.260
$\mathcal{A}_b$	0.882	0.843	0.907
Statistical	0.068	0.046	0.094
Uncorrelated	0.037	0.049	0.092
Correlated	0.021	0.000	0.007
Total Systematic	0.043	0.049	0.092
$a(R_b)$	-0.4302	-0.1308	-0.0218
R_b^{used}	0.2216	0.2180	0.2180
$a(R_c)$	0.0800	0.1328	0.0030
R_c^{used}	0.1600	0.1710	0.1710
$a(\mathcal{A}_c)$		0.0809	-0.1332
$\mathcal{A}_c^{used}$		0.666	0.666
$a(BR(b \to \ell))$	-0.3038		
$BR(b \to \ell)^{used}$ [%]	10.75		
$a(BR(b \to c \to \bar{\ell}))$	0.1095		
$BR(b \to c \to \bar{\ell})^{used}$ [%]	8.10		
$a(\overline{\chi})$	0.4197		0.2229
$\overline{\chi}^{used}$	0.12200		0.13000

Table 32: The measurements of $\mathcal{A}_b$.

	SLD	
Tagging	93-95† lepton [7]	93-95† $D^{*\pm}$ [7]
$\sqrt{s}$ (GeV)	91.260	91.260
$\mathcal{A}_c$	0.612	0.640
Statistical	0.102	0.110
Uncorrelated	0.042	0.053
Correlated	0.050	0.020
Total Systematic	0.065	0.057
$a(R_b)$	0.1173	
R_b^{used}	0.2216	
$a(R_c)$	-0.4864	
R_c^{used}	0.1600	
$a(\mathcal{A}_b)$		-0.1278
$\mathcal{A}_b{}^{used}$		0.935
$a(BR(b \to \ell))$	0.4580	
$BR(b \to \ell)^{used}$ [%]	10.75	
$a(BR(b \to c \to \bar{\ell}))$	-0.4991	
$BR(b \to c \to \bar{\ell})^{used}$ [%]	8.10	

Table 33: The measurements of $\mathcal{A}_c$.

	ALEPH		DELPHI	L3	OPAL
Tagging	90-91 lepton [32]	92-93† multiple [32]	91-92 lepton [34]	90-91† lepton [36]	90-91† lepton [37]
$BR(b \to \ell)$	11.20	11.01	11.30	11.42	10.60
Statistical	0.33	0.10	0.45	0.48	0.60
Uncorrelated	0.32	0.20	0.50	0.30	0.39
Correlated	0.27	0.21	0.46	0.21	0.53
Total Systematic	0.42	0.29	0.68	0.37	0.66
$a(R_c)$				0.6107	0.2236
R_c^{used}				0.1710	0.1710
$a(BR(b \to c \to \bar{\ell}))$				0.4608	
$BR(b \to c \to \bar{\ell})^{used}$ [%]				7.90	
$a(\overline{\chi})$		0.2075			
$\overline{\chi}^{used}$		0.12610			

Table 34: The measurements of $BR(b \to \ell)$ from the lepton-tag analyses.

Tagging	ALEPH 90-91 lepton [32]	ALEPH 92-93† multiple [32]	DELPHI 91-92 lepton [34]	OPAL 90-91† lepton [37]
$BR(b \to c \to \bar{\ell})$	8.81	7.68	7.90	8.40
Statistical	0.25	0.18	0.49	0.40
Uncorrelated	0.40	0.25	0.95	0.57
Correlated	0.69	0.42	0.78	0.38
Total Systematic	0.80	0.49	1.23	0.68
$a(R_c)$				0.3157
R_c^{used}				0.1710
$a(\overline{\chi})$		-0.5108		
$\overline{\chi}^{used}$		0.12610		

Table 35: The measurements of $BR(b \to c \to \bar{\ell})$ from the lepton-tag analyses.

Tagging	ALEPH 90-95† lepton [32]	DELPHI 91-92 lepton [34]	L3 90-93† lepton [36]	OPAL 90-95† lepton [37]
$\overline{\chi}$	0.12461	0.14900	0.12530	0.11390
Statistical	0.00515	0.02000	0.01100	0.00540
Uncorrelated	0.00244	0.01044	0.00516	0.00306
Correlated	0.00403	0.01192	0.00266	0.00324
Total Systematic	0.00471	0.01584	0.00581	0.00446
$a(R_b)$	0.0341		0.0009	
R_b^{used}	0.2192		0.2160	
$a(R_c)$	0.0009		0.0007	
R_c^{used}	0.1710		0.1690	
$a(BR(b \to \ell))$	0.0524		0.0462	0.0170
$BR(b \to \ell)^{used}$ [%]	11.34		10.50	10.90
$a(BR(b \to c \to \ell))$	-0.0440		-0.0342	-0.0318
$BR(b \to c \to \bar{\ell})^{used}$ [%]	7.86		7.90	8.30

Table 36: The measurements of $\overline{\chi}$ from the lepton-tag analyses.

Tagging	DELPHI 91-94† $D^{*\pm}$ [45]	OPAL 90-94† $D^{*\pm}$ [46]
$P(c \to D^{*+}) \times BR(D^{*+} \to \pi^+ D^0)$	0.1678	0.1510
Statistical	0.0069	0.0110
Uncorrelated	0.0065	0.0108
Correlated	0.0011	0.0020
Total Systematic	0.0066	0.0110

Table 37: The measurements of $P(c \to D^{*+}) \times BR(D^{*+} \to \pi^+ D^0)$.

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